

Introduction to the  
Painlevé equations

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## 6 Painlevé equations

$$1898 \leq \begin{array}{c} \text{Painlevé} \\ \text{Gambier} \end{array} \leq 1906$$

Classification of the 2<sup>nd</sup> order  
nonlinear ODE without movable  
critical pts.

ex. 
$$\frac{d^2y}{dx^2} = \frac{k+1}{y} \left( \frac{dy}{dx} \right)^2$$

$$y = (ax+b)^{-\frac{1}{k}}$$

$$a, b \in \mathbb{C}.$$

Motivation:

Find a new-transcendental  
function (in Math. Phy)  
defined by a differential  
equation.

$$e^x$$

$$\sin x$$

$$\log x$$

$$J_\nu(x)$$

$$F(a, b, c; x)$$

...

$$f(x; \omega_1, \omega_3)$$

# Key Words :

Irreducibility

Hamiltonian System

Holonomic Deformation

(Isomonodromic Deformation)

Birational Transformation

(Bäcklund Transformation)

Weyl Group

Roots Systems ..

$\tau$ -function & Bilinear form

Space of initial conditions.

Hamiltonian System :

$$\delta = f(t) \frac{d}{dt} \quad f(t) \in \mathbb{C}(t)$$

$$H(t; q, p) \in \mathbb{C}[t, q, p]$$

$$\delta q = \frac{\partial H}{\partial p}, \quad \delta p = -\frac{\partial H}{\partial q} \quad (*)$$

If  $H$  is of 2<sup>nd</sup> deg. w.r.t.  $p$

then, by eliminating  $p$  from the systems,  
we obtain :

$$\delta^2 q = R(t; q, \delta q)$$

$R$ : rational

Assumption: (\*) is without movable  
critical pts.



Ex.  $\delta = \frac{d}{dt}$ ,  $H = \frac{1}{2}p^2 - 2q^3 + \frac{1}{2}g_2q$   
 $g_2 \in \mathbb{C}$

Hamiltonian system defines  
 Weierstrass  $p$ -fnc.

$K / \mathbb{C}(t)$  differential field extension

$L = K(q, p)$  : diff. field w.r.t.  $\delta$

through  $\delta q = \frac{\partial H}{\partial p}$ ,  $\delta p = -\frac{\partial H}{\partial q}$ .

Theorem (Painlevé)

Under the assumption.

$\text{trans. deg}_{\mathbb{C}} L = 2$

$\Rightarrow$  Hamiltonian system is reduced  
 to one of the Painlevé equations,

$P_I, P_{II}, \dots, P_{VI}$ .

$$\text{Ex. } H = \frac{1}{2}p^2 - (2q^3 + tq)$$

$$\delta = \frac{d}{dt}, \quad \delta q = \frac{\partial H}{\partial p} = p$$

$$\delta p = -\frac{\partial H}{\partial q} = 6q^2 + t$$

$$\Rightarrow \delta^2 q = 6q^2 + t \quad (P_I)$$

$$P_I \quad P_{II} \quad P_{III} \quad P_{IV} \quad P_V \quad P_{VI}$$

$$\delta: \frac{d}{dt} \quad \frac{d}{dt} \quad t \frac{d}{dt} \quad \frac{d}{dt} \quad t \frac{d}{dt} \quad t(t-1) \frac{d}{dt}$$

$$\text{Ex. } H = 2q^2 p^2 - (2tq^2 + (2\theta_0 + 1)q - 2t)p + (\theta_0 + \theta_\infty)tq$$

$$\delta = t \frac{d}{dt}$$

$$\Rightarrow P_{III}'$$

Holonomic deformation

(Monodromy Preserving deformation)

Problem: Given

$$\frac{d^2y}{dx^2} + P_1(x;t) \frac{dy}{dx} + P_2(x;t)y = 0,$$

Find two rational functions in  $x$

$$a_1(x;t), a_2(x;t)$$

such that

$$\left\{ \begin{array}{l} \frac{\partial^2 y}{\partial x^2} + P_1 \frac{\partial y}{\partial x} + P_2 y = 0 \\ \frac{\partial y}{\partial t} = a_1 y + a_2 \frac{\partial y}{\partial x} \end{array} \right.$$

is completely integrable.

←  
extended system.



