

Quasi stationary distributions and Fleming Viot Processes

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Quasi stationary distributions (QSD)

Irreducible jump Markov process with rates $Q = (q(x, y))$ on $\Lambda \cup \{0\}$. $P_t(x, y)$ transition matrix.

Λ countable and 0 absorbing state.

Z_t is ergodic with a unique invariant measure δ_0

Law starting with μ conditioned to non absorption until time t :

$$\varphi_t^\mu(x) = \frac{\sum_{y \in \Lambda} \mu(y) P_t(y, x)}{1 - \sum_{y \in \Lambda} \mu(y) P_t(y, 0)}, \quad x \in \Lambda.$$

A *quasi stationary distribution* (QSD) is a probability measure ν on Λ satisfying

$$\varphi_t^\nu = \nu$$

ν is a **Left eigenvector** for the restriction of the matrix Q to Λ with **eigenvalue** $\lambda_\nu = -\sum_{y \in \Lambda} \nu(y)q(y, 0)$: ν must satisfy the system

$$\sum_{y \in \Lambda} \nu(y) q(y, x) = \left(-\sum_{y \in \Lambda} \nu(y)q(y, 0) \right) \nu(x), \quad \forall x \in \Lambda.$$

$$\nu Q = \lambda_\nu \nu$$

$$\sum_{y \in \Lambda} \nu(y) [q(y, x) + q(y, 0)\nu(x)] = 0, \quad \forall x \in \Lambda.$$

recall
$$q(x, x) = -\sum_{y \in \Lambda \cup \{0\} \setminus \{x\}} q(x, y)$$

$$\sum_{y \in \Lambda \setminus \{x\}} \nu(y) [q(y, x) + q(y, 0)\nu(x)] = \nu(x) \sum_{y \in \Lambda \setminus \{x\}} (q(x, y) + q(x, 0)\nu(y))$$

(balance equations)

Yaglom limit for μ :

$$\lim_{t \rightarrow \infty} \varphi_t^\mu(y), \quad y \in \Lambda$$

if it exists and it is a probability on Λ .

Λ finite, Darroch and Seneta (1967): there exists a unique QSD ν for Q and that the Yaglom limit converges to ν independently of the initial distribution.

Λ infinite: neither existence nor uniqueness of QSD are guaranteed.

Example: asymmetric random walk

$p = q(i, i + 1) = 1 - q(i, i - 1)$, for $i \geq 0$. In this case there are infinitely many QSD when $p < 1/2$ and none when $p \geq 1/2$.

Minimal QSD (for $p < 1/2$):

$$\nu(x) \sim x \left(\frac{p}{1-p} \right)^{x/2}$$

Existence

For $\Lambda = \mathbb{N}$ under the condition

$$\lim_{x \rightarrow \infty} \mathbb{P}(R > t | Z_0 = x) = 0$$

R absorption time, Ferrari, Kesten, Martínez and Picco [6]:

$$\text{existence of QSD} \iff \mathbb{E}e^{\theta R} < \infty$$

for some $\theta > 0$.

Existence

Ergodicity coefficient of Q :

$$\alpha = \alpha(Q) := \sum_{z \in \Lambda} \inf_{x \in \Lambda \setminus \{z\}} q(x, z)$$

Maximal absorbing rate of Q :

$$C = C(Q) := \sup_{x \in \Lambda} q(x, 0)$$

Theorem 1. *If $\alpha > C$ then there exists a unique QSD ν for Q and the Yaglom limit converges to ν for any initial measure μ .*

Jacka and Roberts [10]: under $\alpha > C$ uniqueness and Yaglom limit.

The Fleming-Viot process (FV) .

- System of $N > 0$ particles evolving on Λ .
- Particles move independently with rates Q between absorptions.
- When a particle is absorbed, it chooses one of the other particles uniformly and jumps instantaneously to its position.

