

Jamming percolation and glass transition in lattice models

C.Toninelli, IHES

G.Biroli, D.S.Fisher

Outline

- **Physical motivations:**
phenomenology of liquid/glass transition;
what is an ideal glass transition?

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phenomenology of liquid/glass transition;
what is an ideal glass transition?
- **A new class of kinetically constrained models:**
existence of an ideal glass transition;
ergodicity breaking transition;
identification of ρ_c ;
first order/critical character of transition;
a new percolation transition.

Liquid-Glass transition

- Liquid at $T > T_m$ → fast cooling → avoided crystallization → supercooled liquid: dramatic slowing down without significant structural changes.

Glass "transition" temperature: $\tau(T_g) = 10^3 s$

Liquid-Glass transition

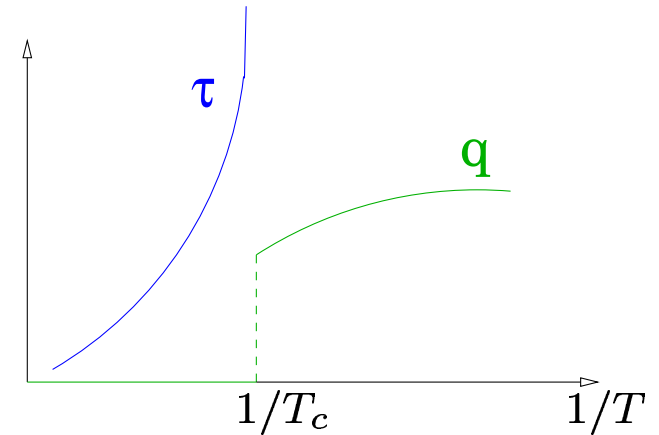
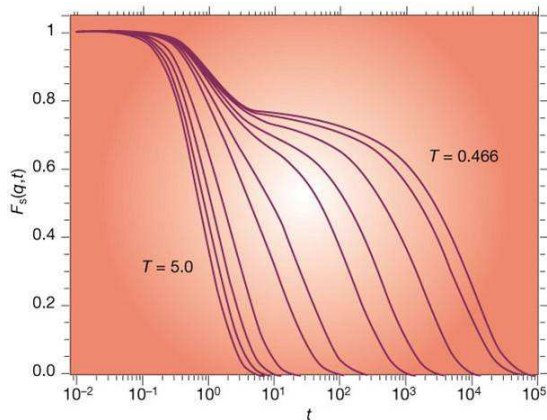
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- ● Is there a **divergence of $\tau(T)$** at finite $T < T_g$?
- Is there an **equilibrium glass transition**?
- Which are the **mechanisms** causing the slowing down of dynamics and its heterogeneous character?
- Which are the typical **length/time scales**?

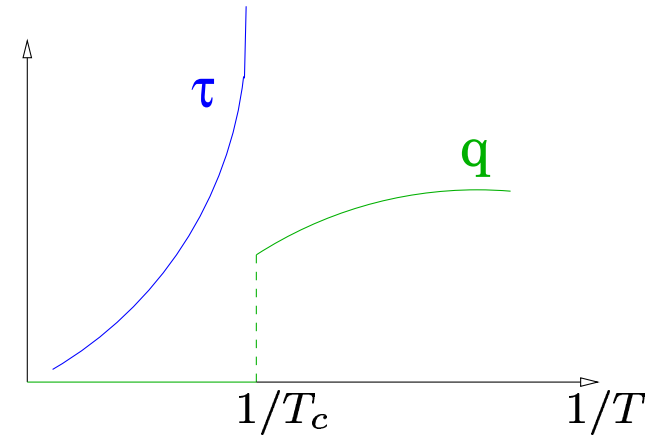
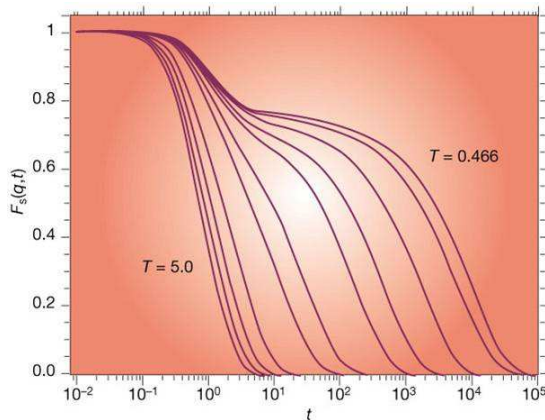
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Mixed first order/critical properties



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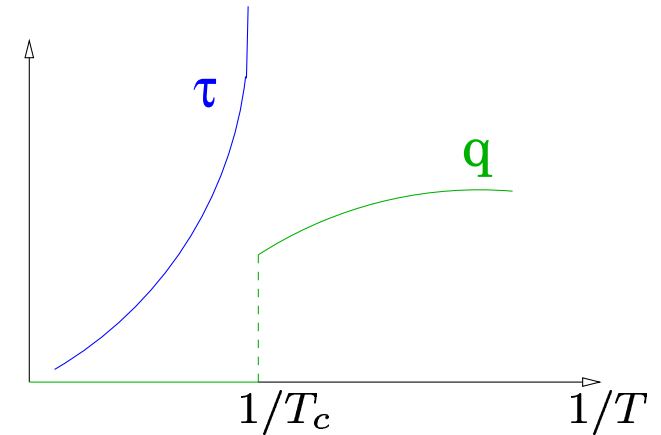
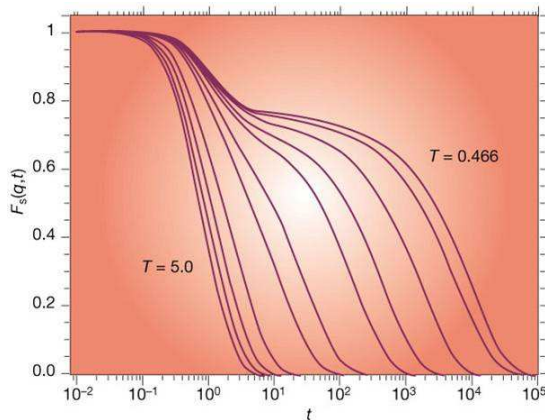


- Discontinuous order parameter:

$$q(T) = \lim_{t \rightarrow \infty} F(k, t) = \lim_{t \rightarrow \infty} F \cdot T \cdot \langle \rho(t, x) \rho(0, x) \rangle_c$$

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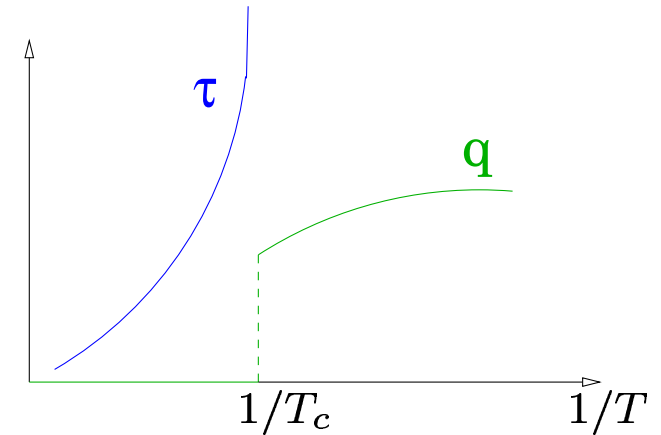
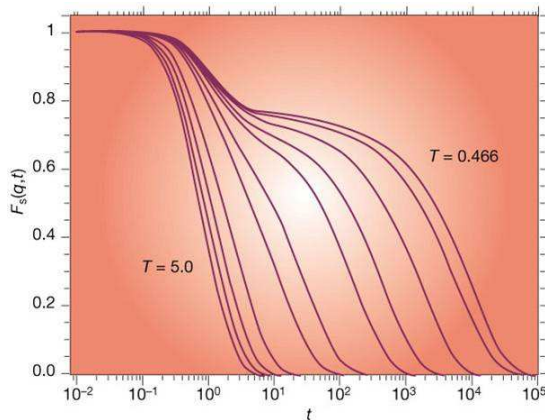
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- **Superharrenius divergence:** $\tau(T) \simeq \exp \frac{c}{T - T_c}$

Kinetically Constrained Models I

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- Systems of particles with **discrete positions** on a lattice and **hard core** constraint;
- **Stochastic conservative or non-conservative dynamics**: sequence of jumps or birth/death of the particles;
- **Additional constraints** (besides hard core) to mimic **cage effect** → blocked configurations, slow dynamics, dynamical arrest, glassy phenomenology,

Kinetically constrained models II

Ex. Kob-Andersen '92:

Constraint: move allowed if jumping particle has ≤ 4 occupied n.n. both before and after moving.

Simulations: slow, glassy dynamics at high density

Conjecture: ergodicity breaking transition at $\rho_c < 1$

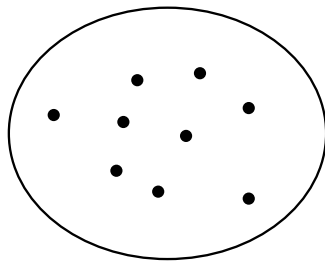
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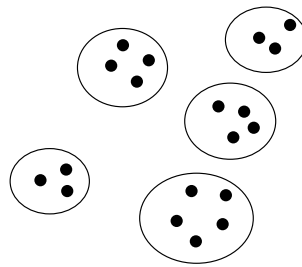
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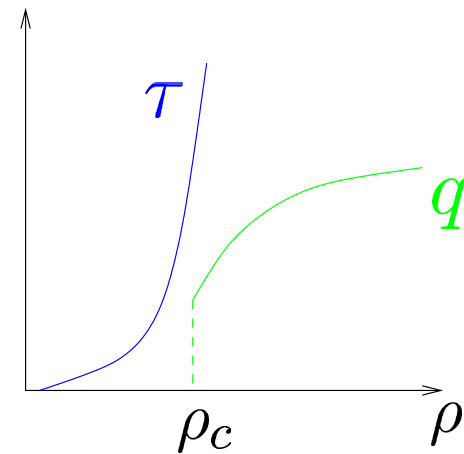


$$\rho < \rho_c$$



$$\rho > \rho_c$$

Configuration space



Kinetically constrained models II

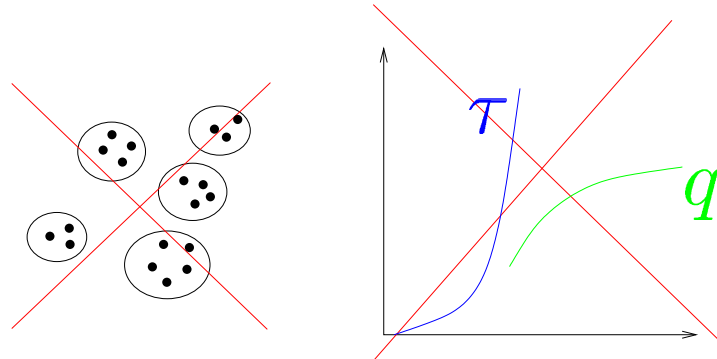
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Rigorous results (G.Biroli, D.S.Fisher, C.T.)



Large finite size effects $L_c(\rho) \propto \exp \exp \frac{c}{1-\rho}$

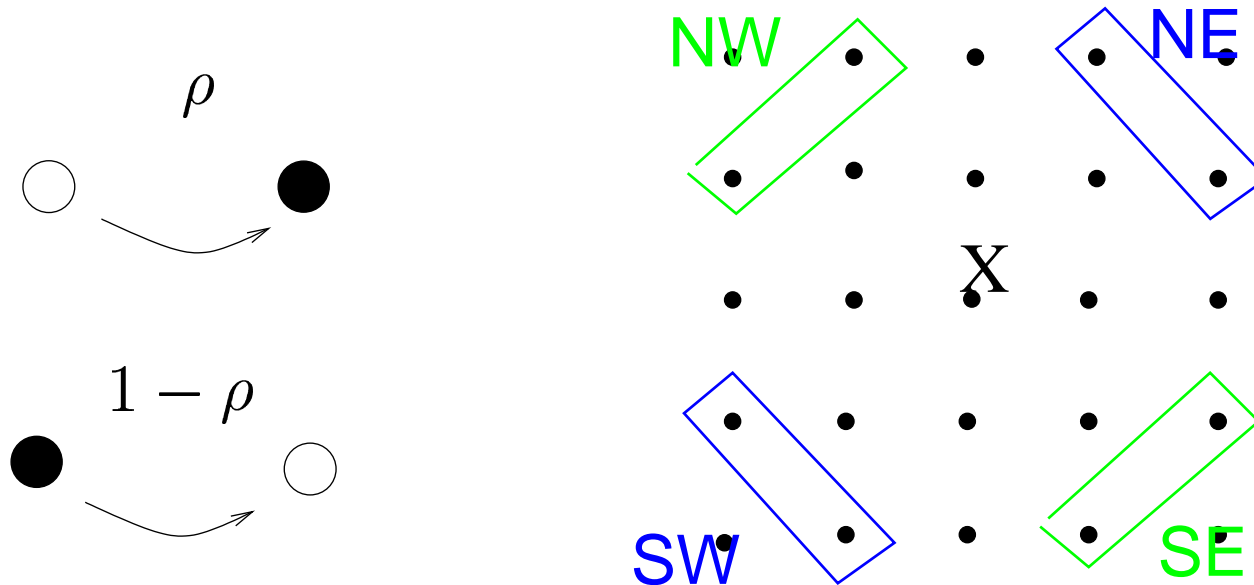
Relaxation: rearrangement of regions with diverging size and distance for $\rho \rightarrow 1$ $\tau \propto L_c(\rho)^2$

