

Bifurcation Analysis for Time Steppers

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THE THREE TOOLS OF COMPUTATIONAL FLUID DYNAMICS

Time stepping: $\partial_t U = LU + N(U)$

Steady state solving: $0 = LU + N(U)$

Linear stability analysis: $\lambda u = Lu + N_U u$

Heat Equation

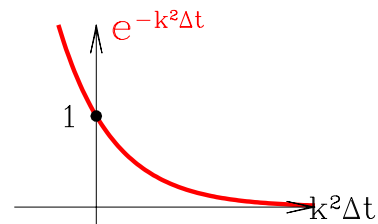
$$\partial_t u = \partial_{xx}^2 u$$

$$u(x, t) = \sum_{k=1}^{k_{max}} u_k(t) \sin kx$$

$$\partial_t u_k = -k^2 u_k$$

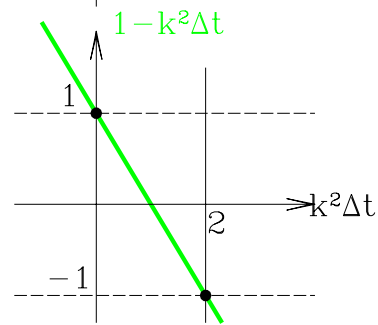
EXACT SOLUTION

$$u_k(t + \Delta t) = e^{-k^2 \Delta t} u_k(t)$$



EXPLICIT EULER

$$\begin{aligned} u_k(t + \Delta t) &= u_k(t) - k^2 \Delta t u_k(t) \\ &= (1 - k^2 \Delta t) u_k(t) \end{aligned}$$



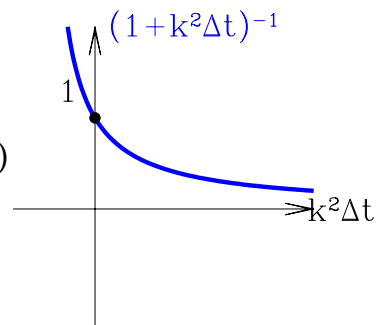
As $k_{max} \rightarrow \infty$, $\Delta t_{max} = \frac{2}{k_{max}^2} \rightarrow 0$

IMPLICIT EULER

$$\begin{aligned} u_k(t + \Delta t) &= u_k(t) \\ &\quad - k^2 \Delta t u_k(t + \Delta t) \\ (1 + k^2 \Delta t) u_k(t + \Delta t) &= u_k(t) \\ u_k(t + \Delta t) &= (1 + k^2 \Delta t)^{-1} u_k(t) \end{aligned}$$

Matrix version:

$$u(t + \Delta t) = (I - \Delta t L)^{-1} u(t)$$



NAVIER-STOKES EQUATIONS

$$\begin{aligned}\partial_t U &= -(U \cdot \nabla)U - \nabla P + \nu \nabla^2 U \\ &= -(I - \nabla \nabla^{-2} \nabla \cdot)(U \cdot \nabla)U + \nu \nabla^2 U \\ &= N(U) + L U\end{aligned}$$

• Time stepping (Direct Numerical Simulation)

$$\partial_t U = N(U) + L U \equiv A(U)$$

Explicit/Implicit Euler:

$$\begin{aligned}U(t + \Delta t) &= U(t) + \Delta t [N(U(t)) + L U(t + \Delta t)] \\ &= (I - \Delta t L)^{-1} (I + \Delta t N) U(t) \equiv B U(t)\end{aligned}$$

• Steady state solving

$$0 = N(U) + L U$$

Newton's method:

$$\begin{cases} A_U u = A(U) \\ U \leftarrow U - u \end{cases}$$

• Linear stability analysis:

$$\lambda u = N_U u + L u \equiv A_U u$$

Arnoldi/block power method:

$$u_{n+1} = A_U^{-1} u_n \quad \text{or} \quad u_{n+1} = e^{A_U \Delta t} u_n$$

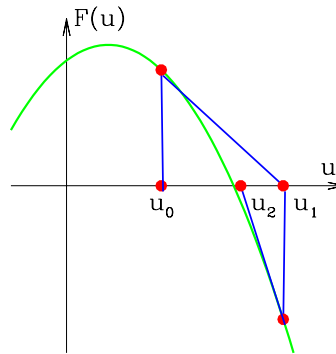
$$N_U u \equiv -(U \cdot \nabla)u - (u \cdot \nabla)U$$

$$A_U u = N_U u + L u$$

STEADY STATE SOLVING

$$0 = F(U)$$

Newton's
method



$$0 = F(U - u) \approx F(U) - DF(U)u$$
$$\begin{cases} DF(U)u = F(U) \\ U \leftarrow U - u \end{cases}$$

How to solve linear systems?

1) **Direct: Gaussian Elimination = LU + Backsolve**

Storage: M^2

Time: M^3 ,

For 3D case with $M_x = M_y = M_z = 10^2$, we have $M = 10^6$

$M^2 = 10^{12}$

$M^3 = 10^{18}$

2) **Iterative: Conjugate Gradient methods**

Use only matrix-vector products $u \rightarrow Au$,

For an arbitrary matrix,

Each product $u \rightarrow Au$ requires M^2 operations
Convergence requires M iterations } M^3

Can gain:

If A is structured or sparse, then $u \rightarrow Au$ takes $\sim M$ ops

If A is well-conditioned, convergence takes few iterations.

