

Conformal Field Theory and Combinatorics

Part III: Applications to Combinatorics

Jesper L. Jacobsen^{1,2}

¹Université Pierre et Marie Curie, Paris 6, France

²Institut de Physique Théorique, CEA/Saclay, France

Friday 18 January, 2008

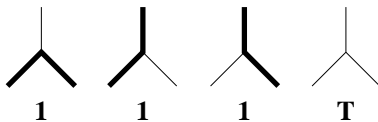
Summary

- 1 **Hamiltonian circuits and walks**
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 **Coulomb Gas construction**
 - Compactification with respect to a lattice
 - Liouville field theory
- 3 **Critical exponents**
 - Central charge and watermelon exponents
 - Standard exponents of polymer physics

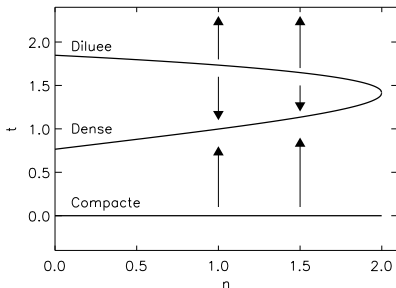
Summary

- 1 **Hamiltonian circuits and walks**
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 Coulomb Gas construction
 - Compactification with respect to a lattice
 - Liouville field theory
- 3 Critical exponents
 - Central charge and watermelon exponents
 - Standard exponents of polymer physics

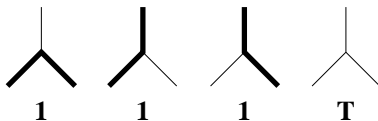
$O(n)$ MODEL ON THE HEXAGONAL LATTICE



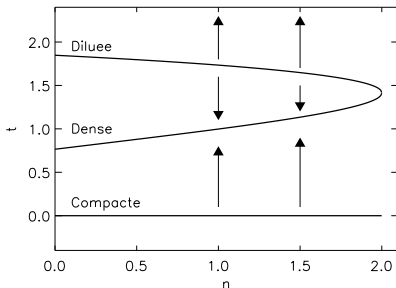
- Low- T branch is RG attractive and $Z = Z(T^2)$
- Hence $T = 0$ is a line of fixed points



$O(n)$ MODEL ON THE HEXAGONAL LATTICE



- Low- T branch is RG attractive and $Z = Z(T^2)$
- Hence $T = 0$ is a line of fixed points



RIGHT AT $T = 0$ DEFINE PARTITION FUNCTION

$$Z = \sum_g n^N$$

- $n \rightarrow 0$ describes Hamiltonian circuits
- Hamiltonian walks obtained by inserting defect pair

MAPPING TO $D = 2$ DIMENSIONAL HEIGHT MODEL

- With $T = 0$, height differences satisfy $\mathbf{A} + \mathbf{B} + \mathbf{C} = \mathbf{0}$

$$\mathbf{A} = \left(\frac{1}{\sqrt{3}}, 0 \right), \quad \mathbf{B} = \left(-\frac{1}{2\sqrt{3}}, \frac{1}{2} \right), \quad \mathbf{C} = \left(-\frac{1}{2\sqrt{3}}, -\frac{1}{2} \right)$$

RIGHT AT $T = 0$ DEFINE PARTITION FUNCTION

$$Z = \sum_{\mathcal{G}} n^N$$

- $n \rightarrow 0$ describes Hamiltonian circuits
- Hamiltonian walks obtained by inserting defect pair

MAPPING TO $D = 2$ DIMENSIONAL HEIGHT MODEL

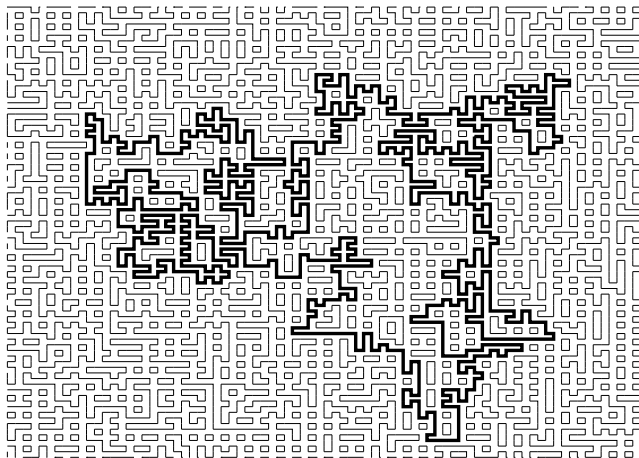
- With $T = 0$, height differences satisfy $\mathbf{A} + \mathbf{B} + \mathbf{C} = \mathbf{0}$