A new class of models

Configurations:

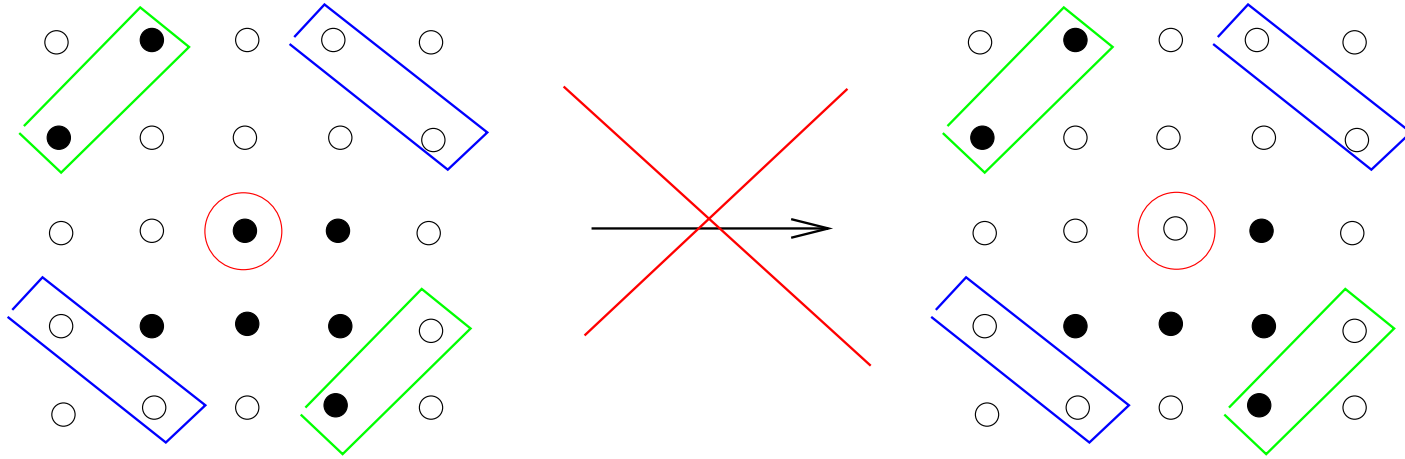
particles on a square lattice with hard core constraint

Dynamic: birth/death

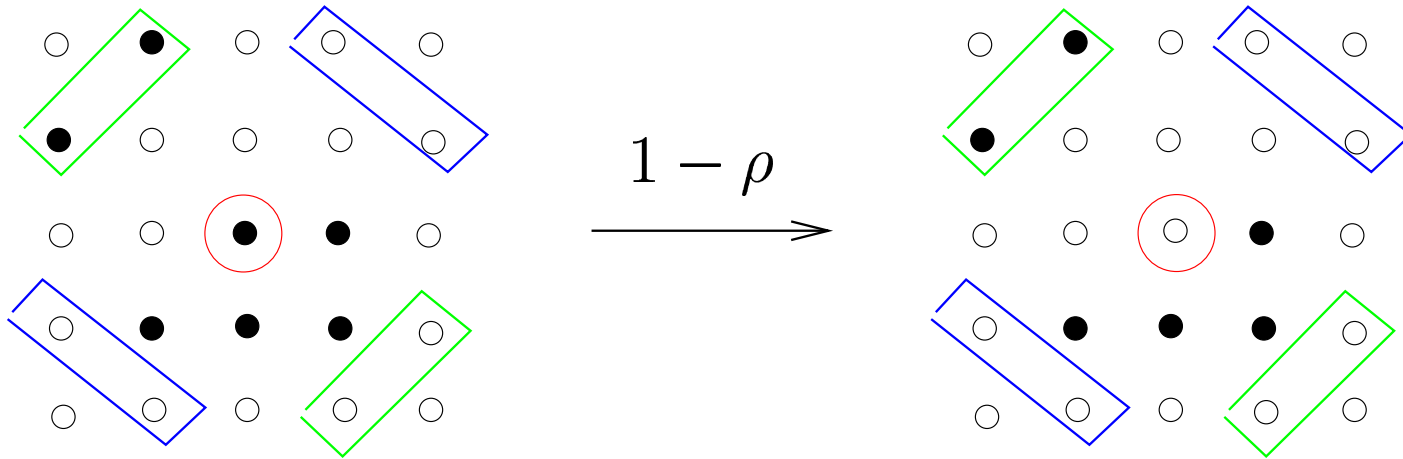
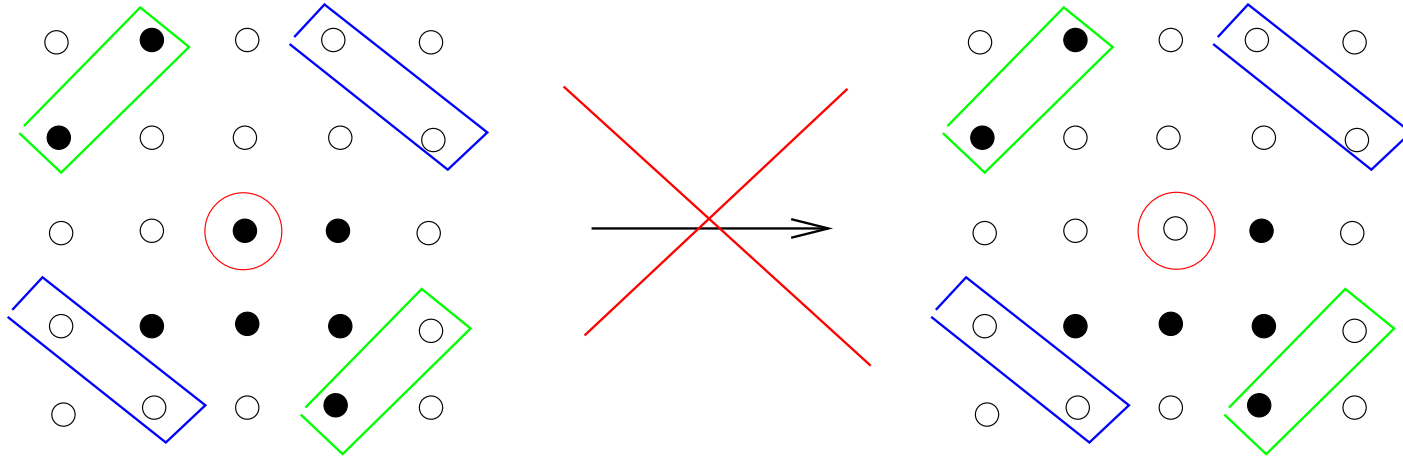


(NE OR SW empty) AND (NW OR SE empty)

Rules



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Results I

- **Trivial statics** (correlation length = 1)
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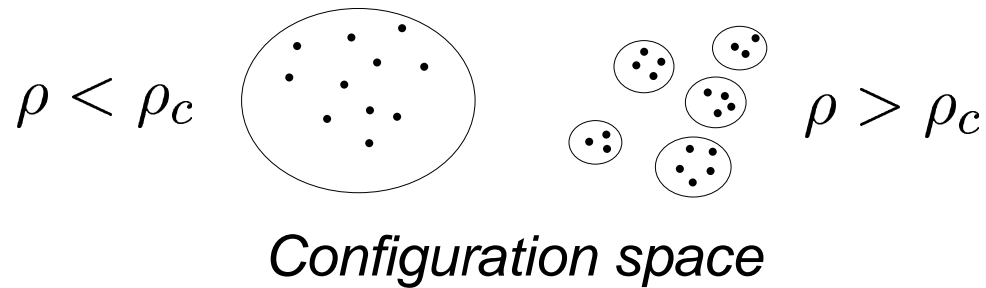
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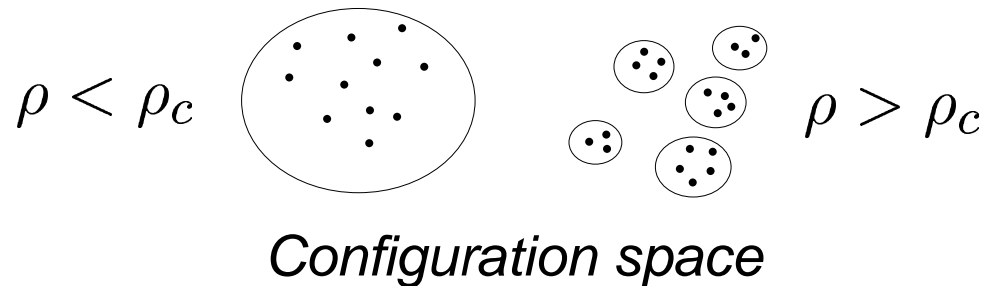
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- **Identification of ρ_c**
= critical density of directed percolation in $d = 2$

Results II

- **Critical properties** $\rho \nearrow \rho_c$

$$L_c \simeq \exp(c \xi_{\parallel}^{1-z}),$$

ξ_{\parallel} = correlation length of DP

$$\xi_{\parallel} \simeq 1/(\rho_c - \rho)^{\nu}, \quad \nu \simeq 1.7, \quad z \simeq 0.63 \quad \text{anisotropy DP}$$

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density of blocked clusters discontinuous at ρ_c
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 $q(\rho_c) > 0 \rightarrow$ plateau for $F(q, t)$

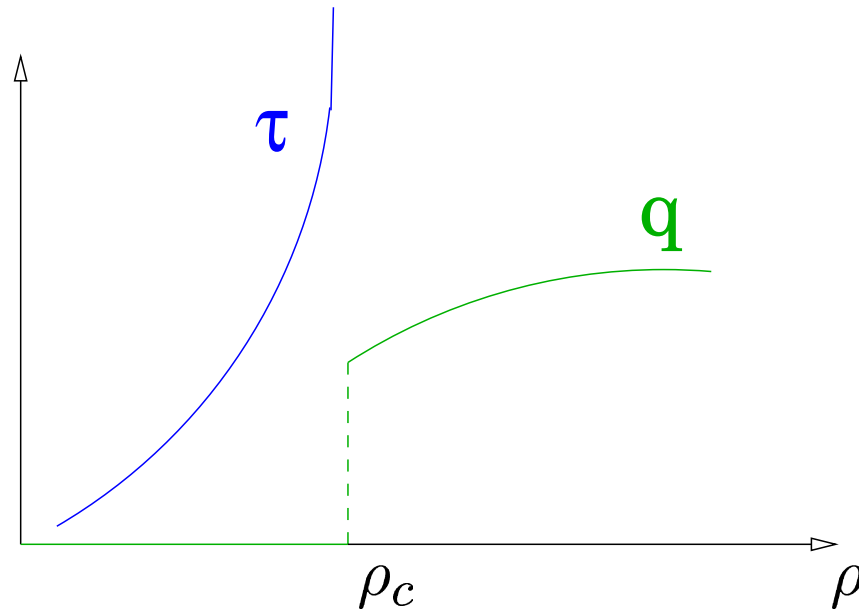
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C.Toninelli, G.Biroli, D.S. Fisher *Phys. Rev. Lett* in press (cond-mat 0509661)

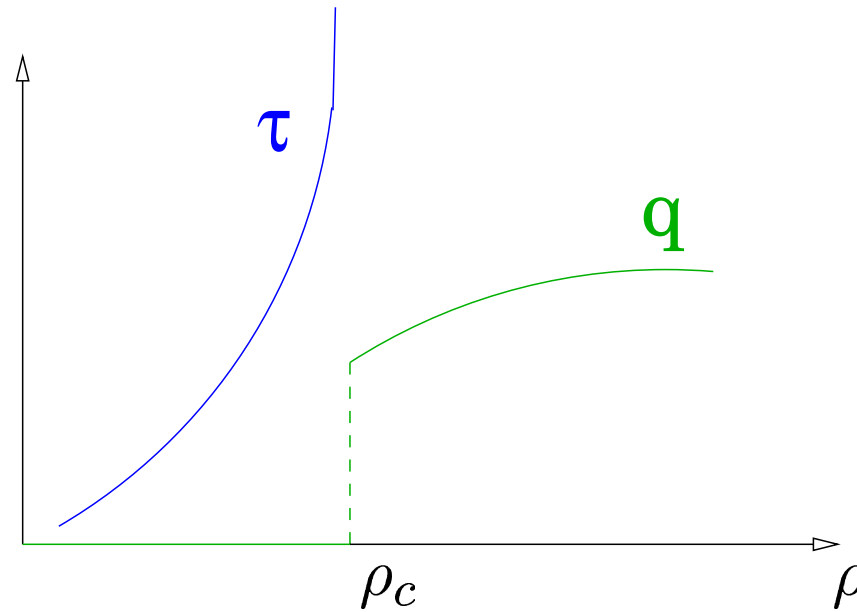
C.Toninelli, G.Biroli (cond-mat 0512335)