# Monodromy Preserving deformation

( R. Fuchs, R. Garnier, K.O )

( L. Scheleginger, T. Miwa & M. Jimbo )

$$\frac{d^2y}{dx^2} + p_1(x,t) \frac{dy}{dx} + p_2(x,t) y = 0$$

$$p_1 = \frac{1-\kappa_0}{x} + \frac{1-\kappa_1}{x-1} + \frac{1-\theta}{x-t} - \frac{1}{x-q}$$

$$p_2 = \frac{\kappa}{x(x-1)} - \frac{H}{x(x-1)(x-t)} + \frac{q(q-1)p}{x(x-1)(x-q)}$$

$x=q$  : apparent singularity

⇓

$$H = q(q-1)(q-t)p^2$$

$$- \{ \kappa_0(q-1)(q-t) + \kappa_1 q(q-t) + (\theta-1)q(q-1) \} p + \kappa(q-t)$$

$$(*) \left\{ \begin{array}{l} \frac{\partial^2 y}{\partial x^2} + p_1 \frac{\partial y}{\partial x} + p_2 y = 0 \\ \frac{\partial y}{\partial t} = a_1(x;t)y + a_2 \frac{\partial y}{\partial x} \end{array} \right.$$

$\exists a_1(x;t), a_2(x;t)$  rational in  $x$

s.t.  $(*)$  is completely integrable



$$\delta q = \frac{\partial H}{\partial p}, \quad \delta p = -\frac{\partial H}{\partial q}. \quad P_H$$

Riemann Scheme

$x=0$	$x=1$	$x=t$	$x=g$	$x=\infty$
$0$	$0$	$0$	$0$	$\downarrow$
$\kappa_0$	$\kappa_1$	$\theta$	$\alpha$	$\nu + \kappa_{\infty}$

$$2\nu + \kappa_0 + \kappa_1 + \theta + \kappa_{\infty} = 1$$

$$\kappa = \nu(\nu + \kappa_{\infty})$$

# Confluence of singularities

$$t \Rightarrow 1 + \epsilon t$$

$$\kappa_1 \Rightarrow \epsilon^{-1} + \theta + \epsilon$$

( $\alpha$ -method)

$$\theta \Rightarrow -\epsilon^{-1}$$

$$\epsilon \rightarrow 0$$

$$P_1(x; t) = \frac{1 - \kappa_0}{x} + \frac{t}{(x-1)^2} + \frac{1 - \theta}{x-1} - \frac{1}{x-2}$$

$$P_2(x; t) = \frac{\kappa}{x(x-1)} - \frac{H}{x(x-1)^2} + \frac{q(q-1)p}{x(x-1)(x-2)}$$

$$H = q(q-1)^2 p^2 - \{ \kappa_0 (q-1)^2 + \theta q(q-1) - tq \} p + \kappa(q-1)$$

$$\delta q = \frac{\partial H}{\partial p}, \quad \delta p = -\frac{\partial H}{\partial q} \quad ; \quad P_{\nabla}$$

$$\delta = t \frac{d}{dt}$$

$$1 + 1 + 1 + 1$$

$$|$$

$$2 + 1 + 1$$

$$P_{\text{IV}} \rightarrow P_{\text{V}} \rightarrow P_{\text{III}}$$

$$\begin{array}{c} \searrow \\ P_{\text{IV}} \end{array} \xrightarrow{\begin{array}{c} \swarrow \\ P_{\text{III}} \end{array}} P_{\text{II}} \rightarrow P_{\text{I}}$$

$$|+|+|+| \rightarrow 2+|+| \rightarrow 2+2$$

$$\begin{array}{c} \searrow \\ 3+1 \end{array} \xrightarrow{\begin{array}{c} \swarrow \\ 2+2 \end{array}} 4 \rightarrow \frac{7}{2}$$

$$H = 2q p^2 - (q^2 + 2tq + 2\kappa_0) p + \theta_{\infty} q$$

$$\delta = \frac{d}{dt}$$

$$\delta q = \frac{\partial H}{\partial p}, \quad \delta p = -\frac{\partial H}{\partial q} \quad P_{TV}$$

$$H = H(\kappa_0, \theta_{\infty})$$

$$q_1 = q \frac{q p - \theta_{\infty}}{q p - \kappa_0}$$

$$H_1 = H$$

$$p_1 = p \frac{q p - \kappa_0}{q p - \theta_{\infty}}$$

$$H_1 = H(\theta_{\infty}, \kappa_0)$$

$$\bar{p} = p - \frac{1}{2} q - t, \quad \bar{A} = q \bar{p} \quad \bar{H} = H + q.$$