Generator (Master equation):

$$\mathcal{L}f(\xi) = \sum_{i=1}^N \sum_{y \in \Lambda \setminus \{\xi(i)\}} \left[q(\xi(i), y) + q(\xi(i), 0) \frac{\eta(\xi, y)}{N-1} \right] (f(\xi^{i,y}) - f(\xi))$$

where $\xi^{i,y}(j) = y$ for $j = i$ and $\xi^{i,y}(j) = \xi(j)$ otherwise and

$$\eta(\xi, y) := \sum_{i=1}^N \mathbf{1}\{\xi(i) = y\}.$$

Empirical profile and conditioned process

ξ_t process in $\Lambda^{(1,\dots,N)}$;

$\eta_t = \in \mathbb{N}^\Lambda$ unlabeled process,

$\eta_t(x)$ = number of ξ particles in state x at time t .

Theorem 2. *Let μ probability on Λ . Assume*

$(\xi_0^{N,\mu}(i), i = 1, \dots, N)$ iid with law μ . Then, for $t > 0$ and $x \in \Lambda$,

$$\lim_{N \rightarrow \infty} \frac{\mathbb{E} \eta_t^{N,\mu}(x)}{N} = \varphi_t^\mu(x)$$

$$\lim_{N \rightarrow \infty} \frac{\eta_t^{N,\mu}(x)}{N} = \varphi_t^\mu(x), \quad \text{in Probability}$$

Fleming and Viot [8], Burdzy, Holyst and March [1], Grigorescu and Kang [9] and Löbus [12] in a Brownian motion setting.

Ergodicity of FV

Λ finite, FV Markov in finite state space

Hence ergodic (there exists unique stationary measure and the process converges to the stationary measure).

For Λ infinite:

Theorem 3. *If $\alpha > 0$, then for each N the FV process with N particles is ergodic.*

Stationary empirical profile and QSD

Assume ergodicity.

Let η^N be distributed with the unique invariant measure.

Theorem 4. $\alpha > C$. For each $x \in \Lambda$, the following limits exist

$$\lim_{N \rightarrow \infty} \frac{\eta^N(x)}{N} = \nu(x), \quad \text{in Probability}$$

and ν is the unique QSD for Q .

Sketch of proofs

Existence part of Theorem 1 corollary of Theorem 4. Uniqueness: Jacka and Robert.

Theorem 3: stationary version of the process “from the past” as in perfect simulation.

Theorems 2 and 4 based on asymptotic independence.

- φ_t unique solution of

$$\frac{d}{dt}\varphi_t^\mu(x) = \sum_{y \in \Lambda} \varphi_t^\mu(y)[q(y, x) + q(y, 0)\varphi_t^\mu(x)], \quad x \in \Lambda$$

- η_t satisfies

$$\frac{d}{dt}\mathbb{E}\left(\frac{\eta_t^{N,\mu}(x)}{N}\right) = \sum_{y \in \Lambda} \mathbb{E}\left(\frac{\eta_t^{N,\mu}(y)}{N} \left(q(y, x) + q(y, 0)\frac{\eta_t^{N,\mu}(x)}{N-1}\right)\right)$$

- We prove:

$$\mathbb{E}[\eta_t^{N,\mu}(y) \eta_t^{N,\mu}(x)] - \mathbb{E}\eta_t^{N,\mu}(y) \mathbb{E}\eta_t^{N,\mu}(x) = O(N)$$

- QSD satisfies

$$\sum_{y \in \Lambda} \nu(y) [q(y, x) + q(y, 0)\nu(x)] = 0, \quad x \in \Lambda.$$

- η^N invariant for FV satisfies:

$$\sum_{y \in \Lambda} \mathbb{E} \left(\frac{\eta^N(y)}{N} \left(q(y, x) + q(y, 0) \frac{\eta^N(x)}{N-1} \right) \right) = 0, \quad x \in \Lambda.$$

- Under $\alpha > C$:

$$\mathbb{E}[\eta^N(y) \eta^N(x)] - \mathbb{E}\eta^N(y) \mathbb{E}\eta^N(x) = O(N)$$

- Variance order $1/N$, setting $x = y$.
- Finally we show $(\varphi_t^{N, \mu}, N \in \mathbb{N})$ and $(\rho^N, N \in \mathbb{N})$ are tight.