Results III



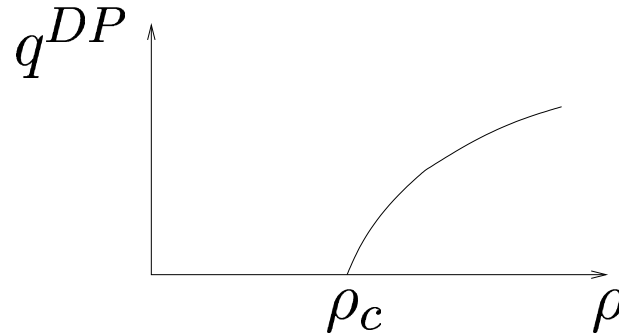
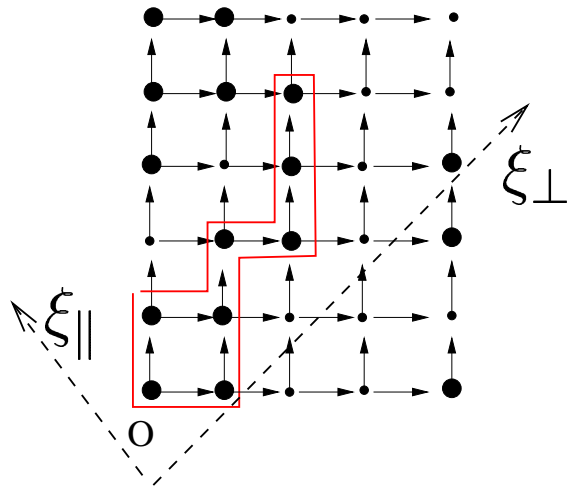
→ Finite dimensional model with ideal glass transition

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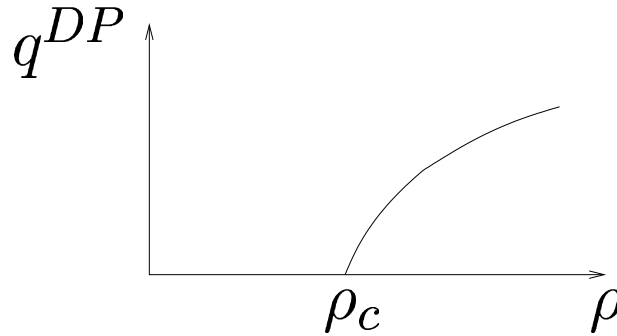
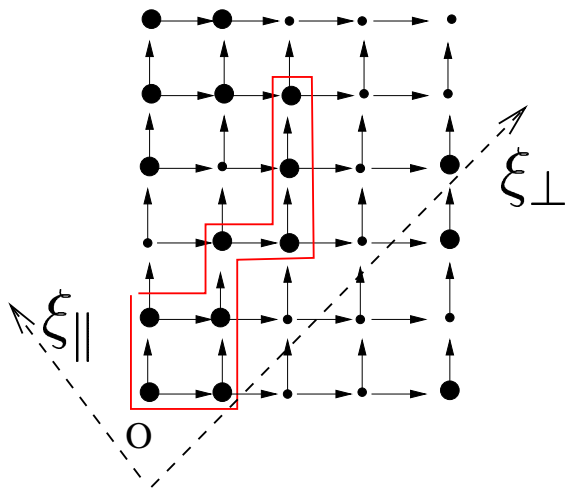
- Finite dimensional model with ideal glass transition
- A new percolation transition:
discontinuous + anomalous critical properties
Jamming percolation

Directed Percolation



q^{DP} = density of infinite clusters following arrows

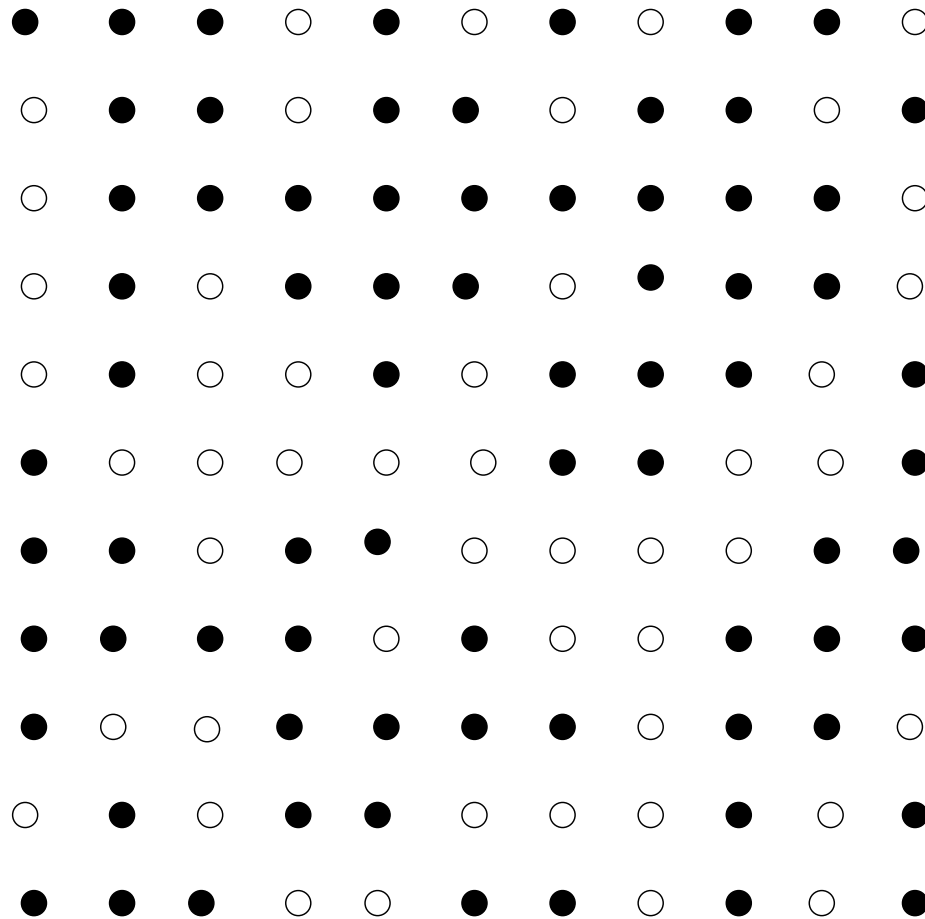
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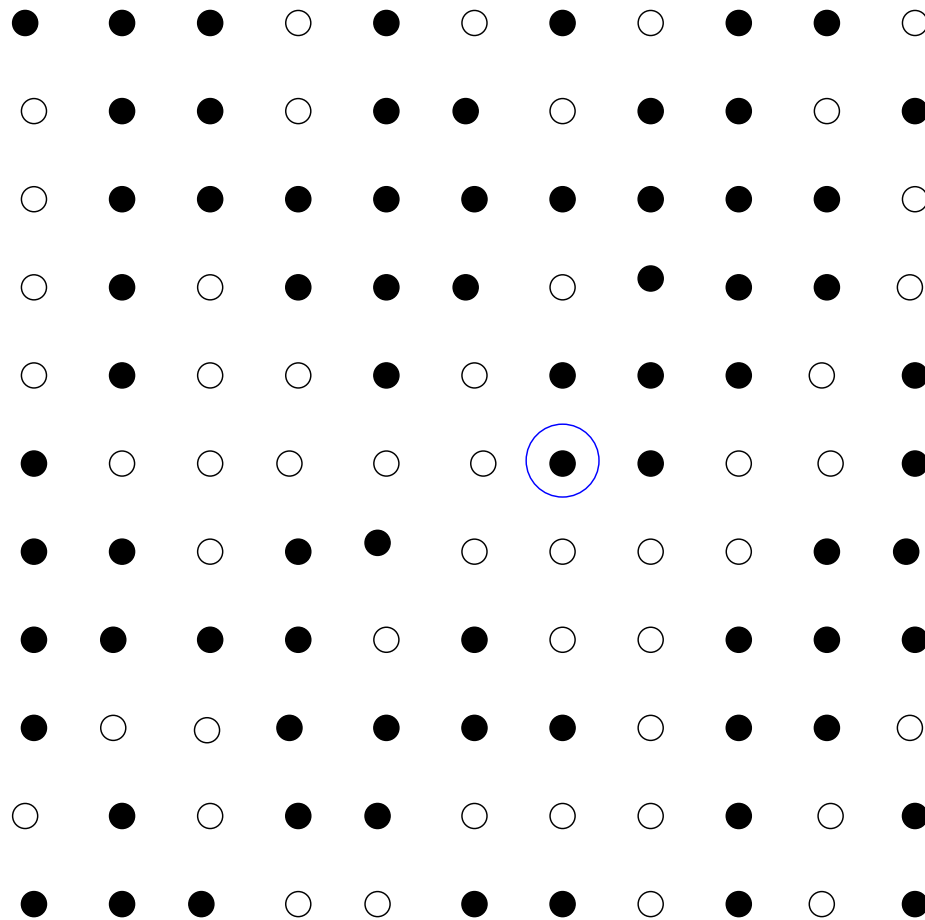
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Continuous transition à $\rho_c \simeq 0.7$