$$q_2 = 2\bar{p} \frac{\bar{A} - \kappa_0}{\bar{A} - \kappa_0 + \theta_{\infty} + 1}$$

$$p_2 = -\frac{1}{2} p \cdot \frac{\bar{A} - \kappa_0 + \theta_{\infty} + 1}{\bar{A}}$$

$$H_2 = \bar{H}$$

$$H_2 = H(\kappa_0, \theta_{\infty} + 1)$$



# Group of Transformations:

For  $P_{II}$

$$\sigma_1: \alpha \mapsto 1 - \alpha, \quad \sigma_2: \alpha \mapsto -\alpha$$

$$G = \langle \sigma_1, \sigma_2 \rangle \cong \tilde{W}_a(A_1)$$

affine Weyl Group.

$P_I$     $P_{II}$     $P_{III}$     $P_{IV}$     $P_V$     $P_{VI}$

-    $A_1$     $A_1 \times A_1$     $A_2$     $A_3$     $D_4$

# Special Solutions

$$H = \frac{1}{2}p^2 - \left(q^2 + \frac{t}{2}\right)p - \left(\alpha + \frac{1}{2}\right)q$$

$P_{II}$

$$\frac{dq}{dt} = 2q^3 + tq + \alpha$$

if  $\alpha \in \mathbb{R}$

$C: [-\frac{1}{2}, 0]$  fundamental region of  $G$ .

$C$ : Weyl chamber

when  $\alpha \in \partial C$ , we have:

$$H(-\frac{1}{2}) \quad p=0 \quad \frac{dq}{dt} = -q^2 - \frac{1}{2}t$$

$$q = \frac{d}{dt} \log u$$

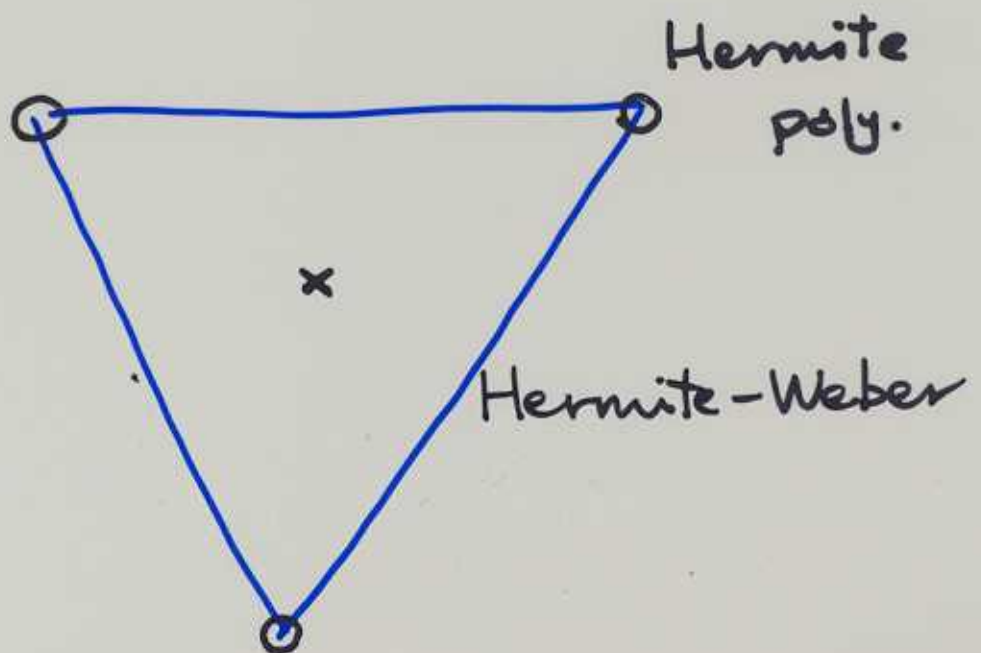
$$\frac{d^2 u}{dt^2} + \frac{1}{2}tu = 0$$

$$H(0) \quad q=0, \quad p = \frac{t}{2}$$

Riccati-typ sol. & rational sol.

$P_I$	x
$P_{II}$	Airy
$P_{III}$	Bessel
$P_{IV}$	Hermite - Weber
$P_V$	Confluent - Hypergeometric
$P_{VI}$	Gauß Hypergeometric.