Comments

- Fleming-Viot permits to show existence of a QSD in the $\alpha > C$ case (new).
- Compared with Brownian motion in a bounded region with absorbing boundary (Burdzy, Holyst and March [1], Grigorescu and Kang [9] and Löbus [12] and other related works):
 - Existence of the FV process immediate here.
 - they prove the convergence without factorization.
 - We prove: vanishing correlations sufficient for convergence of expectations and in probability.
 - To prove tightness classify ξ particles in types.
 - Tightness proof needs $\alpha > C$ as the vanishing correlations proof.

Graphical construction of FV process

To each particle $i = 1, \dots, N$, associate 3 marked Poisson processes:

- Regeneration times. PP (α) : $(a_n^i)_{n \in \mathbb{Z}}$, marks $(A_n^i)_{n \in \mathbb{Z}}$
- Internal times. PP $(\bar{q} - \alpha)$: $(b_n^i)_{n \in \mathbb{Z}}$,
marks $((B_n^i(x), x \in \Lambda), n \in \mathbb{Z})$
- Voter times. PP (C) : $(c_n^i)_{n \in \mathbb{Z}}$,
marks $((C_n^i, (F_n^i(x), x \in \Lambda)), n \in \mathbb{Z})$

Law of marks:

- $\mathbb{P}(A_n^i = y) = \alpha(y)/\alpha, y \in \Lambda;$
- $\mathbb{P}(B_n^i(x) = y) = \frac{q(x, y) - \alpha(y)}{\bar{q} - \alpha}, x \in \Lambda, y \in \Lambda \setminus \{x\};$
 $\mathbb{P}(B_n^i(x) = x) = 1 - \sum_{y \in \Lambda \setminus \{x\}} \mathbb{P}(B_n^i(x) = y).$
- $P(F_n^i(x) = 1) = \frac{q(x, 0)}{C} = 1 - P(F_n^i(x) = 0), x \in \Lambda.$
- $P(C_n^i = j) = \frac{1}{N - 1}, j \neq i.$

Call ω a realization of the marked PP.

Construction of $\xi_{[s,t]}^{N,\xi} = \xi_{[s,t],\omega}^{N,\xi}$

- Order Poisson times.
- Initial configuration ξ at time s .
- Configuration does not change between Poisson events.
- At each regeneration time a_n^i particle i adopts state A_n^i regardless the current configuration.
- If at the internal time b_n^i – the state of particle i is x , then at time b_n^i particle i adopts state $B_n^i(x)$ regardless the state of the other particles.
- If at the voter time c_n^i – the state of particle i is x and $F_n^i(x) = 1$, then at time c_n^i particle i adopts the state of particle C_n^i ; if $F_n^i(x) = 0$, then particle i does not change state.
- The final configuration is $\xi_{[s,t]}^{N,\xi}$.

Lemma 1. *The process $(\xi_{[s,t]}^{N,\xi}, t \geq s)$ is FV with initial condition*

$$\xi_{[s,s]}^{N,\xi} = \xi.$$

Generalized duality Define

$$\omega^i[s, t] = \{m \in \omega : m \text{ involved in the definition of } \xi_{[s, t], \omega}^{N, \xi}(i)\},$$

Generalized duality equation:

$$\xi_{[s, t], \omega}^{N, \xi}(i) = H(\omega^i[s, t], \xi). \quad (1)$$

- No explicit formula for H .
- For any time s , $\xi_{[s, t]}^{N, \xi}(i)$ depends only on the *finite* number of Poisson events contained in $\omega^i[s, t]$.

Theorem 3. *If $\alpha > 0$ the FV process is ergodic.*

Proof If number of marks in $\omega^i[-\infty, t]$ is finite, then

$$\xi_{t,\omega}^N(i) =: \lim_{s \rightarrow -\infty} H(\omega^i[s, t], \xi), \quad i \in \{1, \dots, N\}, \quad t \in \mathbb{R}$$

is well defined and does not depend on ξ .

