

# Reduction of quantum graphs to tight-binding Hamiltonians

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Quantum graphs are “metrizations” of combinatorial graphs, each edge is replaced by a segment. Can one relate the spectral properties of a quantum graph with its combinatorial counterpart? Natural assumption: all edges are identical.

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More general boundary conditions, more general edges

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essentially  $\text{spec } H = \{z : \cos \sqrt{z} \in \text{spec } \Delta\}$ .

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Fourth order operators (networks of beams, elasticity theory), Nicaise, Dekoninck'00 et al.

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Enumeration of edges near each vertex  $v$ : the edges 1 to  $\text{outdeg } v$  are outgoing w.r.t.  $v$ , the edges  $\text{outdeg } + 1$  to  $\text{deg } v$  are ingoing. Notation:  $e_j v$  is the edge with index  $j$  near  $v$ . Clear,  $\iota e_j v = v$  for  $j \leq \text{outdeg } v$  and  $\tau e_j v = v$  for  $j > \text{outdeg } v$ .

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For each edge  $e$ :  $n_\iota(e)$  and  $n_\tau(e)$  are the indices of  $e$  near  $\iota e$  and  $\tau e$ , respectively. Clear,  $n_\iota(e_j v) = j$  for  $j \leq \text{outdeg } v$  and  $n_\tau(e_j v) = j$  for  $j > \text{outdeg } v$ .

# Metrization of edge

Let  $S_0$  be a closed densely defined symmetric operator in  $\mathcal{H}$  with deficiency indices  $(2, 2)$ . Then for  $S := S_0^*$  one can find two maps

$\Gamma = (\Gamma_0, \Gamma_1), \Gamma' = (\Gamma'_0, \Gamma'_1) : \text{dom} S \rightarrow \mathbb{C}^2$  such that

- $\langle f, Sg \rangle - \langle Sf, g \rangle = \langle \Gamma f, \Gamma' g \rangle - \langle \Gamma' f, \Gamma g \rangle$  for all  $f, g$ ,
- $(\Gamma, \Gamma') : \text{dom} S \rightarrow \mathbb{C}^2 \oplus \mathbb{C}^2$  is surjective: i.e.  $(\mathbb{C}^2, \Gamma, \Gamma')$  is a boundary triple for  $S$ .

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Denote  $H_0 := S|_{\ker \Gamma}$ , self-adjoint. For  $z$  outside of  $\text{spec } H_0$  and any  $\varphi \in \mathbb{C}^2$  one can find a unique  $f$  solving  $(S - z)f = 0$  and  $\Gamma f = \xi$ . Consider the matrix  $m$  defined by  $\Gamma' f = m(z)\Gamma f$ .

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Typical example:  $S = -\partial^2 + V$ ,  $V \in L^2[0, 1]$ , in  $L^2[0, 1]$  with  $\text{dom } S = H^2[0, 1]$ . One can take  $\Gamma f = (f(0), f(1))$ ,  $\Gamma' f = (f'(0), -f'(1))$ ,

$$m(z) = \frac{1}{s(1; z)} \begin{pmatrix} c(1; z) & -1 \\ -1 & s'(1; z) \end{pmatrix},$$

where  $s, c$  solve  $-y'' + Vy = zy$  with  $s(0; z) = c'(0; z) = 1 - s'(0; z) = 1 - c(0; z) = 0$ .

# Quantization

Replace each edge of the graph by a copy of  $S$ ,  $\Gamma_0, \Gamma'_0$  are “glued” to  $\iota e$ ,  $\Gamma_1, \Gamma'_1$  are glued to  $\tau e$ .

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Self-adjointness:  $f_e \in \text{dom } S + \text{boundary conditions}$ . The most general local boundary conditions:  $(1 - U(v))f(v) = i(1 + U(v))f'(v)$  for all  $v$ . Here  $U(v)$  is a unitary matrix of size  $\text{deg } v$ ,

$$f(v) := (\Gamma_0 f_{e_1 v}, \dots, \Gamma_0 f_{e_{\text{outdeg } v} v}, \Gamma_1 f_{e_{\text{outdeg } v+1} v}, \dots, \Gamma_1 f_{e_{\text{deg } v} v}),$$
$$f'(v) := (\Gamma'_0 f_{e_1 v}, \dots, \Gamma'_0 f_{e_{\text{outdeg } v} v}, \Gamma'_1 f_{e_{\text{outdeg } v+1} v}, \dots, \Gamma'_1 f_{e_{\text{deg } v} v}),$$

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Denote  $N(v) := \text{deg } v - \dim \ker(1 + U(v))$ .