CONDITIONING

A is well conditioned if its eigenvalues lie close together.

The best conditioned matrix is a multiple of the identity.

The condition number is, roughly,

$$\kappa(A) \sim \left| \frac{\max \text{ eig of } A}{\min \text{ eig of } A} \right|$$

PRECONDITIONING

$$Au = v$$

$$PAu = Pv$$

where:

P is easy to act with

PA is better conditioned than A

Extreme cases:

$P = I$ (easy to act with but no improvement)

$P = A^{-1}$ (perfect preconditioner but impossible)

Our case:

$$\begin{aligned} A_U u &= L u + N_U u \\ &= \nabla^2 U - (U \cdot \nabla) u - (u \cdot \nabla) U \end{aligned}$$

For 3D case with $M_x = M_y = M_z = 100$,

eigs of L range from ~ -1 to $-(M_x^2 + M_y^2 + M_z^2) = -30000$.

Idea:

L is the main source of difficulty \implies Use $P = L^{-1}$

Question: Where do we get L^{-1} ?

Answer: Already present in a timestepping code!

$$\begin{aligned}U(t + \Delta t) - U(t) &= [(I - \Delta t L)^{-1}(I + \Delta t N) - I] U(t) \\&= (I - \Delta t L)^{-1} [I + \Delta t N - (I - \Delta t L)] U(t) \\&= (I - \Delta t L)^{-1} \Delta t (N + L) U(t)\end{aligned}$$

$(B - I)U(t) \equiv U(t + \Delta t) - U(t)$ has same roots as $(N + L)$!

- In time-stepping, Δt must be small ($\sim 10^{-2}$) to insure $(I - \Delta t L)^{-1}(1 + \Delta t N_U) \approx \exp((L + N_U)\Delta t)$.
- Here, Δt plays only algebraic role, and can (should) be taken large ($\gtrsim 10^2$).
- Δt interpolates between $P = I$ and $P = L^{-1}$.
- Called **Stokes preconditioning**

ONE NEWTON STEP

$$(I - \Delta t L)^{-1} \Delta t (N_U + L) u = (I - \Delta t L)^{-1} \Delta t (N + L) U$$

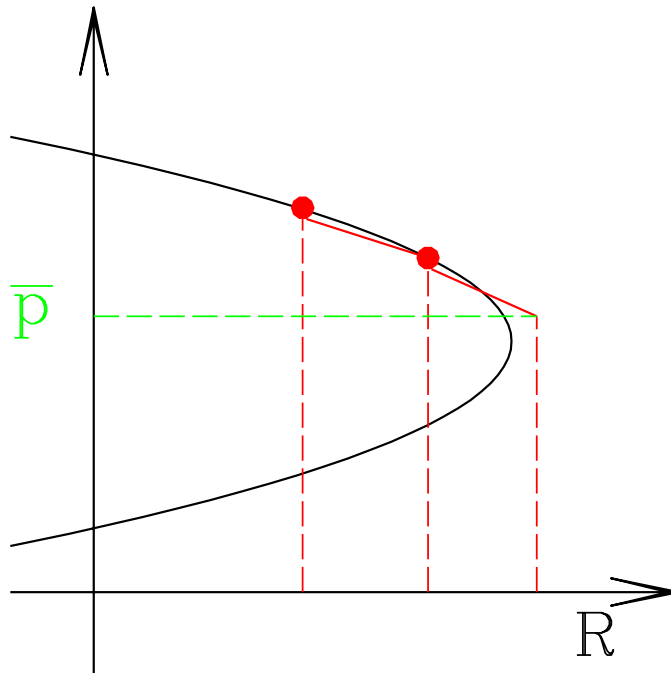
$$\underbrace{[(I - \Delta t L)^{-1}(I + \Delta t N_U) - I] u}_{\substack{\text{difference between two} \\ \text{widely spaced consecutive} \\ \text{linearized timesteps}}} = \underbrace{[(I - \Delta t L)^{-1}(I + \Delta t N) - I] U}_{\substack{\text{difference between two} \\ \text{widely spaced consecutive} \\ \text{timesteps}}}$$

Solve linear system with BI-CGSTAB

H.A. van der Vorst, *Bi-CGSTAB: A fast and smoothly converging variant of Bi-CG for the solution of nonsymmetric linear systems*, **SIAM J. Sci. Stat. Comput.** **13**, 631 (1992)

~ 30 lines of code

CONTINUATION



Goal

$$0 = RN(U) + LU$$

$$0 = p(U, R) - \bar{p} \text{ where } \left\{ \begin{array}{l} U_i \text{ some component} \\ R \end{array} \right\}$$

Newton step

(U, R) not solution, so try $(U - u, R - r)$

$$\begin{aligned} 0 &= (R - r)N(U - u) + L(U - u) \\ &= RN(U) + LU - RN_U u - rN(U) - Lu + O(r, u)^2 \end{aligned}$$

$$0 = p(U - u, R - r) - \bar{p} = \left\{ \begin{array}{l} U_i - \bar{p} - u_i \\ R - \bar{p} - r \end{array} \right\}$$

$$\underbrace{\left[\begin{array}{c|c} RN_U + L & N(U) \\ \hline 0 \ 0 \ \dots \ 0 \ 1 \ 0 \ \dots \ 0 & 1 \end{array} \right]} \begin{bmatrix} u \\ r \end{bmatrix} = \begin{bmatrix} RN(U) + LU \\ \left\{ \begin{array}{c} U_i - \bar{p} \\ R - \bar{p} \end{array} \right\} \end{bmatrix}$$

or

If $p(U, R) = R$ (i.e. set Reynolds number),
then set $R = \bar{p}$, $r = 0$ and get previous case:

$$L^{-1} [RN_U + L] [u] = L^{-1} [RN(U) + LU]$$

If $p(U, R) = U_i$, then must solve extended system for
(u, r).