$P_{IV}$



# Bäcklund Transformation

$$H = \frac{1}{2}p^2 - \left(q^2 + \frac{t}{2}\right)p - \left(\alpha + \frac{1}{2}\right)q$$

$$H = H(\alpha), \quad \delta = \frac{d}{dt}$$

$$\delta q = \frac{\partial H}{\partial p}, \quad \delta p = -\frac{\partial H}{\partial q}$$

$P_{II}$

$$\frac{d^2 q}{dt^2} = 2q^3 + tq + \alpha.$$

$$q_1 = q + \frac{\alpha + \frac{1}{2}}{p} \quad p_1 = p$$

$$\delta q_1 = \frac{\partial H_1}{\partial p_1} \quad \delta p_1 = -\frac{\partial H_1}{\partial q_1} \quad H_1 = H$$

$$H_1 = H(\alpha - 1)$$

$$p = \bar{p} + 2q^2 + t, \quad H = \bar{H} - q$$

$$q_2 = -q + \frac{\alpha - \frac{1}{2}}{p} \quad p_2 = -\bar{p} \quad H_2 = \bar{H}$$

$$H_2 = H(\alpha - 1)$$

$\tau$ -functions

$$\delta^* \log \tau = H.$$

$$P_I : \quad q = -\frac{d^2}{dt^2} \log \tau$$

$$P_{II} : \quad q = \frac{d}{dt} \log \frac{\tau(\alpha-1)}{\tau(\alpha)}$$

$\ell$ : translation in  $G$ .

$$\ell^* : \tau_m \rightarrow \tau_{m+1}$$

$\tau$ -sequence.

$$\delta^2 \log \tau_m = \frac{\tau_{m-1} \tau_{m+1}}{\tau_m^2}$$

up to elementary change  
of Hamiltonian function.



# Painlevé Analysis.

$$\frac{d^2 q}{dt^2} = 2q^3 + tq + \alpha$$

$$q = \frac{1}{z} - \frac{t_0}{6}z - \frac{\alpha+1}{4}z^2 + hz^3 + \dots$$

$$z = t - t_0, \quad h \in \mathbb{C}$$

$$\left( q = -\frac{1}{z} + \frac{t_0}{6}z - \frac{\alpha-1}{4}z^2 + hz^3 + \dots \right)$$

$$q = Q^{-1},$$

$$Q = z \left[ 1 + \frac{t_0}{6}z^2 + \frac{\alpha+1}{4}z^3 + \left( \frac{t_0^2}{36} - h \right) z^4 + \dots \right]$$

$$z = Q \left[ 1 - \frac{t_0}{6}Q^2 - \frac{\alpha+1}{4}Q^3 + \left( \frac{t_0^2}{18} + h \right) Q^4 + \dots \right]$$

$$\frac{dq}{dt} + q^2 + \frac{t}{2} = -\left(\alpha + \frac{1}{2}\right)z + \left(\frac{t_0^2}{36} + 5h\right)z^2 + \dots$$

$$= -\left(\alpha + \frac{1}{2}\right)Q - Q^2 P$$

$$\frac{dQ}{dt} = 1 + \frac{t}{2}Q^2 + (\alpha + \frac{1}{2})Q^3 + Q^4P.$$

$$\downarrow \quad \frac{d^2Q}{dt^2} = \frac{2}{Q} \left(\frac{dQ}{dt}\right)^2 - \frac{2}{Q} - tQ - \alpha Q$$

$$\frac{dQ}{dt} = \frac{\partial H}{\partial P}, \quad \frac{dP}{dt} = -\frac{\partial H}{\partial Q}.$$

$$H = \frac{1}{2}Q^4P^2 + \left[\alpha + \frac{1}{2}Q^3 + \frac{t}{2}Q^2 + 1\right]P \\ + \frac{2\alpha+1}{4} \left[ (\alpha + \frac{1}{2})Q^2 + tQ \right].$$

$$q = Q^{-1}, \quad p = -(\alpha + \frac{1}{2})Q - Q^2P$$

$$H = \frac{1}{2}p^2 - \left(q^2 + \frac{t}{2}\right)p - (\alpha + \frac{1}{2})q.$$

Space of initial conditions

$$H = \frac{1}{2}p^2 - \left(q^2 + \frac{t}{3}\right)p - \left(\alpha + \frac{1}{3}\right)q$$

$$q(t_0) = q_0, p(t_0) = p_0 \quad (q_0, p_0) \in \mathbb{C}^2$$

$$q(t) = \frac{1}{3} - \frac{t_0}{6}z - \frac{\alpha+1}{4}z^2 + hz^3 + \dots$$

$$\mathbb{C}^2 \cup \{h\}_{h \in \mathbb{C}}.$$

$\bar{X}$ : compactification of  $\mathbb{C}^2$

Hamiltonian System defined

foliation in  $\bar{X}$

(with sing.)