$$\mathbf{A} = \left(\frac{1}{\sqrt{3}}, 0 \right), \quad \mathbf{B} = \left(-\frac{1}{2\sqrt{3}}, \frac{1}{2} \right), \quad \mathbf{C} = \left(-\frac{1}{2\sqrt{3}}, -\frac{1}{2} \right)$$

Summary

- 1 **Hamiltonian circuits and walks**
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 Coulomb Gas construction
 - Compactification with respect to a lattice
 - Liouville field theory
- 3 Critical exponents
 - Central charge and watermelon exponents
 - Standard exponents of polymer physics

FPL² = MODEL OF TWO TYPES OF FULLY-PACKED LOOPS



PARTITION FUNCTION OF FPL² MODEL

$$Z = \sum_{\mathcal{G}} n_b^{N_b} n_g^{N_g}$$

- Here N_b (resp. N_g) is # black (resp. green) loops

MAPPING TO $D = 3$ DIMENSIONAL HEIGHT MODEL

$$\mathbf{A} + \mathbf{B} + \mathbf{C} + \mathbf{D} = \mathbf{0}$$

$$\mathbf{A} = (-1, 1, 1) \quad \mathbf{B} = (1, 1, -1)$$

$$\mathbf{C} = (-1, -1, -1) \quad \mathbf{D} = (1, -1, 1)$$

PARTITION FUNCTION OF FPL² MODEL

$$Z = \sum_{\mathcal{G}} n_b^{N_b} n_g^{N_g}$$

- Here N_b (resp. N_g) is # black (resp. green) loops

MAPPING TO $D = 3$ DIMENSIONAL HEIGHT MODEL

$$\mathbf{A} + \mathbf{B} + \mathbf{C} + \mathbf{D} = \mathbf{0}$$

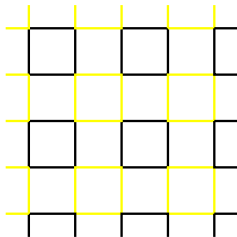
$$\begin{array}{ll} \mathbf{A} = (-1, 1, 1) & \mathbf{B} = (1, 1, -1) \\ \mathbf{C} = (-1, -1, -1) & \mathbf{D} = (1, -1, 1) \end{array}$$

Summary

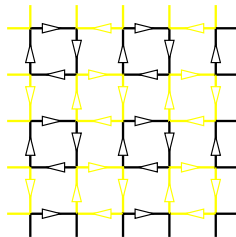
- 1 Hamiltonian circuits and walks
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 **Coulomb Gas construction**
 - **Compactification with respect to a lattice**
 - Liouville field theory
- 3 Critical exponents
 - Central charge and watermelon exponents
 - Standard exponents of polymer physics

Ideal states = MACROSCOPICALLY FLAT STATES

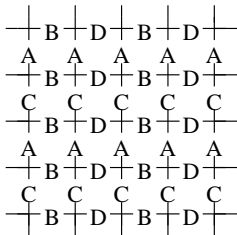
a)



b)



c)

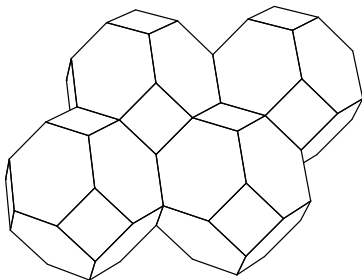


d)

	0	A	0	A
	-B	A+D	-B	A+D
	0	A	0	A
	-B	A+D	-B	A+D

CONSTRUCTION OF *ideal state graph* \mathcal{I}

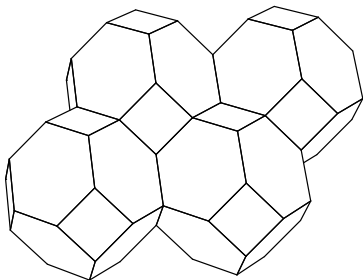
- Transition between ideal states: transpose colour pair
- Associated change of average height $\in \mathbb{R}^D$
- In \mathcal{I} , vertex = ideal state, edge = transition