$$\left[\begin{array}{c|c} L^{-1}(RN_u + L) & L^{-1}N(U) \\ \hline 0 \ 0 \ \dots \ 0 \ 1 \ 0 \ \dots \ 0 & 0 \end{array} \right] \begin{bmatrix} u \\ r \end{bmatrix} = \begin{bmatrix} L^{-1}(RN(U) + LU) \\ U_i - \bar{p} \end{bmatrix}$$

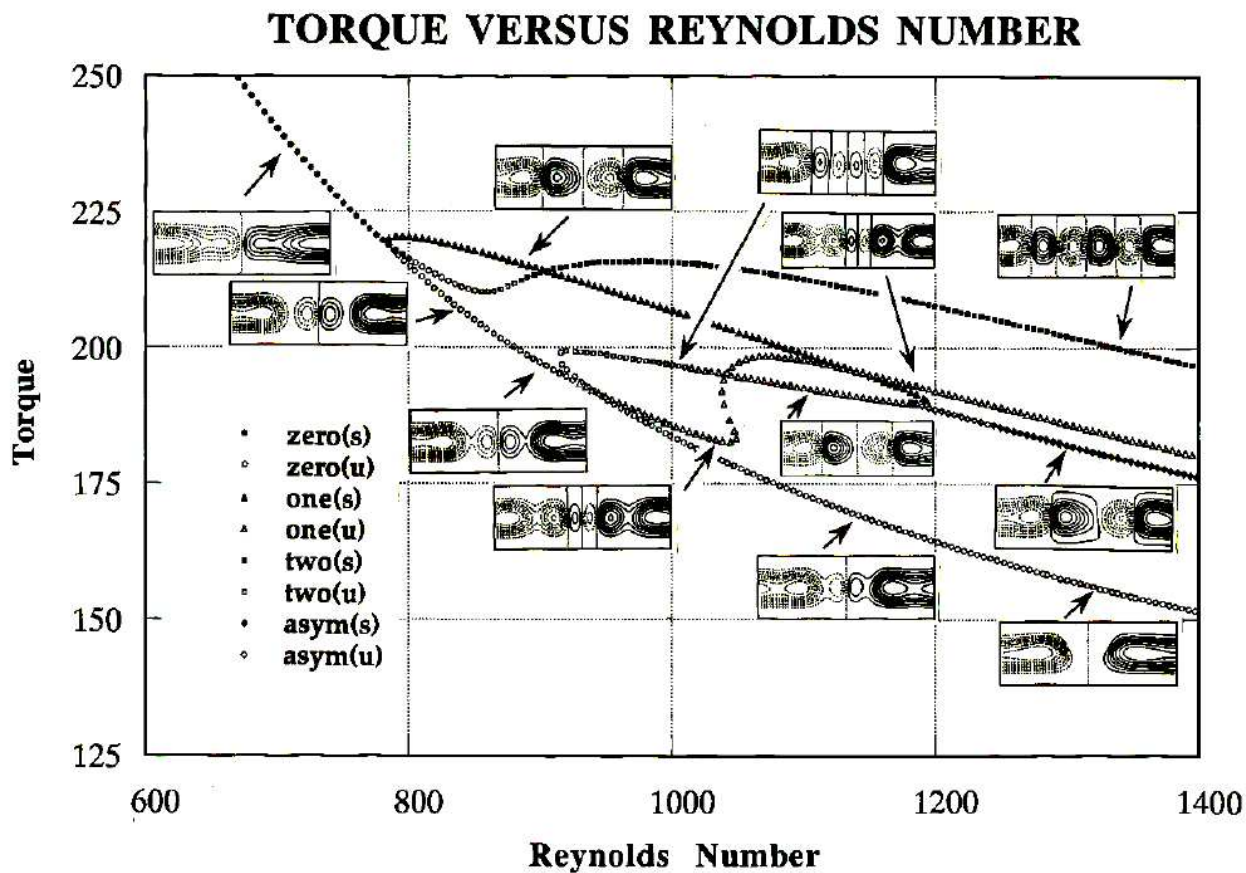
Set $u_i = U_i - \bar{p}$
Calculate $L^{-1}(RN_U + L)u$
Add $L^{-1}N(U)r$ } Use only vectors and operators of length M