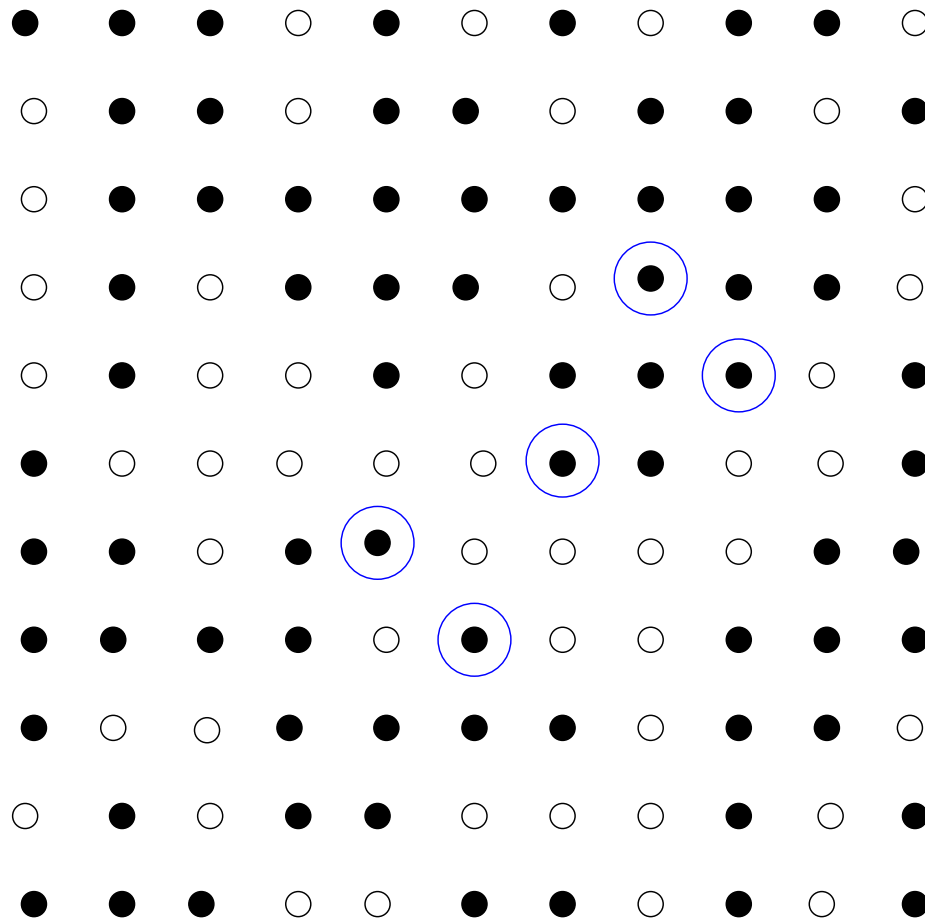
Ergodicity breaking $\rho > \rho_c$



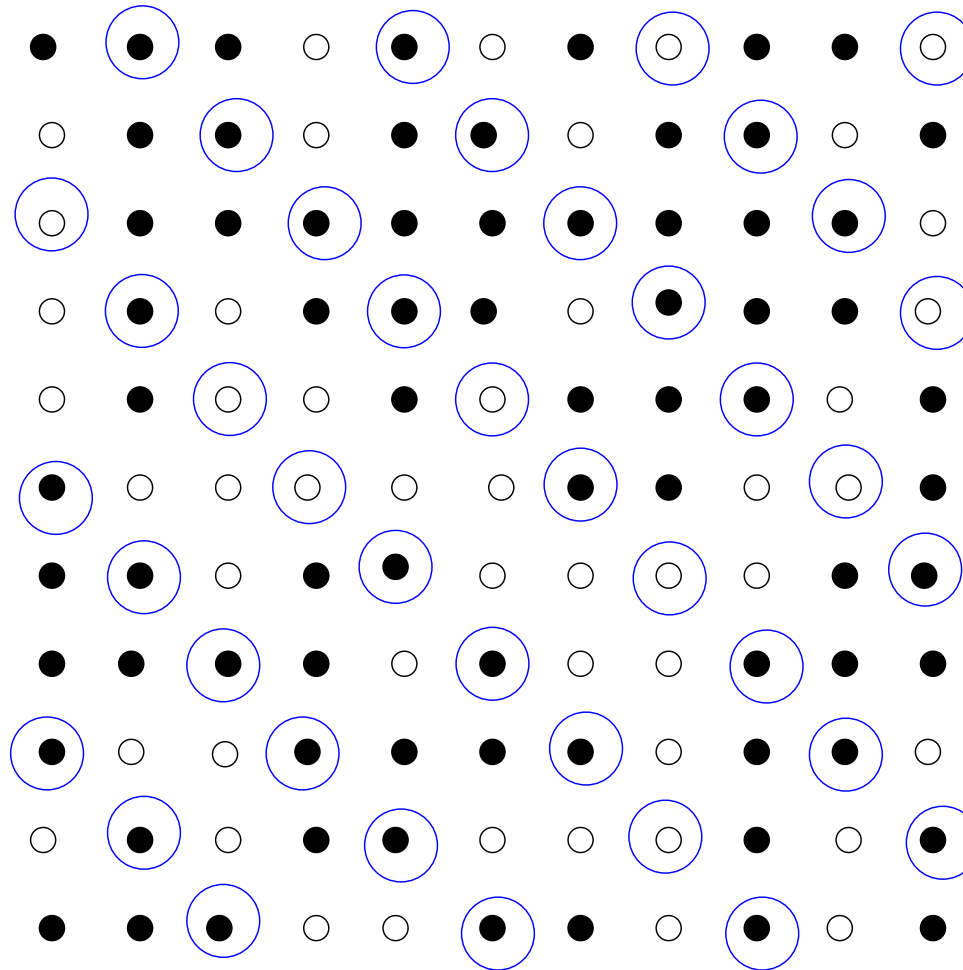
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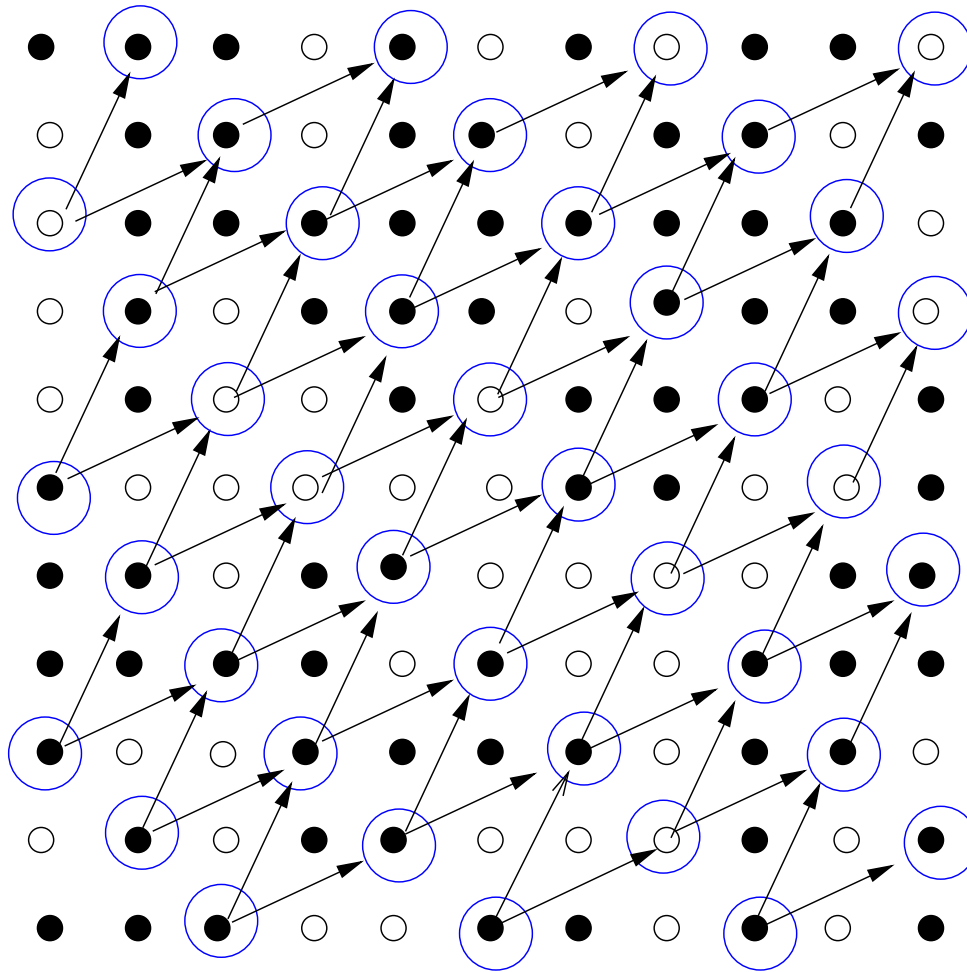
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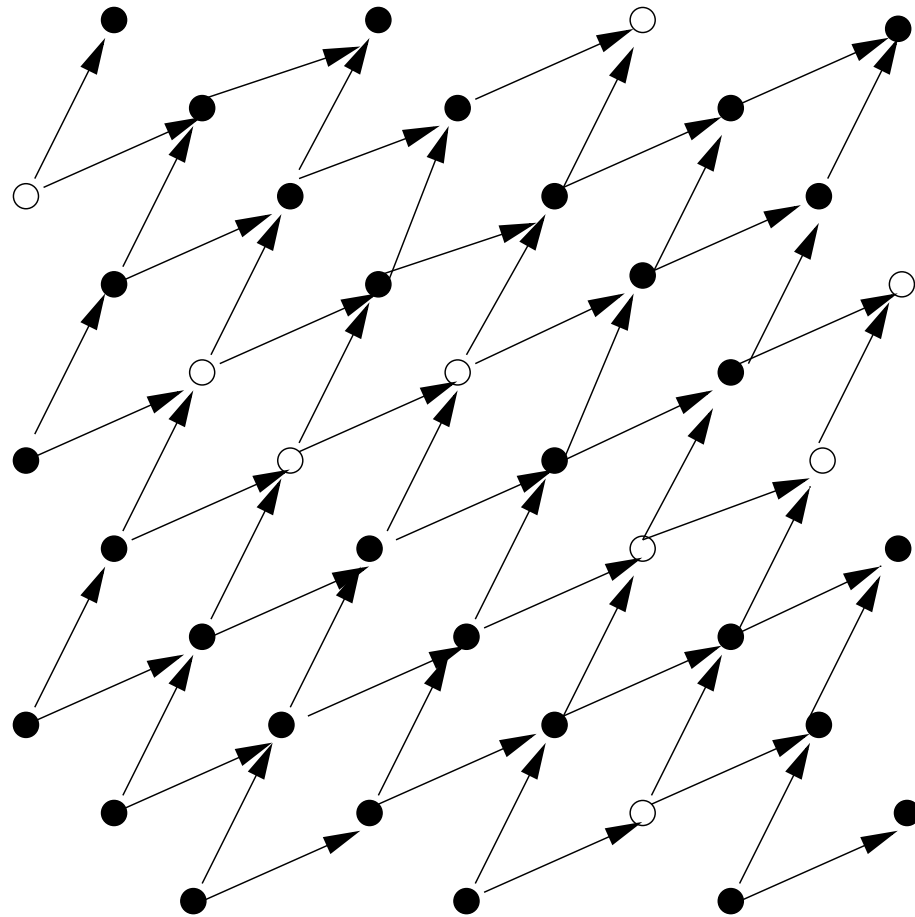
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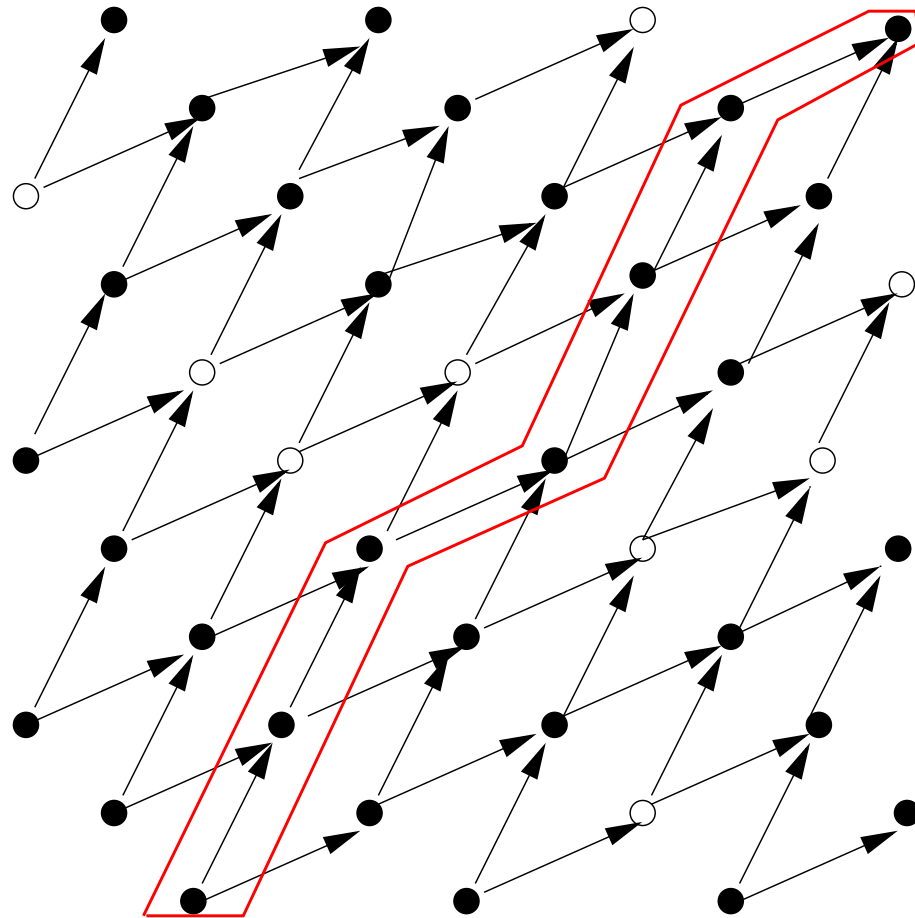
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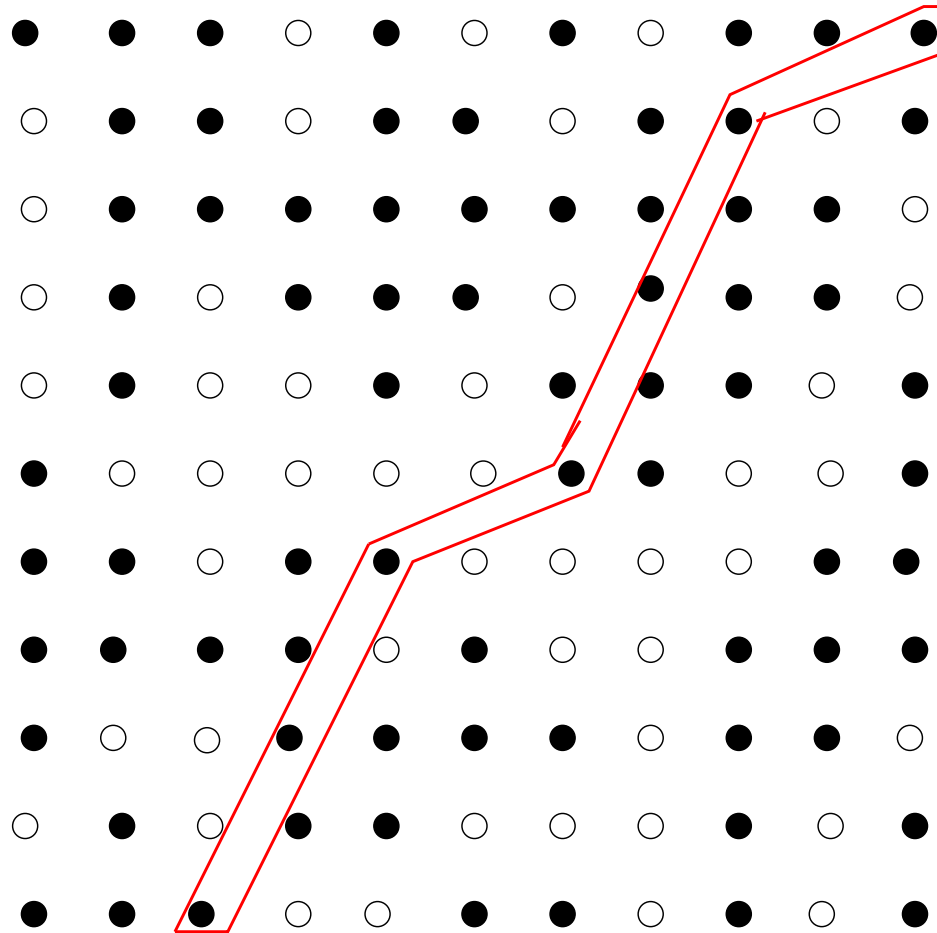
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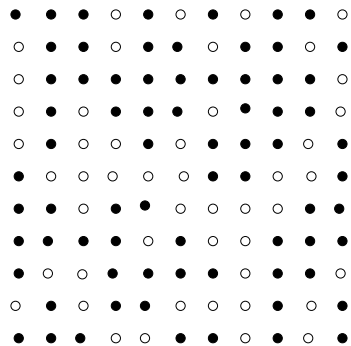
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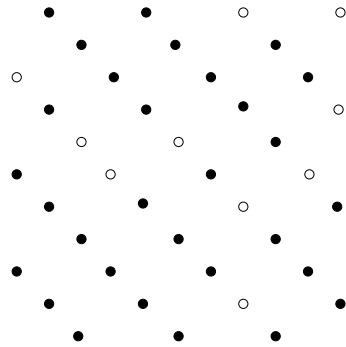
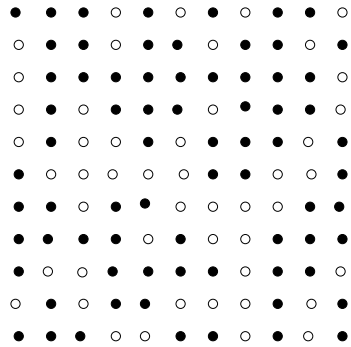
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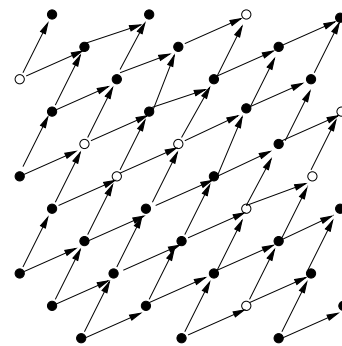
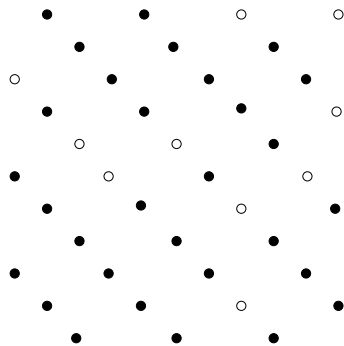
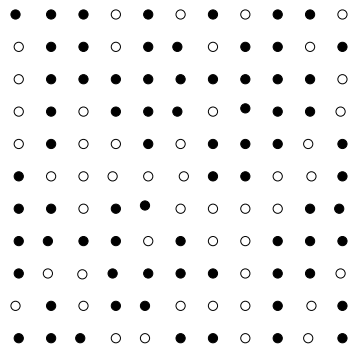
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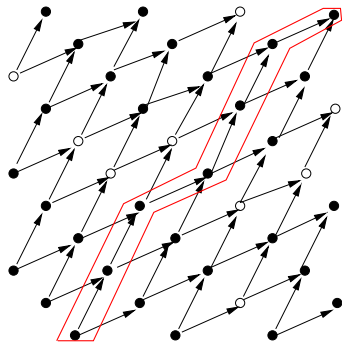
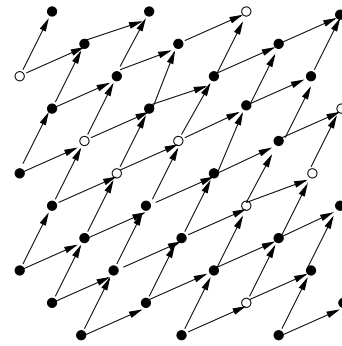
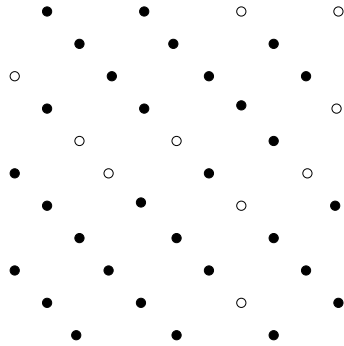
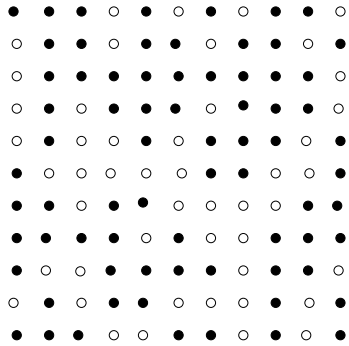
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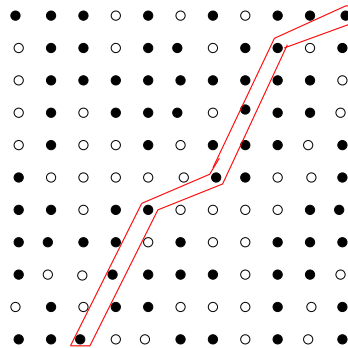
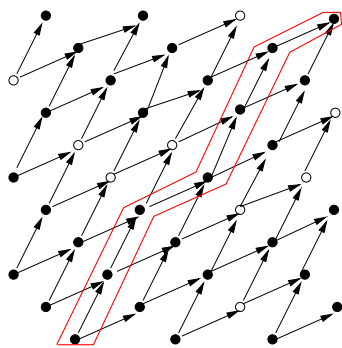
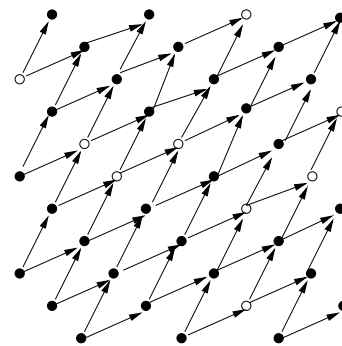
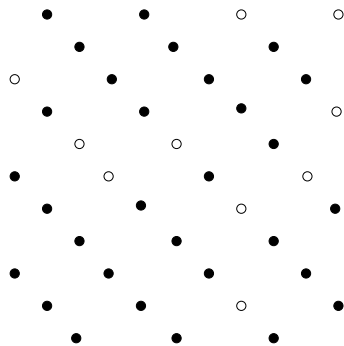
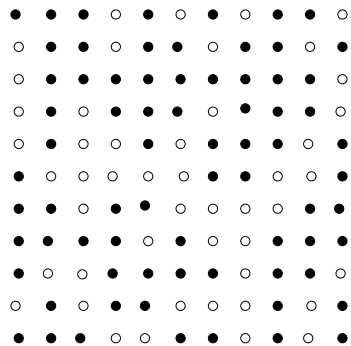
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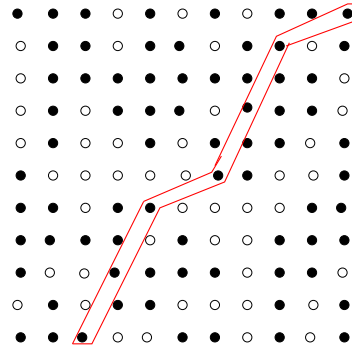
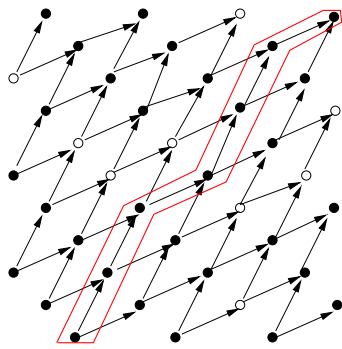
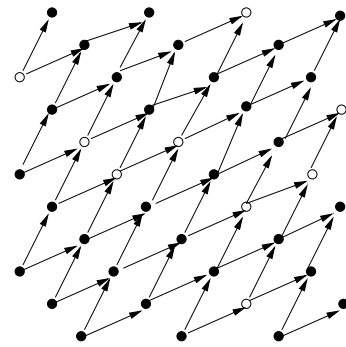
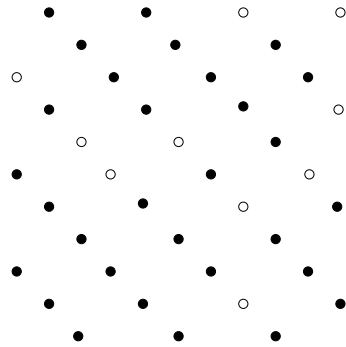
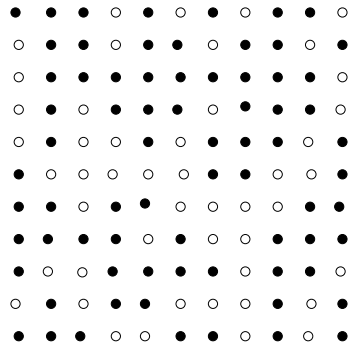
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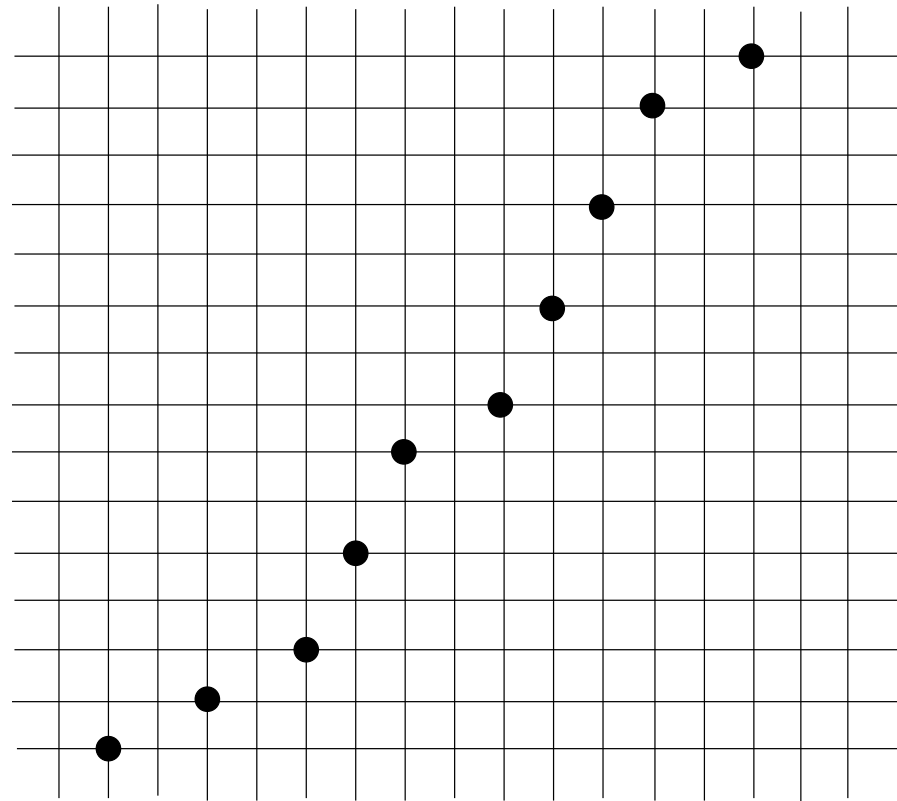
Percolation of blocked clusters \rightarrow ergodicity breaking