- A fixed ideal state occurs ∞ many times in \mathcal{I}
- Defines *repeat lattice* \mathcal{R}
- Here FCC lattice with cubic cell of edge length 4

CONSTRUCTION OF *ideal state graph* \mathcal{I}

- Transition between ideal states: transpose colour pair
- Associated change of average height $\in \mathbb{R}^D$
- In \mathcal{I} , vertex = ideal state, edge = transition



- A fixed ideal state occurs ∞ many times in \mathcal{I}
- Defines *repeat lattice* \mathcal{R}
- Here FCC lattice with cubic cell of edge length 4

Summary

- 1 Hamiltonian circuits and walks
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 **Coulomb Gas construction**
 - Compactification with respect to a lattice
 - **Liouville field theory**
- 3 Critical exponents
 - Central charge and watermelon exponents
 - Standard exponents of polymer physics

COMPACTIFIED HEIGHT MODEL

- $\mathbf{h}(\mathbf{x}) \in \mathbb{R}^D/\mathcal{R}$ with EM charges $\mathbf{e} \in \mathcal{R}^*$ and $\mathbf{m} \in \mathcal{R}$

ELASTIC TERM S_E

- Constrained by rotational invariance...

$$S_E = \frac{1}{2} \int d^2\mathbf{x} K_{\alpha\beta} \partial h^\alpha \cdot \partial h^\beta$$

- ...and by loop reversal symmetries ($\mathbf{A} \leftrightarrow \mathbf{B}$ and $\mathbf{C} \leftrightarrow \mathbf{D}$)

$$S_E = \frac{1}{2} \int d^2\mathbf{x} \left\{ K_{11} [(\partial h^1)^2 + (\partial h^3)^2] + 2K_{13} (\partial h^1 \cdot \partial h^3) + K_{22} (\partial h^2)^2 \right\}$$

- Diagonalised by a change of basis for \mathbf{h}

COMPACTIFIED HEIGHT MODEL

- $\mathbf{h}(\mathbf{x}) \in \mathbb{R}^D/\mathcal{R}$ with EM charges $\mathbf{e} \in \mathcal{R}^*$ and $\mathbf{m} \in \mathcal{R}$

ELASTIC TERM S_E

- Constrained by rotational invariance...

$$S_E = \frac{1}{2} \int d^2\mathbf{x} K_{\alpha\beta} \partial h^\alpha \cdot \partial h^\beta$$

- ...and by loop reversal symmetries ($\mathbf{A} \leftrightarrow \mathbf{B}$ and $\mathbf{C} \leftrightarrow \mathbf{D}$)

$$S_E = \frac{1}{2} \int d^2\mathbf{x} \left\{ K_{11} [(\partial h^1)^2 + (\partial h^3)^2] + 2K_{13} (\partial h^1 \cdot \partial h^3) + K_{22} (\partial h^2)^2 \right\}$$

- Diagonalised by a change of basis for \mathbf{h}

COMPACTIFIED HEIGHT MODEL

- $\mathbf{h}(\mathbf{x}) \in \mathbb{R}^D/\mathcal{R}$ with EM charges $\mathbf{e} \in \mathcal{R}^*$ and $\mathbf{m} \in \mathcal{R}$

ELASTIC TERM S_E

- Constrained by rotational invariance...

$$S_E = \frac{1}{2} \int d^2\mathbf{x} K_{\alpha\beta} \partial h^\alpha \cdot \partial h^\beta$$

- ...and by loop reversal symmetries ($\mathbf{A} \leftrightarrow \mathbf{B}$ and $\mathbf{C} \leftrightarrow \mathbf{D}$)

$$S_E = \frac{1}{2} \int d^2\mathbf{x} \left\{ K_{11} [(\partial h^1)^2 + (\partial h^3)^2] + 2K_{13} (\partial h^1 \cdot \partial h^3) + K_{22} (\partial h^2)^2 \right\}$$

- Diagonalised by a change of basis for \mathbf{h}

BOUNDARY TERM S_B

$$S_B = \frac{i}{4\pi} \int d^2\mathbf{x} (\mathbf{e}_0 \cdot \mathbf{h}) \mathcal{R}(\mathbf{x})$$