TRAVELING WAVES: $U(x - Ct, y, z)$

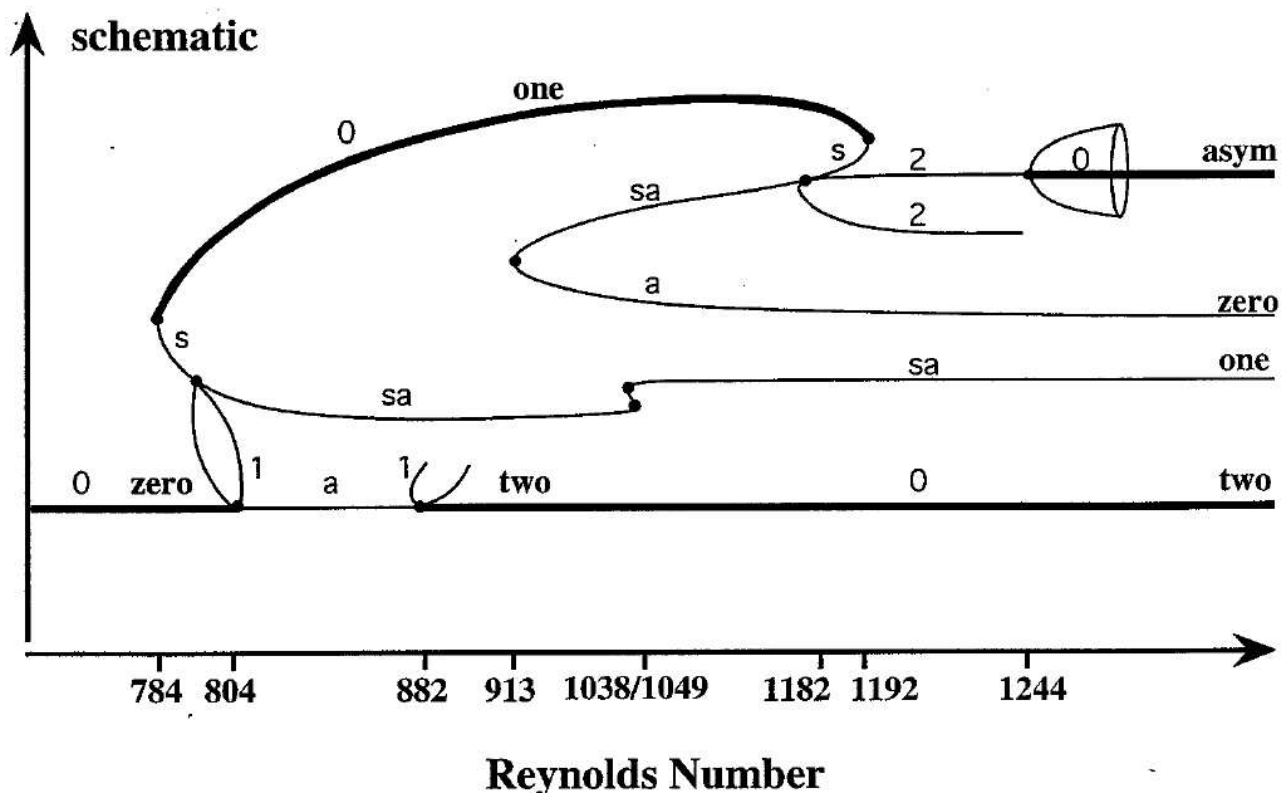
$$\text{Goal } \begin{cases} 0 = C\partial_x U + N(U) + LU \\ 0 = p(U) - \bar{p} \end{cases}$$

$$\left[\begin{array}{c|c} C\partial_x + N_U + L & \partial_x U \\ \hline 0 \ 0 \ \dots \ 0 \ 1 \ 0 \ \dots \ 0 & 0 \end{array} \right] \begin{bmatrix} u \\ c \end{bmatrix} = \begin{bmatrix} C\partial_x U + N(U) + LU \\ U_i - \bar{p} \end{bmatrix}$$

AXISYMMETRIC SPHERICAL COUETTE FLOW WITH $\sigma = 0.18$



AXISYMMETRIC SPHERICAL COUETTE FLOW WITH $\sigma = 0.18$



LINEAR STABILITY ANALYSIS

$$\lambda u = Lu + N_U u$$

How to calculate eigenpairs (λ, u) ?

1) Direct: Diagonalisation = QR decomposition

Storage: M^2

Time: M^3

For 3D case with $M_x = M_y = M_z = 10^2$, we have $M = 10^6$

$$M^2 = 10^{12}$$

$$M^3 = 10^{18}$$

2) Iterative: Calculate a few desired eigenpairs.

Use only matrix-vector products $u \rightarrow Au$

To diagonalise an arbitrary matrix,

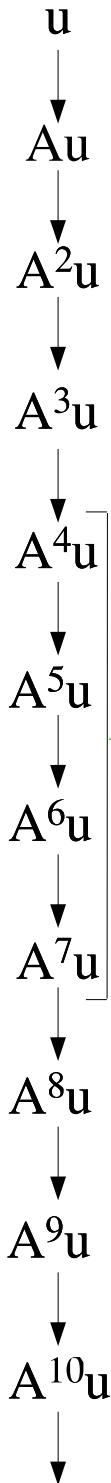
Each product $u \rightarrow Au$ requires M^2 operations
Generating M eigenpairs requires M iterations } M^3

Can gain:

If A is structured or sparse, then $u \rightarrow Au$ takes $\sim M$ ops.

Aim method at desired eigenvalues.