T-Junctions

DP spanning clusters \rightarrow blocked
Blocked cluster \nrightarrow DP spanning cluster

T-Junctions

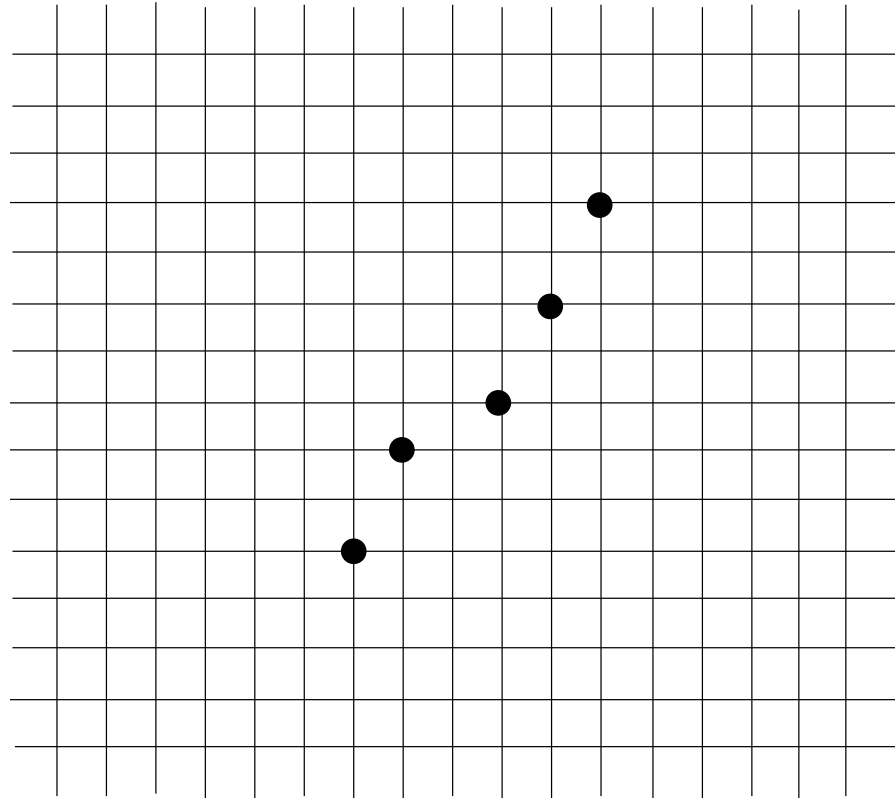
DP spanning clusters \rightarrow blocked
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Percolating NE-SW cluster \rightarrow blocked