- By construction $(\xi_t^N, t \in \mathbb{R})$ is a stationary FV process.
- The law of ξ_t^N is unique invariant measure.
- Number of points in $\omega^i[-\infty, t]$ is finite if there is $[s(\omega), s(\omega) + 1]$ in the past of t with one regeneration mark for each k and no voter marks. \square

Particle correlations in the FV process

Proposition 1. *Let $x, y \in \Lambda$. For all $t > 0$*

$$\left| \mathbb{E} \left(\frac{\eta_t^N(x) \eta_t^N(y)}{N^2} \right) - \mathbb{E} \left(\frac{\eta_t^N(x)}{N} \right) \mathbb{E} \left(\frac{\eta_t^N(y)}{N} \right) \right| < \frac{1}{N} e^{2Ct} \quad (2)$$

Assume $\alpha > C$. Let η^N be distributed according to the unique invariant measure for the FV process with N particles. Then

$$\left| \mathbb{E} \left(\frac{\eta^N(x) \eta^N(y)}{N^2} \right) - \mathbb{E} \left(\frac{\eta^N(x)}{N} \right) \mathbb{E} \left(\frac{\eta^N(y)}{N} \right) \right| < \frac{1}{N} \frac{\alpha}{\alpha - C} \quad (3)$$

Coupling

- 4-fold coupling $(\omega^i[0, t], \omega^j[0, t], \hat{\omega}^i[0, t], \hat{\omega}^j[0, t])$
- $\omega^i[0, t] = \hat{\omega}^i[0, t]$
- $\hat{\omega}^j[0, t] \cap \omega^i[0, t] = \emptyset$ implies $\omega^j[0, t] = \hat{\omega}^j[0, t]$
- marginal process $(\hat{\omega}^i[0, t], \hat{\omega}^j[0, t])$ have the same law as two independent processes with the same marginals as $(\omega^i[0, t], \omega^j[0, t])$.

$$\begin{aligned} & \mathbb{P}(\xi_t^{N,\xi}(j) = x, \xi_t^{N,\xi}(i) = y) - \mathbb{P}(\xi_t^{N,\xi}(j) = x)\mathbb{P}(\xi_t^{N,\xi}(i) = y) \\ &= \mathbb{E}\left(\mathbf{1}\{H(\omega^j, \xi) = x, H(\omega^i, \xi) = y\} - \mathbf{1}\{H(\hat{\omega}^j, \xi) = x, H(\hat{\omega}^i, \xi) = y\}\right) \end{aligned}$$

• If

$$\omega^i \cap \omega^j = \emptyset$$

then

$$\omega^j(s, t) = \hat{\omega}^j(s, t) \text{ and } \omega^i(s, t) = \hat{\omega}^i(s, t)$$

Hence,

$$\begin{aligned} & |\mathbb{P}(\xi_t^{N,\xi}(j) = x, \xi_t^{N,\xi}(i) = y) - \mathbb{P}(\xi_t^{N,\xi}(j) = x)\mathbb{P}(\xi_t^{N,\xi}(i) = y)| \\ & \leq \mathbb{P}(\omega^i \cap \omega^j \neq \emptyset). \end{aligned}$$

Lemma 2.

$$\mathbb{P}(\omega^i \cap \omega^j \neq \emptyset) \leq \frac{1}{N-1} \frac{C}{\alpha - C} (1 - e^{2(C-\alpha)t}) \quad (4)$$

Proof:

$$\mathbb{P}(\omega^i \cap \omega^j \neq \emptyset) \leq \frac{2C}{N-1} \int_0^t \mathbb{E} \hat{\Psi}^i[s, t] \mathbb{E} \hat{\Psi}^j[s, t] ds$$

$\hat{\Psi}^i[s, t]$ Random walk that grows with rate C and decreases with rate α . Expectation is bounded by $e^{(t-s)(C-\alpha)}$.

$$\mathbb{P}(\omega^i \cap \omega^j \neq \emptyset) \leq \frac{2C}{N-1} \int_0^t e^{2(C-\alpha)s} ds$$

which gives the result. \square

Proof of Proposition 1 Take η and ξ such that $\eta(x) = \sum_j \mathbf{1}\{\xi(j) = x\}$. Then

$$\mathbb{E}\left(\frac{\eta_t^{N,\eta}(x)\eta_t^{N,\eta}(y)}{N^2}\right) = \frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \mathbb{P}(\xi_t^{N,\xi}(i) = x, \xi_t^{N,\xi}(j) = y)$$

$$\frac{\mathbb{E}\eta_t^{N,\eta}(x) \mathbb{E}\eta_t^{N,\eta}(y)}{N^2} = \frac{1}{N^2} \left(\sum_{i=1}^N \sum_{j=1}^N \mathbb{P}(\xi_t^{N,\xi}(i) = x) \mathbb{P}(\xi_t^{N,\xi}(j) = y) \right)$$

Using this, and (4) with $\alpha = 0$ we get (2).