$(\bar{X}, D)$

$$X = \bar{X} \setminus D$$

$\forall x \in X \quad \exists_1$  sol. of the System.

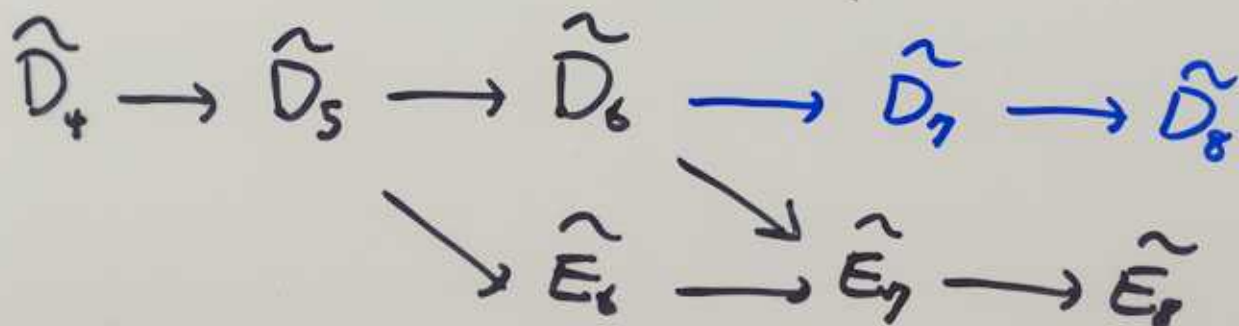
$\forall$  sol. of the system,  $\exists_1 x \in X$

$D$ : set of vertical leaves.

$P_I$	$P_{II}$	$P_{III}$	$P_{IV}$	$P_V$	$P_{VI}$
$\hat{E}_8$	$\hat{E}_7$	$\hat{D}_6$	$\hat{E}_6$	$\hat{D}_5$	$\hat{D}_4$



Sakai-san's theory.





$$H = \frac{1}{2}p^2 - (q^2 + \frac{t}{2})p - (\alpha + \frac{1}{2})q$$

$$\frac{dH}{dt} = -\frac{1}{2}p$$

$$\frac{d^2H}{dt^2} = -qp - \frac{1}{2}(\alpha + \frac{1}{2})$$

$$\left(\frac{d^2H}{dt^2}\right)^2 + 4\left(\frac{dH}{dt}\right)^3 - 2\frac{dH}{dt}\left(H - t\frac{dH}{dt}\right) - \frac{1}{4}\left(\alpha + \frac{1}{2}\right)^2 = 0$$

$$\bar{H} = \frac{1}{2}\bar{p}^2 + (q^2 + \frac{t}{2})\bar{p} - (\alpha - \frac{1}{2})q$$

$$\left(\frac{d\bar{H}}{dt}\right)^2 + 4\left(\frac{d\bar{H}}{dt}\right)^3 - 2\frac{d\bar{H}}{dt}\left(\bar{H} - t\frac{d\bar{H}}{dt}\right) - \frac{1}{4}\left(\alpha - \frac{1}{2}\right)^2 = 0$$



# Bilinear forms

$$\frac{d^3 H}{dt^3} + 6 \left( \frac{dH}{dt} \right)^2 + 2 \frac{dH}{dt} - H = 0$$

$$\mathcal{D}^4 \tau \cdot \tau + 2 \mathcal{D}^2 \tau \cdot \tau - 2 \tau' \cdot \tau = 0$$

$$H = \frac{d}{dt} \log \tau.$$

$$\bar{H} = \frac{d}{dt} \log \bar{\tau}$$

$$\mathcal{D}^2 \bar{\tau} \cdot \tau + \frac{1}{2} \bar{\tau} \cdot \tau = 0$$

$$\mathcal{D}^3 \bar{\tau} \cdot \tau + \frac{1}{2} \mathcal{D} \bar{\tau} \cdot \tau = \alpha \bar{\tau} \cdot \tau$$