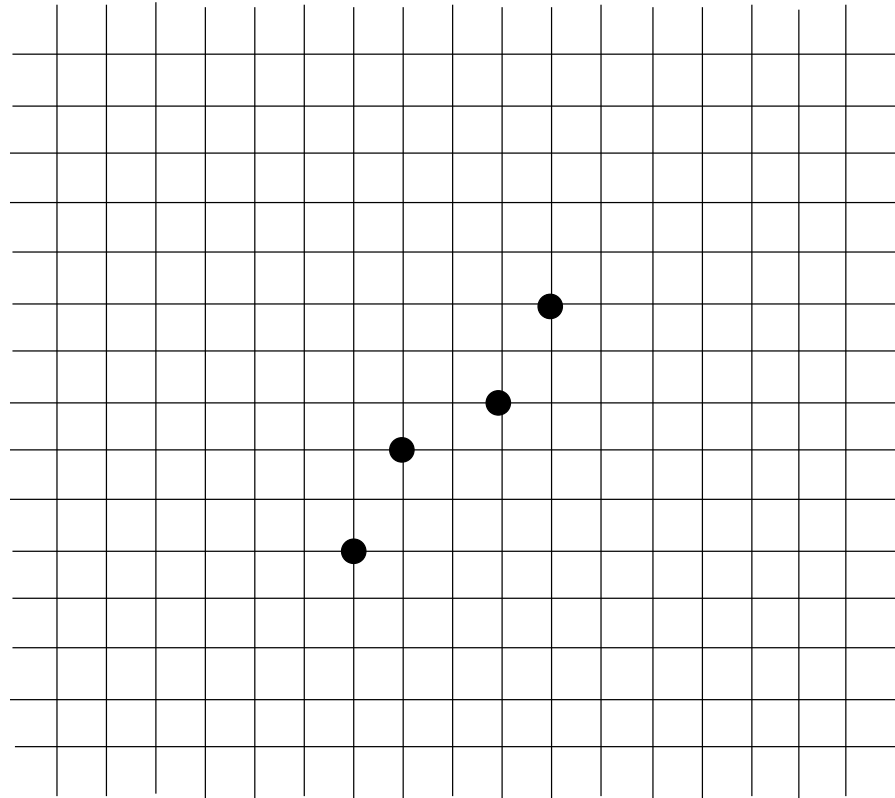
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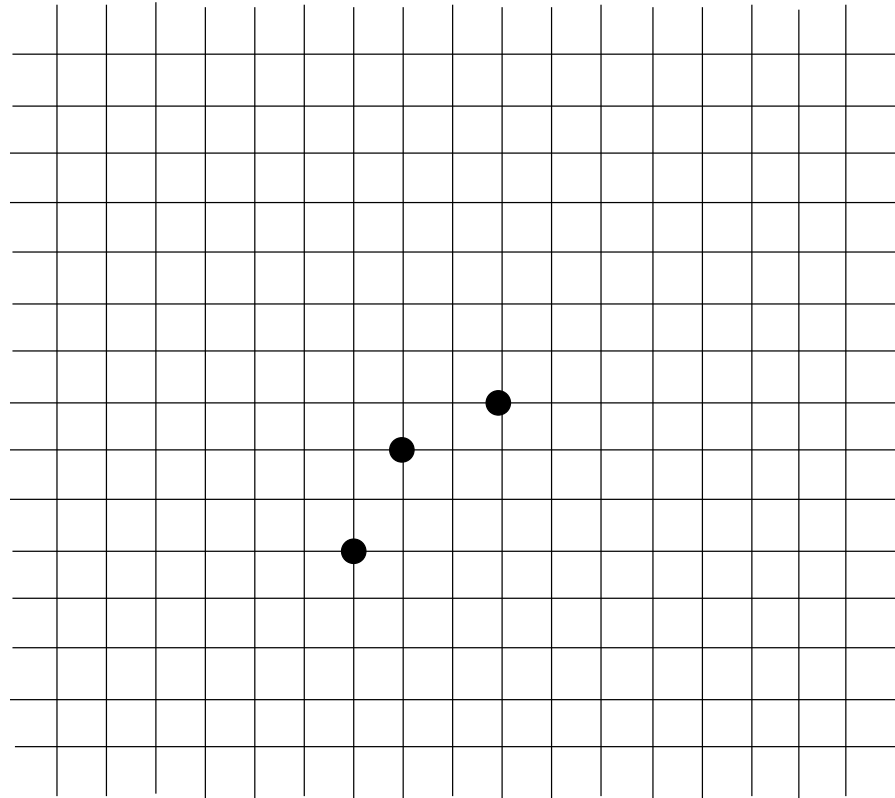
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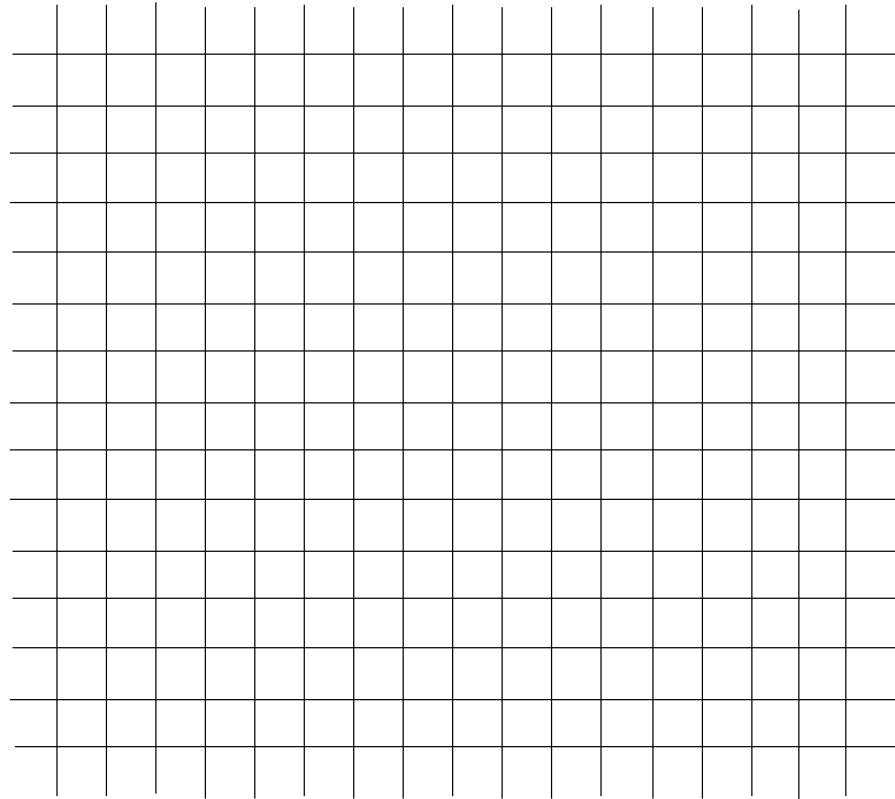
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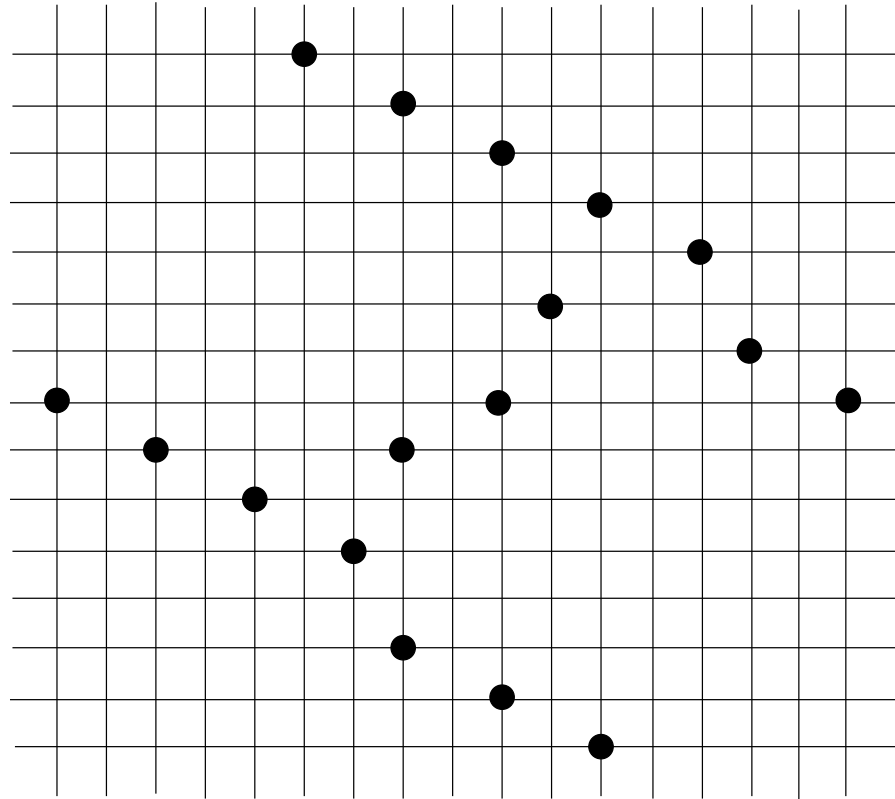
DP spanning clusters \rightarrow blocked
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Non percolating NE-SW cluster \rightarrow non blocked

T-Junctions

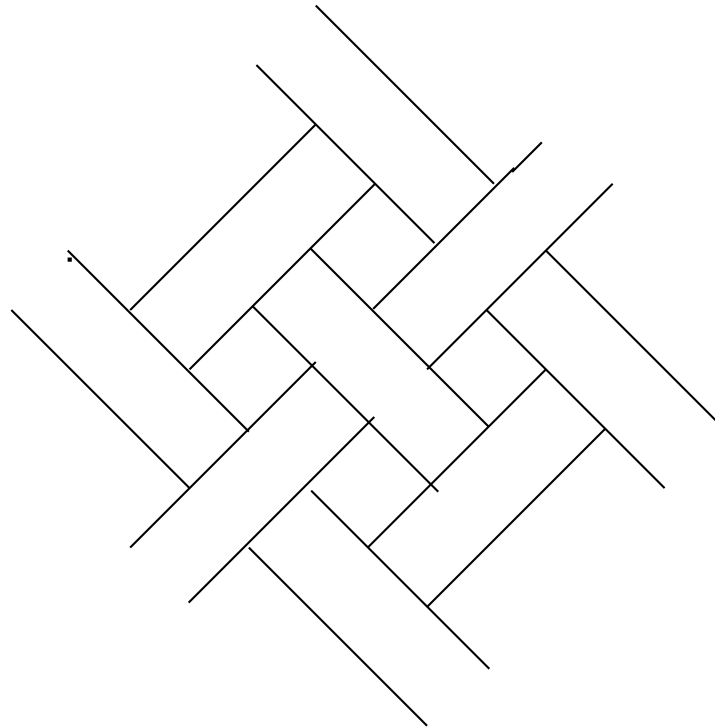
DP spanning clusters \rightarrow blocked
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NE-SW cluster blocked by T-junction with NW-SE
percolating cluster

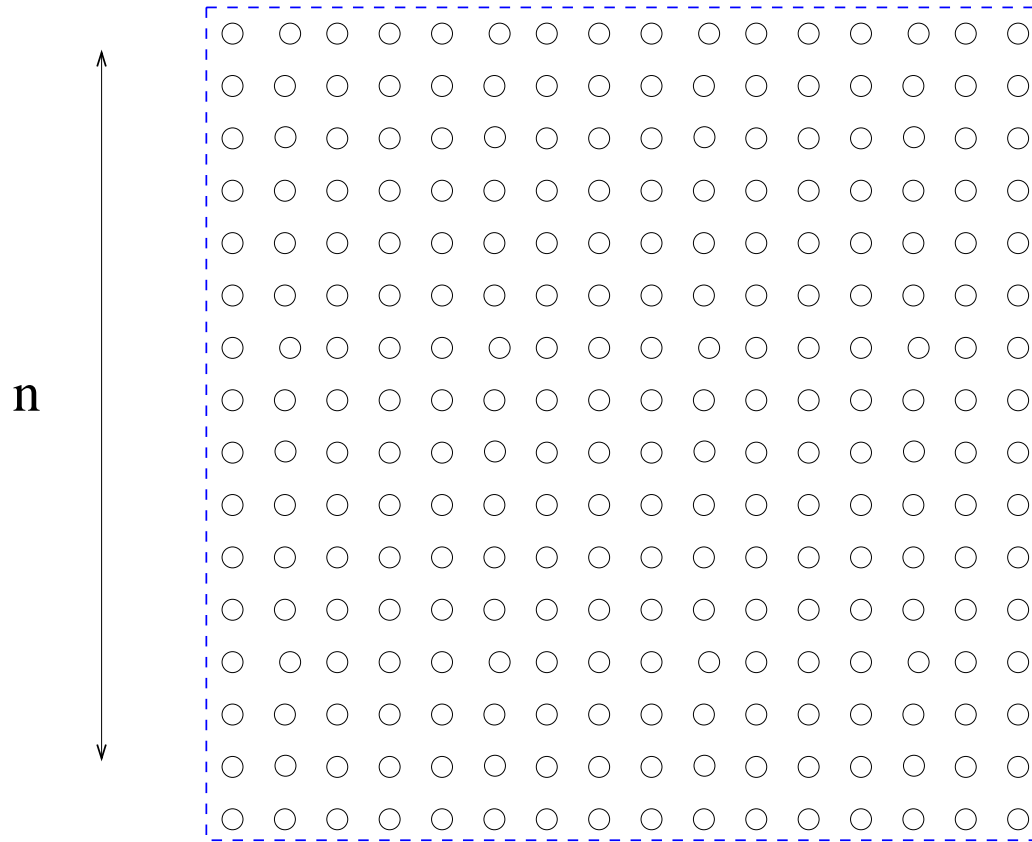
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Lines = NE-SW or NW-SE clusters
Ex. Blocked structures without NW-SE nor NE-SW
percolating clusters