Leading Eigenvalues: $A\Psi = \lambda\Psi$



Power Method

$$v = \frac{A^n u}{\|A^n u\|}$$

$$\lambda \approx \langle v, Av \rangle \quad \Psi \approx v$$

Arnoldi Decomposition

$$H = V^T A V \quad H_{ij} \equiv \langle v_i, A v_j \rangle$$

$K \times K \quad K \times N \quad N \times N \quad N \times K$

$$\begin{bmatrix} v_1 & v_2 & \dots & v_K \\ \downarrow & \downarrow & & \downarrow \\ \downarrow & \downarrow & & \downarrow \end{bmatrix}$$

$$A \approx V H V^T \quad A V \approx V H$$

$N \times N \quad N \times K \quad K \times K \quad K \times N$

$$H \phi = \lambda \phi \quad A V \phi \approx V H \phi = \lambda V \phi$$

orthonormalize
 $\text{span}\{v_1, v_2, \dots, v_K\} =$
 Krylov subspace

H (Hessenberg matrix)

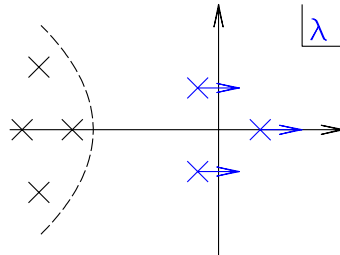
$$K \approx 4$$

$$N \approx 10^6$$

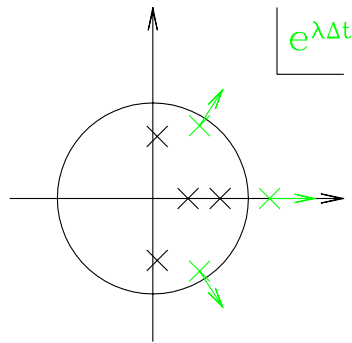
MATRIX TRANSFORMATIONS

If $A u = \lambda u$
then $f(A) u = f(\lambda) u$

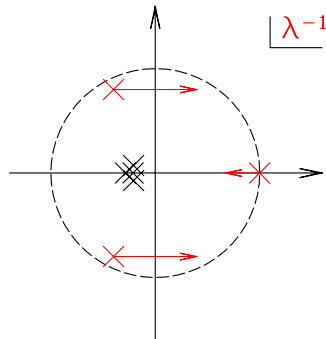
$$f(A) = A$$



$$f(A) = e^{A\Delta t}$$



$$f(A) = A^{-1}$$



$$f(A) = \sum_j f_j A^j$$

f_j chosen dynamically to extract desired eigenvalues ; principle of ARPACK
(Sorensen et al.)

EXPONENTIAL POWER METHOD

$$\begin{aligned} \mathbf{u}_{n+1} &= (\mathbf{I} - \Delta t \mathbf{L})^{-1} (\mathbf{I} + \Delta t \mathbf{N}_U) \mathbf{u}_n \\ &\approx e^{\Delta t (\mathbf{L} + \mathbf{N}_U)} \mathbf{u}_n \end{aligned}$$

Approximation valid for $\Delta t \ll 1$

Time-stepping linearized evolution equation

Enhancement factor at each iteration is

$$\left| \frac{e^{\Delta t \lambda_1}}{e^{\Delta t \lambda_2}} \right| \gtrsim 1 \quad \text{where } \lambda_1 > \lambda_2 > \dots$$

INVERSE POWER METHOD

$$\mathbf{u}_{n+1} = (\mathbf{L} + \mathbf{N}_U)^{-1} \mathbf{u}_n$$

Stokes preconditioning: $(\mathbf{L} + \mathbf{N}_U) \mathbf{u}_{n+1} = \mathbf{u}_n$

$$\begin{aligned} (\mathbf{I} - \Delta t \mathbf{L})^{-1} \Delta t (\mathbf{L} + \mathbf{N}_U) \mathbf{u}_{n+1} &= (\mathbf{I} - \Delta t \mathbf{L})^{-1} \Delta t \mathbf{u}_n \\ (\mathbf{I} - \Delta t \mathbf{L})^{-1} [\mathbf{I} + \Delta t \mathbf{N}_U - (\mathbf{I} - \Delta t \mathbf{L})] \mathbf{u}_{n+1} &= (\mathbf{I} - \Delta t \mathbf{L})^{-1} \Delta t \mathbf{u}_n \\ \underbrace{[(\mathbf{I} - \Delta t \mathbf{L})^{-1} (\mathbf{I} + \Delta t \mathbf{N}_U) - \mathbf{I}] \mathbf{u}_{n+1}}_{\text{difference between two widely spaced consecutive linearized timesteps}} &= \underbrace{(\mathbf{I} - \Delta t \mathbf{L})^{-1} \Delta t \mathbf{u}_n}_{\text{one Stokes timestep}} \end{aligned}$$

Solve with Conjugate Gradient (Bi-CGSTAB) method.