Orthonormal basis of eigenvectors  $\xi_j(v)$ ,  $j = 1, \dots, N(v)$ , in  $\ker(1 + U(v))^\perp$ .

Discrete space  $\mathcal{L}$ , “reduced  $l^2$  space”,  $(\varphi(v))$ ,  $\varphi \in \mathbb{C}^{N(v)}$ ,  $\sum \|\varphi(v)\|^2 < \infty$

# Basic result

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If  $\bigcup_{v \in V} \text{spec } U(v) \setminus \{-1\} = \{e^{i\Phi}\}$ , then outside the spectrum of  $H_0$  one has  $\text{spec}_\bullet H = \eta^{-1}(\text{spec}_\bullet P)$ , where  $\bullet \in \{\text{p, pp, disc, ess, ac, sc}\}$ ,

$\eta(z) := \frac{m_{11}(z) + \alpha}{m_{12}(z)}$ ,  $\alpha := \frac{1 - e^{i\Phi}}{i(1 + e^{i\Phi})}$ , and  $P$  is an operator acting on  $\mathcal{L}$  as

$$\begin{aligned} (P\varphi)_j(v) = & \sum_{k=1}^{\text{outdeg } v} \sum_{s=1}^{N(\tau e_k v)} \xi_{jk}(v) \xi_{sn_\tau(e_k v)}(\tau e_k v) \varphi_s(\tau e_k v) \\ & + \sum_{k=\text{outdeg } v+1}^{\text{deg } v} \sum_{s=1}^{N(\iota e_k v)} \xi_{jk}(v) \xi_{sn_\tau(e_k v)}(\iota e_k v) \varphi_s(\iota e_k v). \end{aligned}$$

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$$\bullet \in \{p, pp, \text{disc}, \text{ess}, \text{ac}, \text{sc}\}, \eta(z) := c(1; z) + \alpha s(1; z), \alpha := \frac{1 - e^{i\Phi}}{i(1 + e^{i\Phi})},$$

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Note that  $\eta$  is the Hill discriminant for  $L := -\partial^2 + V + \alpha \sum_{n \in \mathbb{Z}} \delta(\cdot - n)$  and that  $\|P\| \leq 1$ . This provides (up to the Dirichlet levels) the inclusion  $\text{spec } H \subset \text{spec } L$ .

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$L$  has infinitely many gaps  $\Rightarrow$  the same for  $H$  (independent of the underlying graph!)

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 $\sum_{\iota e=v} f'_e(0) - \sum_{\tau e=v} f'_e(1) = \alpha(v)f(v)$ ,  $\alpha(v)$  real. One has

$$U(v) = -E_{\deg v} + \frac{2}{\deg v + \alpha(v)} J_{\deg v}$$

( $J_{\deg v}$  is the matrix whose all entries are 1. For  $\alpha(v) = \alpha \deg v$  reduces to the transition operator  $Pf(v) = D\Delta D$  in  $l^2(V)$ . Here  $D$  is the multiplication by  $1/\sqrt{\deg v}$  and  $\Delta f(v) = \sum_{w \sim v} f(w)$ ).

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**Magnetic case:**  $(i\delta + a_e)^2 + V$  on each edge,  $f$  is continuous at each vertex,  $\sum_{\iota e=v} (\partial - ia_e)f'_e(0) - \sum_{\tau e=v} (\partial - ia_e)f'_e(1) = \alpha(v)f(v)$ ,  $\alpha(v)$  real. By the transformation  $f_e(t) = \exp(\int_0^t a_e(s)ds)g_e(t)$  reduction to the considered situation. Again, for  $\alpha(v) = \alpha \deg v$  reduction to a discrete operator (magnetic Laplacian)  $Pf(v) = D\Delta D$  in  $l^2(V)$  with  $\Delta f(v) = \sum_{\iota e=v} e^{-i\beta(e)} f(\tau e) + \sum_{\tau e=v} e^{i\beta(e)} f(\iota e)$ ,  $\beta(e) := \int_0^1 a_e(s)ds$ , on  $l^2(V)$ .