Ergodicity $\rho < \rho_c$



Ergodicity $\rho < \rho_c$



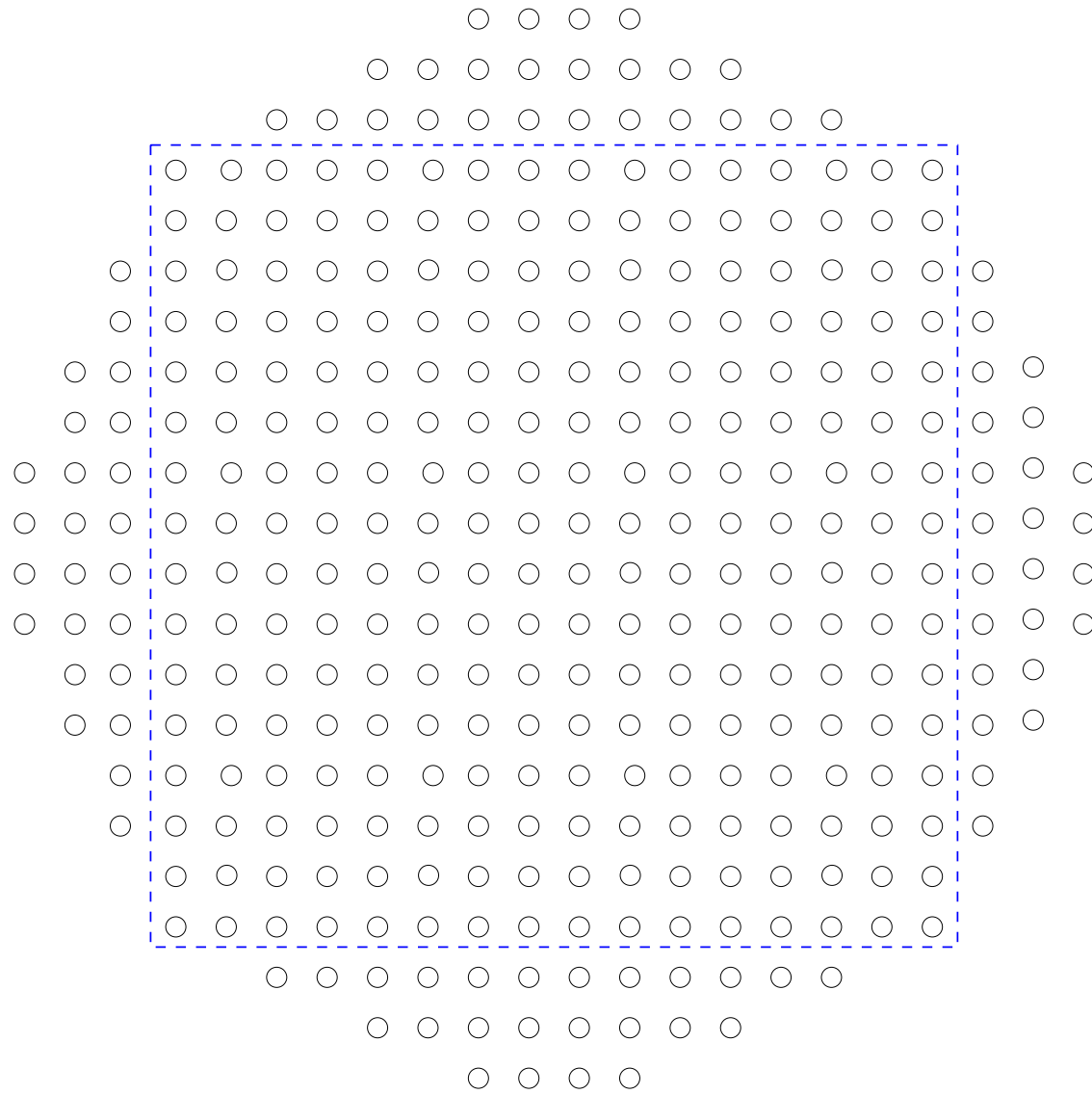
Ergodicity $\rho < \rho_c$



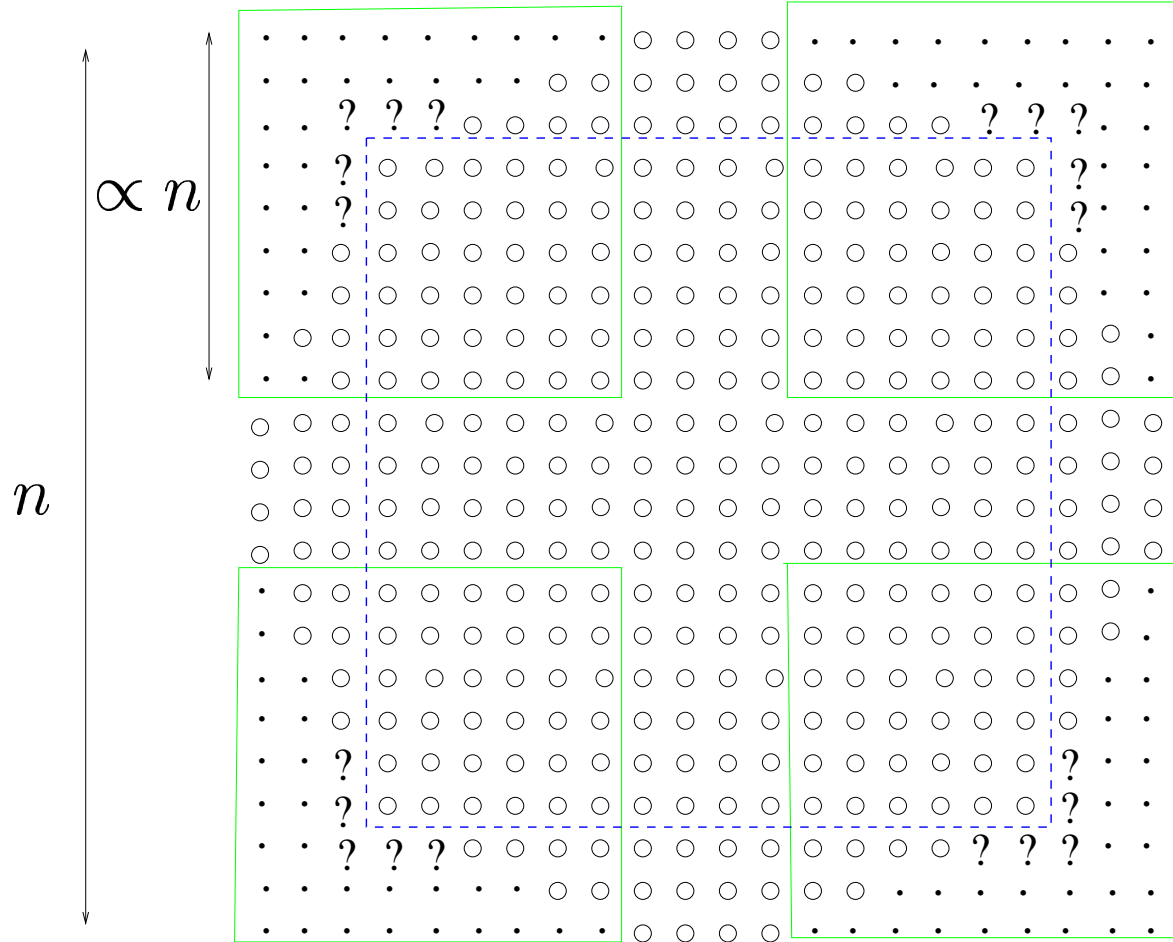
Ergodicity $\rho < \rho_c$



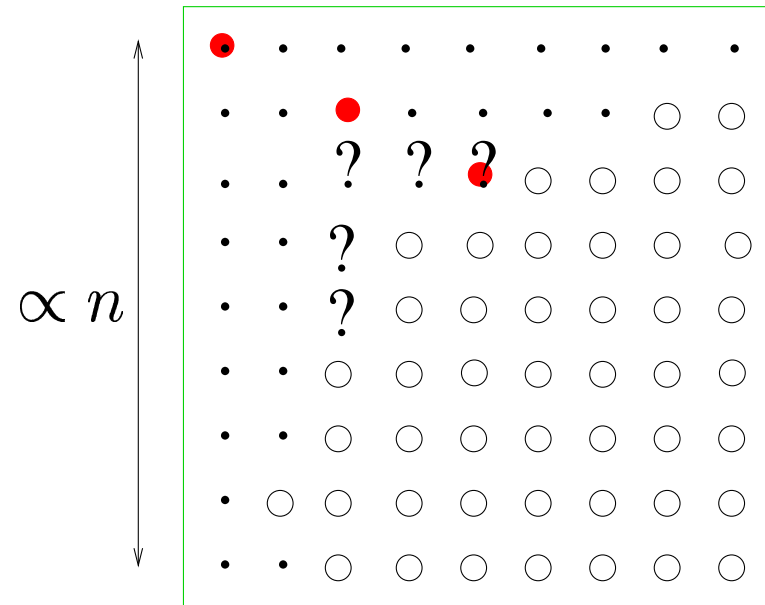
Ergodicity $\rho < \rho_c$



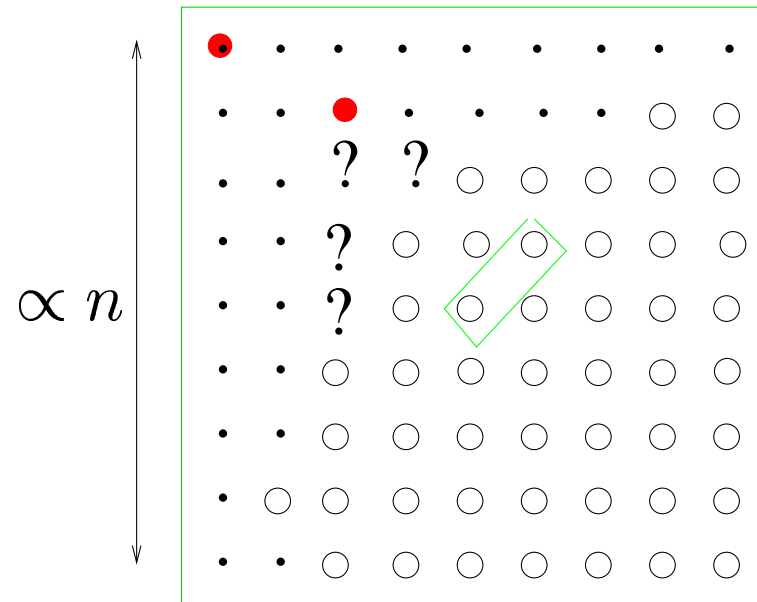
Ergodicity $\rho < \rho_c$



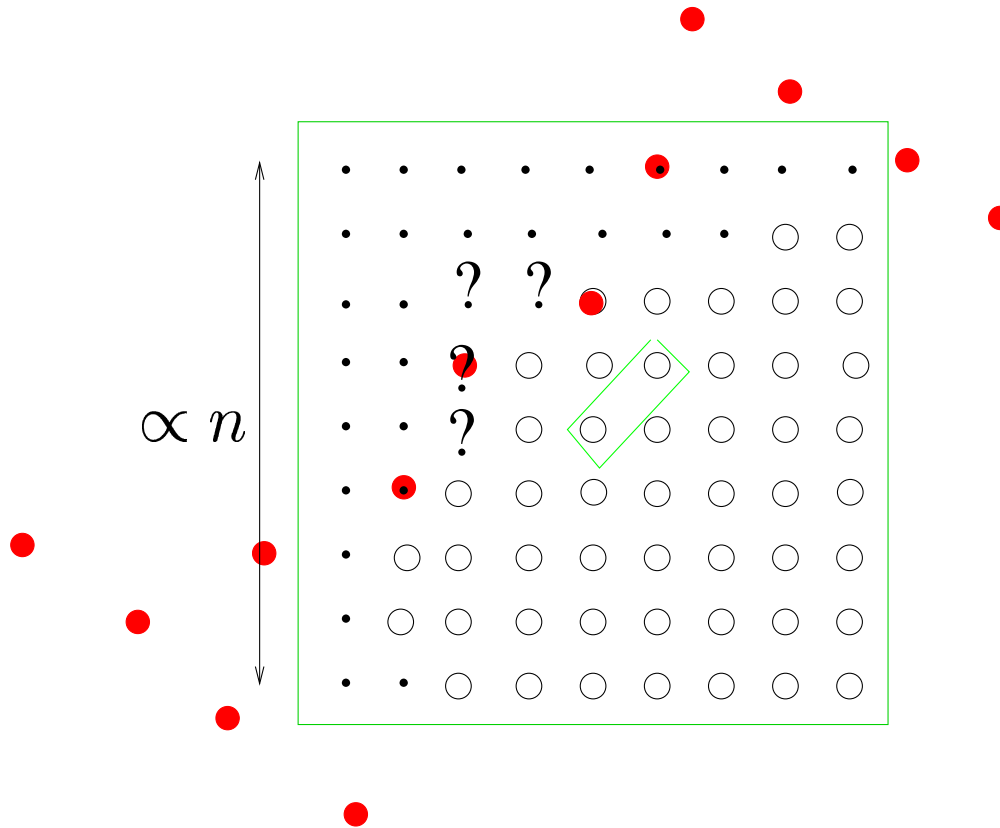
Ergodicity $\rho < \rho_c$



Ergodicity $\rho < \rho_c$



Ergodicity $\rho < \rho_c$



To block ? we need a NE-SW of length $\propto n$

Ergodicity $\rho < \rho_c$

Proba (empty square $n \times n \rightarrow (n + 2) \times (n + 2)$) \propto

Ergodicity $\rho < \rho_c$

Proba (empty square $n \times n \rightarrow (n + 2) \times (n + 2)$) \propto

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$(1 - A \exp^{-n/\xi_{\parallel}})$, $\xi_{\parallel}(\rho) =$ DP correlation length