Enhancement factor at each iteration is

$$\left| \frac{\lambda_2}{\lambda_1} \right| \gg 1 \quad \text{for } \lambda_1 \approx 0$$

Can shift to find eigenvalues closest to s

$$\left| \frac{\lambda_2 - s}{\lambda_1 - s} \right| \gg 1 \quad \text{for } \lambda_1 \approx s$$

BOSE-EINSTEIN CONDENSATION

Ultra-cold coherent state of matter

Predicted by Bose (1924) and Einstein (1925)

Realized experimentally by Cornell, Ketterle, Wieman (1995)

Nobel prize (2001)

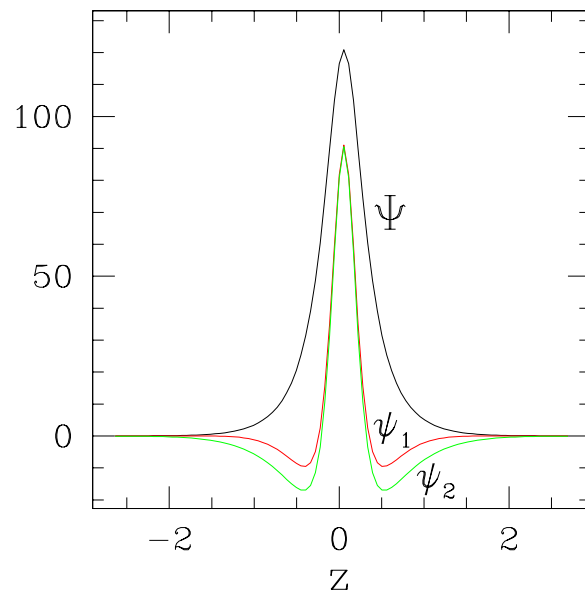
Gross-Pitaevskii / Nonlinear Schrödinger Equation

$$\partial_t \Psi = i \left[\underbrace{\frac{1}{2} \nabla^2}_L + \underbrace{\mu - V(r) - a|\Psi|^2}_N \right] \Psi$$

$$V(\mathbf{x}) = \frac{1}{2} |\boldsymbol{\omega} \cdot \mathbf{x}|^2 = \frac{1}{2} (\omega_r r^2 + \omega_z z^2) \quad (\text{cylindrical trap})$$

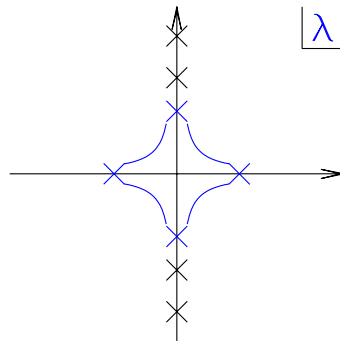
Spatial discretisation up to $M = 10^2 \times 10^2 \times 10^2 = 10^6$

Eigenvalues, energies determine decay rates of condensate.

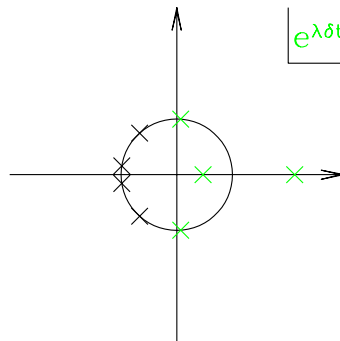


Hamiltonian Systems

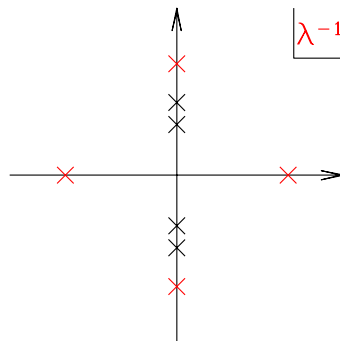
$$f(A) = A$$



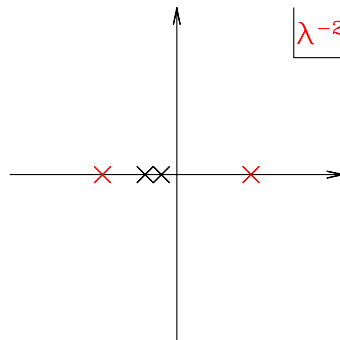
$$f(A) = e^{A\Delta t}$$



$$f(A) = A^{-1}$$



$$f(A) = A^{-2}$$



STEADY STATE SOLVING:

$$0 = L\Psi + N(\Psi)$$

Newton's method + Stokes preconditioning + BICGSTAB

LINEAR STABILITY OF STEADY STATE Ψ :

$$\partial_t \psi = i \left[\left(\frac{1}{2} \nabla^2 + \mu - V(r) \right) \psi - a \Psi^2 (2\psi + \psi^*) \right]$$

$$A \begin{pmatrix} \psi^R \\ \psi^I \end{pmatrix} \equiv \begin{bmatrix} 0 & -(L + DN^I) \\ L + DN^R & 0 \end{bmatrix} \begin{pmatrix} \psi^R \\ \psi^I \end{pmatrix}$$