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$f'_e(0) - f'_{e'}(0) = \frac{\alpha}{\deg v} (f_e(0) - f_{e'}(0))$ . One has

$$U(v) = \frac{\deg v - i\alpha(v)}{\deg v + i\alpha(v)} E_{\deg v} - \frac{2}{\deg v + \alpha(v)} J_{\deg v}.$$

For  $\alpha(v) = \alpha \deg v$  reduction to a discrete operator. Each vertex  $v$  has “degree of freedom”  $\deg v - 1$ .

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$\delta'$  and  $\delta'_s$ -type boundary conditions can be considered. One has to change the boundary triple for  $S$ :  $\Gamma f = (-f'(0), f'(1))$ ,  $\Gamma' f = (f(0), f(1))$ . The function  $\eta$  becomes  $c(1; z) + \alpha c'(1; z)$ , i.e. the Hill discriminant for  $-\partial^2 + V + \alpha \sum_{n \in \mathbb{Z}} \delta'(\cdot - n)$ . The same estimates for the gaps as above, and the Dirichlet spectrum is replaced by the Neumann spectrum.

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Further generalizations (for non-even  $V$ ): bipartite graphs, uniformly directed graphs etc.

# Hybrid spaces

Still  $S_0 = -\Delta$  on a 2D sphere  $M$  with the domain  $\{f \in \text{dom}\Delta : f(a_0) = f(a_1) = 0\}$  ( $a_0, a_1 \in M$  are fixed). Take  $S := S_0^*$ .

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For any  $f \in \text{dom}S$  one can define  $\Gamma_j f = -\lim_{x \rightarrow a_j} \frac{2\pi}{\log d(x, a_j)} f(x)$  and

$\Gamma'_j f = \lim_{x \rightarrow a_j} \left( f(x) + \frac{\Gamma_j f}{2\pi} \log d(x, a_j) \right)$ . It is known that  $(\mathbb{C}^2, \Gamma, \Gamma')$  is a boundary triple for  $S$ . The corresponding matrix  $m$  is

$$m(z) = \begin{pmatrix} G^r(a_0, a_0; z) & G(a_0, a_1; z) \\ G(a_1, a_0; z) & G^r(a_1, a_1; z) \end{pmatrix},$$

where  $G$  is the Green function and  $G^r(x, x; z) := \lim_{y \rightarrow x} \left( G(x, y; z) + \frac{1}{2\pi} \log d(x, y) \right)$ .

# Hybrid spaces

Still  $S_0 = -\Delta$  on a 2D sphere  $M$  with the domain  $\{f \in \text{dom}\Delta : f(a_0) = f(a_1) = 0\}$  ( $a_0, a_1 \in M$  are fixed). Take  $S := S_0^*$ .

For any  $f \in \text{dom}S$  one can define  $\Gamma_j f = -\lim_{x \rightarrow a_j} \frac{2\pi}{\log d(x, a_j)} f(x)$  and

$\Gamma'_j f = \lim_{x \rightarrow a_j} \left( f(x) + \frac{\Gamma_j f}{2\pi} \log d(x, a_j) \right)$ . It is known that  $(\mathbb{C}^2, \Gamma, \Gamma')$  is a boundary triple for  $S$ . The corresponding matrix  $m$  is

$$m(z) = \begin{pmatrix} G^r(a_0, a_0; z) & G(a_0, a_1; z) \\ G(a_1, a_0; z) & G^r(a_1, a_1; z) \end{pmatrix},$$

where  $G$  is the Green function and  $G^r(x, x; z) := \lim_{y \rightarrow x} \left( G(x, y; z) + \frac{1}{2\pi} \log d(x, y) \right)$ .

Replace each edge by a copy of  $S$ . This results is a hybrid space (coupled spheres).

The study is as above, the function  $\eta(z) := \frac{m_{11}(z) + \alpha}{m_{12}(z)}$  has poles, but the same discrete Hamiltonian!