Assume $\alpha > C$. Taking $t = \infty$ in (4) we get (3). \square

Tightness

Proposition 2. *For all $t > 0$, $x \in \Lambda$, $i = 1, \dots, N$ and μ ,*

$$\frac{\mathbb{E}\eta_t^{N,\mu}(x)}{N} \leq e^{Ct} \sum_{z \in \Lambda} \mu(z) P_t(z, x).$$

As a consequence $(\mathbb{E}\eta_t^{N,\mu}/N, N \in \mathbb{N})$ is tight.

Assume $\alpha > 0$ and define μ_α on Λ by

$$\mu_\alpha(x) = \frac{\alpha_x}{\alpha}, \quad x \in \Lambda,$$

where $\alpha_x = \inf_z q(z, x)$. For $z, x \in \Lambda$ define

$$R_\lambda(z, x) = \int_0^\infty \lambda e^{-\lambda t} P_t(z, x) dt.$$

Proposition 3. *Assume $\alpha > C$ and let η^N distributed with invariant measure for FV. Then for $x \in \Lambda$,*

$$\rho^N(x) \leq \frac{C}{\alpha - C} \mu_\alpha R_{(\alpha - C)}(x)$$

As a consequence, the family of measures $(\eta^N / N, N \in \mathbb{N})$ is tight.

Types

- Particle i is *type* 0 at time t if it has not been absorbed in the time interval $[0, t]$.
- If at absorption time s particle i jumps over particle j which has type k , then at time s particle i changes its type to $k + 1$.

$$\mathbb{P}(\xi_t^{N,\mu}(i) = x, \text{type}(i, t) = 0) = \sum_{z \in \Lambda} \mu(z) P_t(z, x).$$

$$A_t(x, k) =: \mathbb{P}(\xi_t^{N,\mu}(i) = x, \text{type}(i, t) = k)$$

Proof of Proposition 2 Recursive hypothesis:

$$A_t(x, k) \leq \frac{(Ct)^k}{k!} \sum_{z \in \Lambda} \mu(z) P_t(z, x) \quad (5)$$

By (5) the statement is true for $k = 0$.

$$A_t(x, k + 1) \leq \int_0^t C \sum_{y \in \Lambda} A_s(y, k) P_{t-s}(y, x) ds.$$

Using recursive hypothesis,

$$\begin{aligned} &= \int_0^t C \frac{(Cs)^k}{k!} \sum_{z \in \Lambda} \mu(z) \sum_{y \in \Lambda} P_s(z, y) P_{t-s}(y, x) ds \\ &= \frac{(Ct)^{k+1}}{(k+1)!} \sum_{z \in \Lambda} \mu(z) P_t(z, x). \end{aligned}$$

by Chapman-Kolmogorov. This proves (5). \square

Proof of Proposition 3 Under the hypothesis $\alpha > C$ the process

$$((\xi_t^N(i), \text{type}(i, t)), i = 1, \dots, N), t \in \mathbb{R})$$

is Markovian constructed in a stationary way

$$A(x, k) := \mathbb{P}(\xi_s^N(i) = x, \text{type}(i, s) = k)$$

does not depend on s .