$$DN^R \equiv \mu - V(x) - 3a\Psi^2$$

$$DN^I \equiv \mu - V(x) - a\Psi^2$$

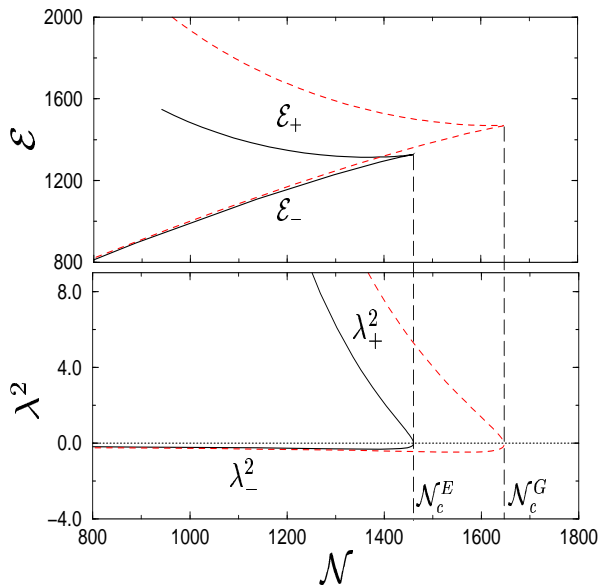
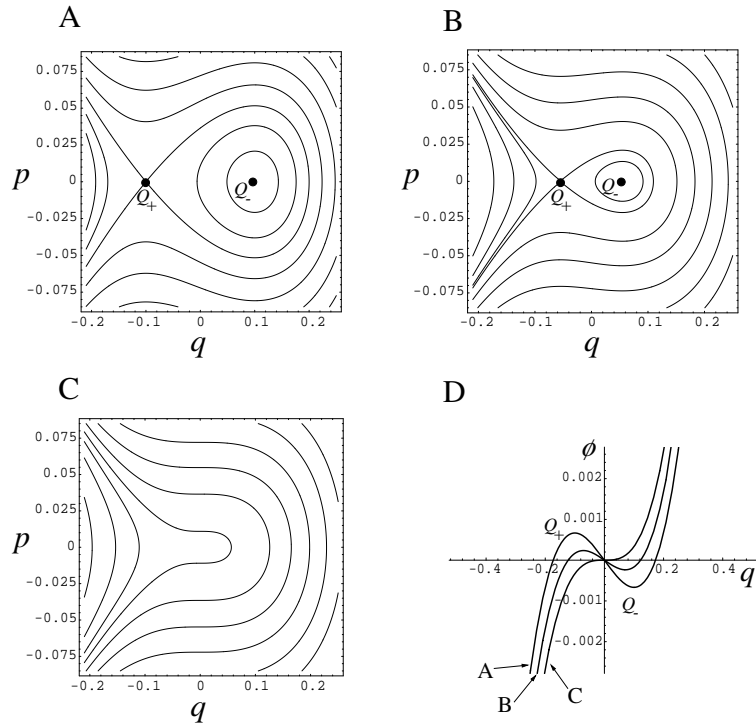
$$A^2 \begin{pmatrix} \psi^R \\ \psi^I \end{pmatrix} = \begin{bmatrix} -(L + DN^I)(L + DN^R) & 0 \\ 0 & -(L + DN^R)(L + DN^I) \end{bmatrix} \begin{pmatrix} \psi^R \\ \psi^I \end{pmatrix}$$

**Inverse square power method
with Stokes preconditioning and shift:**

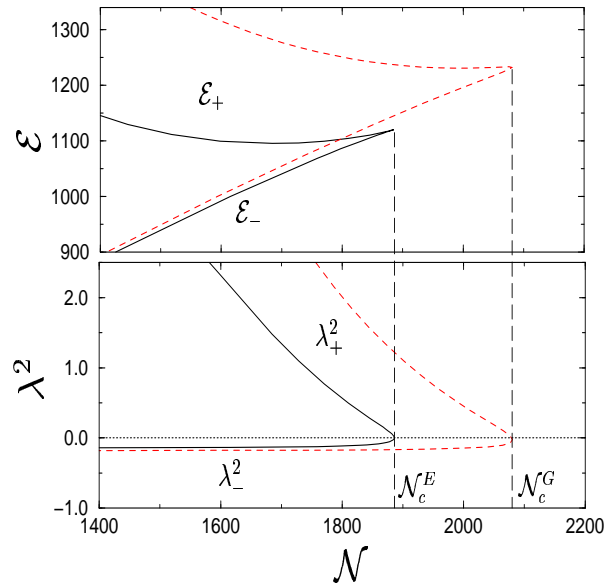
$$\begin{aligned} (A^2 - s^2 I) \psi_{n+1} &= \psi_n \\ L^{-2} (A^2 - s^2 I) \psi_{n+1} &= L^{-2} \psi_n \end{aligned}$$

Solve with BICGSTAB

Hamiltonian saddle-node bifurcation of hyperbolic and elliptic fixed points



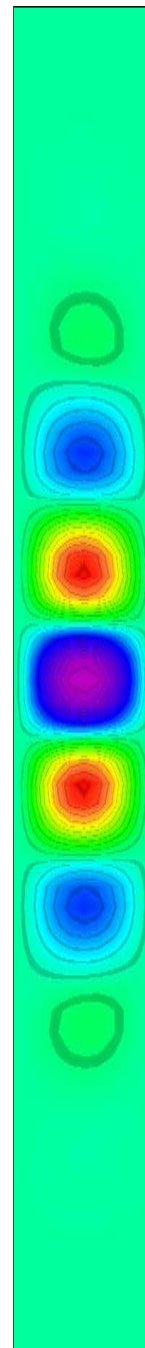
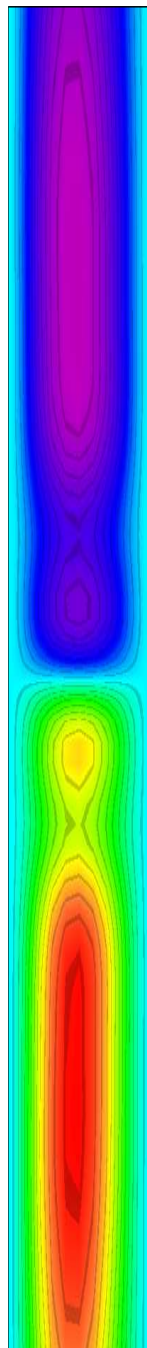
$$\omega_z = \omega_r/5 \text{ (cigar)}$$