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$\prod_n^{\infty} = P$ (empty $n \times n \rightarrow$ empty lattice) $> \exp(-c \xi_{\parallel}^{1-z})$

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- Ergodicity in infinite volume $\rho < \rho_c$
- Ergodicity breaking at ρ_c
- Critical length $\rho \nearrow \rho_c$: $L_c(\rho) \simeq \exp(\xi_{\parallel}^{1-z})$
 $\xi \propto 1/(\rho_c - \rho)^{\nu}$, $\nu \simeq 1.73$, $z \simeq 0.63$

Critical properties different from DP
Enormous finite size effects

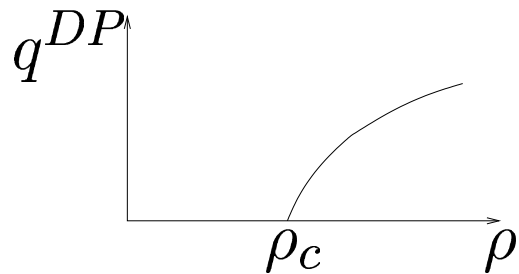
Discontinuity I

Order parameter: density of blocked clusters

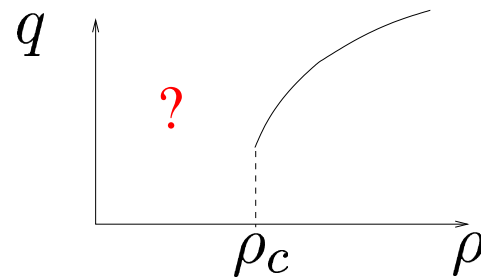
$$\rho < \rho_c: q(\rho) = 0; \quad \rho > \rho_c: q(\rho) > 0$$

$$q(\rho_c) > 0?$$

$q(\rho_c) \geq q^{DP}(\rho_c) = 0$ ($q^{DP}(\rho_c) = 0$: DP cluster fractal)

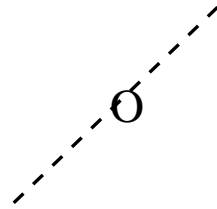


DP: continuous

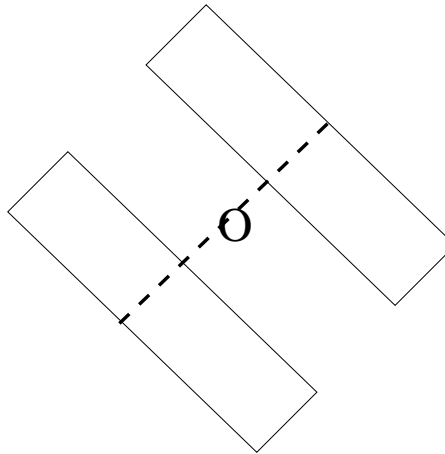


Jamming: discontinuous?

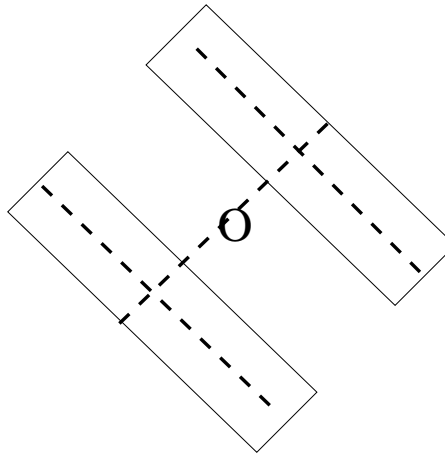
Discontinuity II



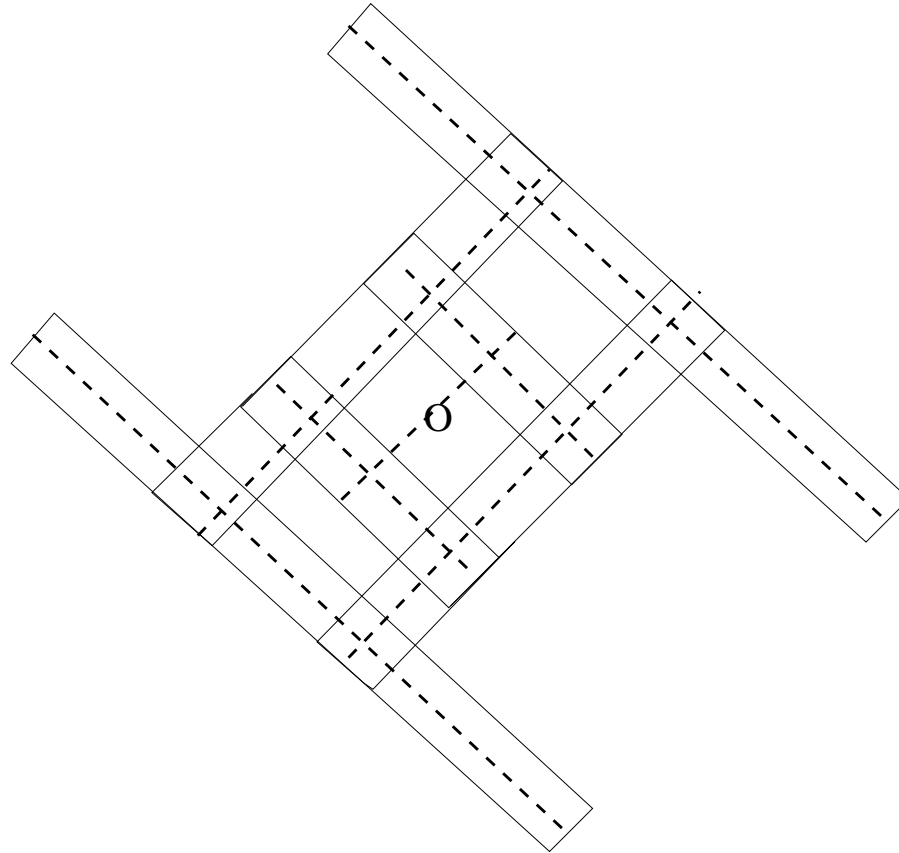
Discontinuity II



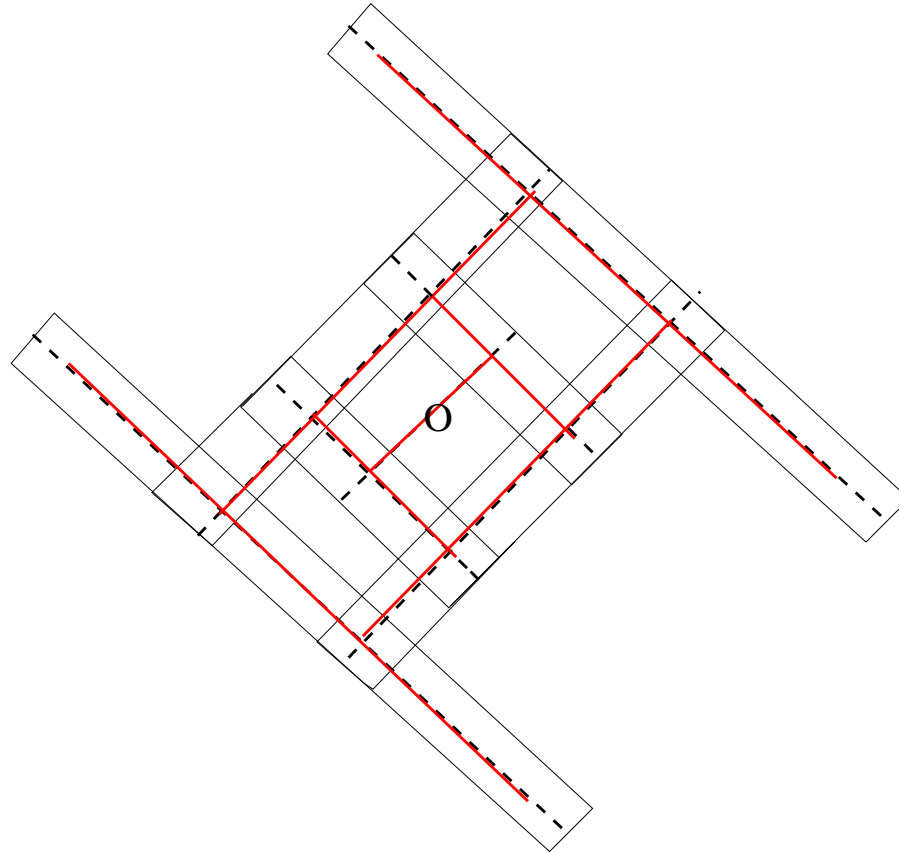
Discontinuity II



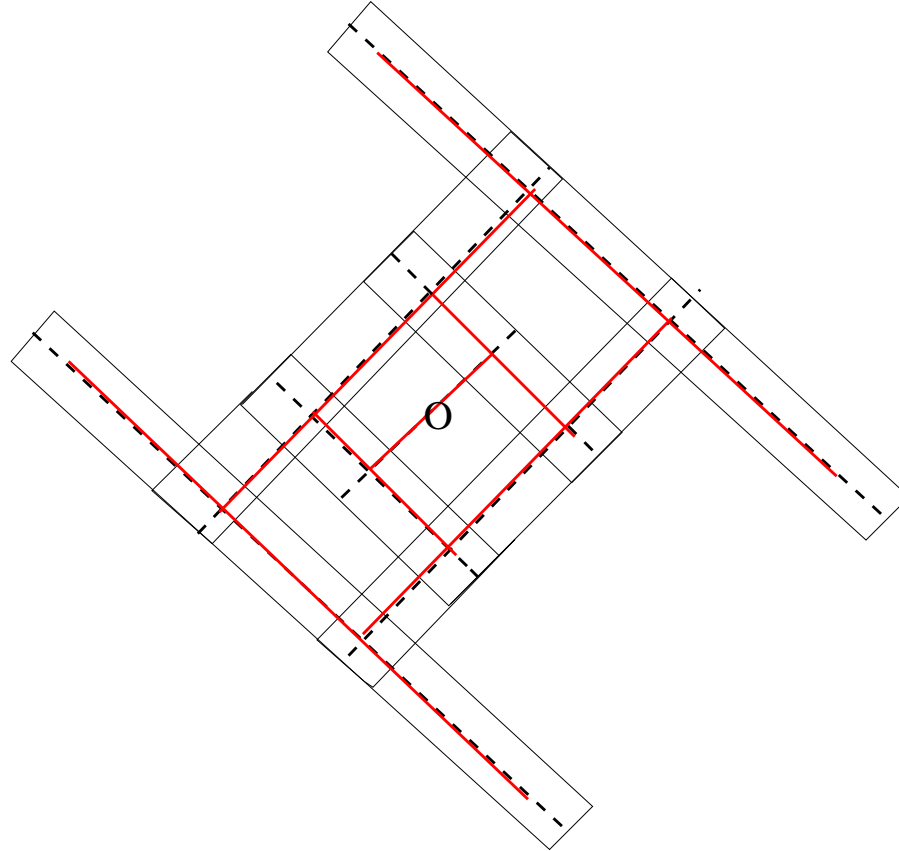
Discontinuity II



Discontinuity II



Discontinuity II



$$\lim_{L \rightarrow \infty} \text{Proba} (\exists \text{ cluster DP in } L \times cL) \rightarrow 1$$
$$\rightarrow q(\rho) = \prod_L^\infty > 0$$

Clusters are compact at $\rho_c \rightarrow$ discontinuous percolation

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- First finite dimensional model with an ideal glass transition (first order + critical)

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- First finite dimensional model with an ideal glass transition (first order + critical)
- New percolation transition: discontinuous + anomalous critical properties
- Relation with experimental data (colloidal suspensions, jamming for granular media ...);
- Out of equilibrium dynamics: aging, driven non-equilibrium.