Last regeneration mark of site i before time s happened at time $s - T_\alpha^i$, where T_α^i is exponential of rate α . Then,

$$A(x, 0) = \int_0^\infty \alpha e^{-\alpha t} \sum_{z \in \Lambda} \mu_\alpha(z) P_t(z, x) dt = \mu_\alpha R_\alpha(x).$$

Similar reasoning implies

$$\begin{aligned}
 A(x, k) &\leq \int_0^\infty e^{-\alpha t} C \sum_{z \in \Lambda} A(z, k-1) P_s(z, x) dt. \\
 &= \frac{C}{\alpha} A_{k-1} R_\alpha(x) \leq \left(\frac{C}{\alpha}\right)^k \mu_\alpha R_\alpha^{k+1}(x).
 \end{aligned}$$

$R_\lambda^k(z, x)$ expectation of $P_{\tau_k}(z, x)$, τ_k sum of k independent exponential λ . Multiplying and dividing by $(\alpha - C)$,

$$\mathbb{P}(\xi_s^N(i) = x) \leq \frac{C}{\alpha - C} \sum_{k=0}^\infty \left(\frac{C}{\alpha}\right)^k \left(1 - \frac{C}{\alpha}\right) \mu_\alpha R_\alpha^{k+1}(x)$$

Expectation of $\mu_\alpha R_\alpha^K$, K geometric with $p = 1 - (C/\alpha)$.

$$\mathbb{P}(\xi_s^N(i) = x) \leq \frac{C}{\alpha - C} \mu_\alpha R_{\alpha - C}(x). \quad \square$$

References

- [1] Burdzy, K., Holyst, R., March, P. (2000) A Fleming-Viot particle representation of the Dirichlet Laplacian, *Comm. Math. Phys.* **214**, 679-703.
- [2] Cavender, J. A. (1978) Quasi-stationary distributions of birth and death processes. *Adv. Appl. Prob.* **10**, 570-586.
- [3] Darroch, J.N., Seneta, E. (1967) On quasi-stationary distributions in absorbing continuous-time finite Markov chains, *J. Appl. Prob.* **4**, 192-196.
- [4] Ferrari, P.A. (1990) Ergodicity for spin systems with stirrings, *Ann. Probab.* **18**, 4:1523-1538.
- [5] Fernández, R., Ferrari, P.A., Garcia, N. L. (2001) Loss network representation of Peierls contours. *Ann. Probab.* **29**, no. 2, 902–937.

- [6] Ferrari, P.A, Kesten, H., Martínez,S., Picco, P. (1995) Existence of quasi stationary distributions. A renewal dynamical approach, *Ann. Probab.* **23**, 2:511–521.
- [7] Ferrari, P.A., Martínez, S., Picco, P. (1992) Existence of non-trivial quasi-stationary distributions in the birth-death chain, *Adv. Appl. Prob* **24**, 795-813.
- [8] Fleming,.W.H., Viot, M. (1979) Some measure-valued Markov processes in population genetics theory, *Indiana Univ. Math. J.* **28**, 817-843.
- [9] Grigorescu, I., Kang,M. (2004) Hydrodynamic limit for a Fleming-Viot type system, *Stoch. Proc. Appl.* **110**, no.1: 111-143.
- [10] Jacka, S.D., Roberts, G.O. (1995) Weak convergence of conditioned processes on a countable state space, *J. Appl. Prob.* **32**, 902-916.

- [11] Kipnis, C., Landim, C. (1999) *Scaling Limits of Interacting Particle Systems*, Springer-Verlag, Berlin.
- [12] Löbus, J.-U. (2006) A stationary Fleming-Viot type Brownian particle system. Preprint.
- [13] Nair, M.G., Pollett, P. K. (1993) On the relationship between μ -invariant measures and quasi-stationary distributions for continuous-time Markov chains, *Adv. Appl. Prob.* **25**, 82-102.
- [14] Seneta, E. (1981) *Non-Negative Matrices and Markov Chains*. Springer-Verlag, Berlin.
- [15] Seneta, E. , Vere-Jones, D. (1966) On quasi-stationary distributions in discrete-time Markov chains with a denumerable infinity of states, *J. Appl. Prob.* **3**, 403-434.
- [16] Vere-Jones, D. (1969) Some limit theorems for evanescent processes, *Austral. J. Statist* **11**, 67-78.

- [17] Yaglom, A. M. (1947) Certain limit theorems of the theory of branching stochastic processes (in Russian). *Dokl. Akad. Nauk SSSR* (n.s.) **56**, 795-798.