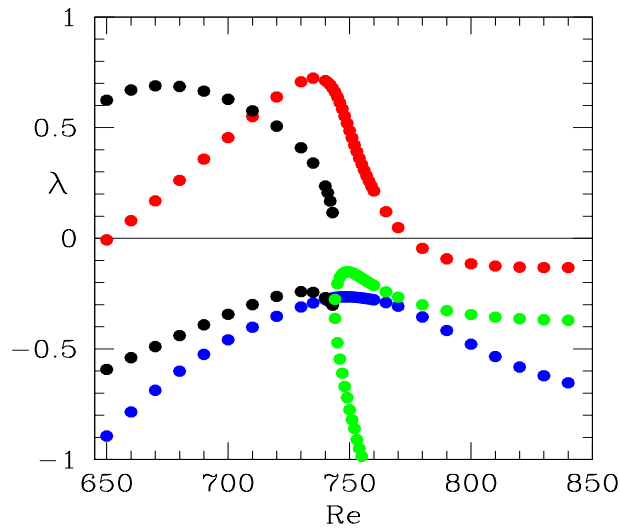
$$\omega_r = \omega_z/5 \text{ (pancake)}$$

AXISYMMETRIC SPHERICAL COUETTE FLOW WITH $\sigma = 0.18$

Basic flow at $Re = 650$ Leading eigenvector

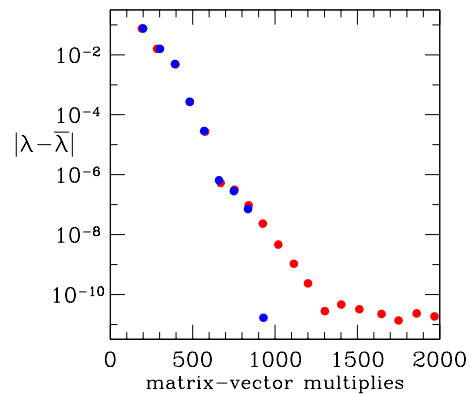
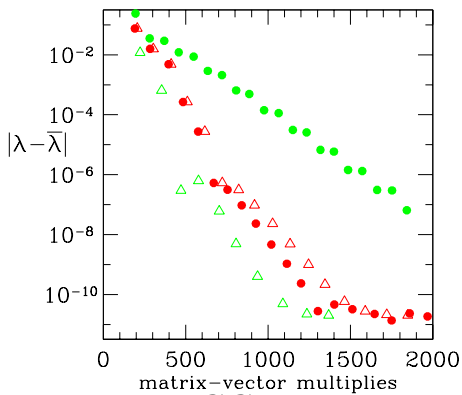
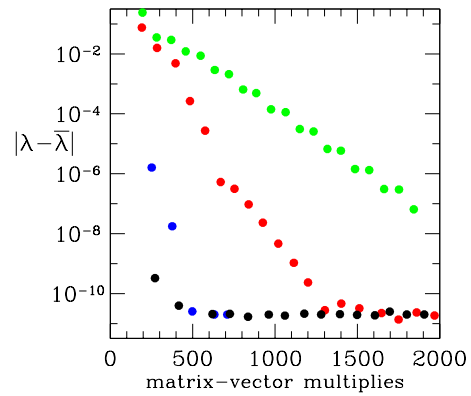
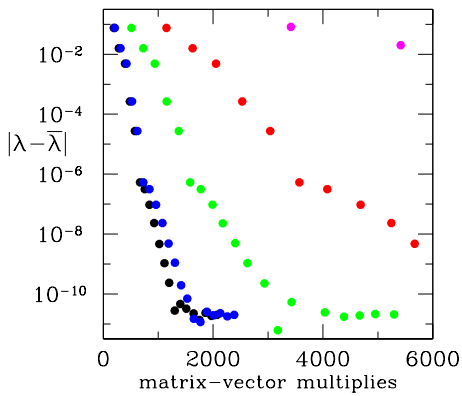


Inverse Power Method on Spherical Couette Flow



$\Delta t = 100, 10, 1, 0.1, 0.01$

$s = -0.152, -0.15, -0.1, 0$



$CG_{crit} = 10^{-7}(\bullet, \bullet), 10^{-9}(\triangle, \triangle)$
 $s = -0.1, 0$

$M = 4096, 16384$

SUMMARY

Time stepping	Steady-state solving	Linear stability analysis
$\partial_t U = (N + L)U$	$0 = (N + L)U$	$\lambda u = (N_U + L)u$
Implicit/explicit Euler	Newton	Inverse power/Arnoldi
$U(t + \Delta t) = BU(t)$ $= (I - \Delta t L)^{-1}$ $(I + \Delta t N)U(t)$	$(N_U + L)u$ $= (N + L)U$ $U \leftarrow U - u$	$(N_U + L)u_{n+1} = u_n$
$\equiv P(I + \Delta t N)U(t)$	$A_U u = AU$ $PA_U u = PAU$	$A_U u_{n+1} = u_n$ $PA_U u_{n+1} = Pu_n$
	3 – 4 Newton steps	3 – 4 Inverse Arnoldi steps
	200 BiCGSTAB iters/step	200 BiCGSTAB iters/step