

## Is there an Interesting Index Theory for Quantum Graphs?

*Roth* (1983):      Kirchhoff BC

$$\sum_{n=0}^{\infty} e^{-\lambda_n t} = \text{Tr } K = \int_{\Gamma} K(t, x, x) dx$$

= sum over closed orbits =  $K_1 + K_2 + K_3$ .

1. Zero length:       $K_1 = \frac{L}{\sqrt{4\pi t}}$       (Weyl's law).
2. Periodic:       $K_2 = \frac{1}{\sqrt{4\pi t}} \sum_C \mathcal{A}(C) L(C_p) e^{-L(C)^2/4t}$ .
3. Closed:      “Ce calcul est un peu plus délicat.”

$$K_3 = \frac{1}{2}(V - E) \quad (\text{the **only** constant term}).$$

*Justin Wilson* (2007):      any real S-matrix

$$\begin{aligned} K_3 &= \frac{1}{2} \sum_{\alpha} S_{\bar{\alpha}\alpha} \int_0^{\infty} K_0(t, x, 0) dx \\ &= \frac{1}{4} \sum_{\alpha} S_{\bar{\alpha}\alpha} \quad (\text{sum all bond backscatter amplitudes}). \end{aligned}$$

$$\text{Kirchhoff:} \quad K_3 = \frac{1}{4} \sum_v \left( \frac{2}{N_v} - 1 \right) N_v = \frac{1}{2}V - \frac{1}{2}E.$$

$V - E$  is

- an integer.
- “topological” (Euler characteristic of a 1-complex).

So what? ...

Gilkey, *Asymptotic Formulae in Spectral Geometry*:

Interval example generalizes to:

! index theorem for DeRham complex of manifold with boundary.

? index interpretation of Roth–Wilson formula.

$H_{D,N} = -\frac{d^2}{dx^2}$  on  $(0, L)$  with Dirichlet/Neumann BC.

$$\dim \ker H_D = 0, \quad \dim \ker H_N = 1.$$

$$\int_0^L K_{D,N}(t, x, x) dx = \frac{L}{\sqrt{4\pi t}} \mp \frac{1}{2}.$$

Therefore,

$$\dim \ker H_N - \dim \ker H_D = 1 = \text{Tr } K_N - \text{Tr } K_D. \quad (*)$$

Better:  $\begin{cases} d: \text{functions (N)} \rightarrow \text{one-forms (D)}, \\ d^\dagger: \text{one-forms} \rightarrow \text{functions.} \end{cases}$

$$d + d^\dagger = \begin{pmatrix} 0 & d^\dagger \\ d & 0 \end{pmatrix}: \Lambda^0(0, L) \oplus \Lambda^1(0, L) \rightarrow \text{self.}$$

$$H_N = d^\dagger d, \quad H_D \simeq \bar{d} d^\dagger.$$

$$\ker H_N = \ker d, \quad \ker H_D = \ker d^\dagger.$$

Therefore,  $(*) = \text{index } d!$

## EXTENSION TO GRAPHS

Kuchment (2004) “Quantum Graphs: I”:

Self-adjoint BC on  $F(v) = \{f_j(v)\}$  and  $F'(v)$ :

- Dirichlet part:  $P_v F = 0$ ,  $P_v = \text{OG projector}$ ;
- Neumann part:  $(I - P_v)F' = -L_v F$ . ( $L_v = 0$ ).

Form domain:  $\{f \in H^1(e): \text{Dirichlet}\}$ .

Operator domain:  $\{f \in H^2(e): \text{both}\}$ .

These provide natural “large” and “small” domains for  $\frac{d}{dx}$  !

Let  $d = \frac{d}{dx}$  with Kirchhoff conditions —

- Dirichlet part:  $f_j(v)$  independent of  $j$ ;
- Neumann part:  $\sum_j f'_j(v) = 0$ .

Then  $d^\dagger = -\frac{d}{dx}$  with “dual” conditions —

- Dirichlet part:  $\sum_j f_j(v) = 0$ ;
- Neumann part:  $f'_j(v)$  independent of  $j$ .

$$d^\dagger d = \text{Kirchhoff Laplacian}, \quad dd^\dagger = \text{dual Laplacian}$$

(outer operator on form domain, inner operator on operator domain).

Using one-forms removes ambiguity in direction of  $\frac{d}{dx}$  in dual case! Using functions we would need

$$\sum \epsilon_j f_j(v) = 0, \quad \epsilon_j f'_j(v) \text{ indep. of } j$$

for some arbitrary orientation of each edge,

$$\epsilon_j \equiv \begin{cases} +1 & \text{if } v = 0, \\ -1 & \text{if } v = L. \end{cases}$$

For dual Laplacian,  $\frac{1}{4} \sum_{\alpha} S_{\bar{\alpha}\alpha} = -\frac{1}{2}(V - E)$ .  
Therefore,

**Index theorem:**

$$\text{index } d = \dim \ker H_K - \dim \ker H_{\text{dual}} = V - E.$$

So what? Well,

$$\begin{aligned} \dim \ker H_K &= \text{number of components} \\ &= 1 \quad \text{for connected graph,} \end{aligned}$$

so

$$\begin{aligned} \dim \ker H_{\text{dual}} &= E - V + \text{number of components} \\ &= E - V + 1 \quad \text{for any connected graph} \\ &= (\text{e.g.}) \begin{cases} 0 & \text{for a tree (e.g., a star),} \\ 1 & \text{for a cycle.} \end{cases} \end{aligned}$$

Open problem: Is this interesting?