- $\mathcal{R}(\mathbf{x})$ is the scalar curvature
- Parametrise weights: $n_b = 2 \cos(\pi e_b)$ and $n_g = 2 \cos(\pi e_g)$
- Constraints:
 $\mathbf{e}_0 \cdot \mathbf{A} = \pi e_b$, $\mathbf{e}_0 \cdot \mathbf{B} = -\pi e_b$, $\mathbf{e}_0 \cdot \mathbf{C} = \pi e_g$, $\mathbf{e}_0 \cdot \mathbf{D} = -\pi e_g$
- Solution: $\mathbf{e}_0 = -\pi(e_b, 0, e_g)$

BOUNDARY TERM S_B

$$S_B = \frac{i}{4\pi} \int d^2\mathbf{x} (\mathbf{e}_0 \cdot \mathbf{h}) \mathcal{R}(\mathbf{x})$$

- $\mathcal{R}(\mathbf{x})$ is the scalar curvature
- Parametrise weights: $n_b = 2 \cos(\pi \mathbf{e}_b)$ and $n_g = 2 \cos(\pi \mathbf{e}_g)$
- Constraints:
 $\mathbf{e}_0 \cdot \mathbf{A} = \pi \mathbf{e}_b$, $\mathbf{e}_0 \cdot \mathbf{B} = -\pi \mathbf{e}_b$, $\mathbf{e}_0 \cdot \mathbf{C} = \pi \mathbf{e}_g$, $\mathbf{e}_0 \cdot \mathbf{D} = -\pi \mathbf{e}_g$
- Solution: $\mathbf{e}_0 = -\pi(\mathbf{e}_b, 0, \mathbf{e}_g)$

LIOUVILLE TERM S_L

$$S_L = \int d^2\mathbf{x} w[\mathbf{h}(\mathbf{x})]$$

- $\exp(-w[\mathbf{h}(\mathbf{x})])$ is scaling limit of vertex weights in oriented loop representation
- Let $(\sigma_1, \sigma_2, \sigma_3, \sigma_4)$ be colours around a vertex. Then $w(\mathbf{h}) = \frac{i}{16} \mathbf{e}_0 \cdot [(\sigma_1 - \sigma_3) \times (\sigma_2 - \sigma_4)]$

FOURIER ANALYSIS

$$w(\mathbf{h}) = \sum_{\mathbf{e} \in \mathcal{R}_w^*} \tilde{w}_{\mathbf{e}} \exp(i\mathbf{e} \cdot \mathbf{h}) .$$

- \mathcal{R}_w^* spanned by next-shortest vectors in \mathcal{R}^*

LIOUVILLE TERM S_L

$$S_L = \int d^2\mathbf{x} w[\mathbf{h}(\mathbf{x})]$$

- $\exp(-w[\mathbf{h}(\mathbf{x})])$ is scaling limit of vertex weights in oriented loop representation
- Let $(\sigma_1, \sigma_2, \sigma_3, \sigma_4)$ be colours around a vertex. Then $w(\mathbf{h}) = \frac{i}{16} \mathbf{e}_0 \cdot [(\sigma_1 - \sigma_3) \times (\sigma_2 - \sigma_4)]$

FOURIER ANALYSIS

$$w(\mathbf{h}) = \sum_{\mathbf{e} \in \mathcal{R}_w^*} \tilde{w}_{\mathbf{e}} \exp(i\mathbf{e} \cdot \mathbf{h}) .$$

- \mathcal{R}_w^* spanned by next-shortest vectors in \mathcal{R}^*

DIMENSIONS OF ELECTROMAGNETIC OPERATORS

- In basis where coupling constant tensor g^α is diagonal:

$$\Delta_{\mathbf{e},\mathbf{m}} = \frac{1}{4\pi} \left[\frac{1}{g_\alpha} \mathbf{e}_\alpha (\mathbf{e}_\alpha - 2\mathbf{e}_{0\alpha}) + g_\alpha (m^\alpha)^2 \right]$$

MARGINALITY REQUIREMENT

- Most relevant \mathbf{e} -charges in \mathcal{R}_w^* are marginal

$$g_1 = \frac{\pi}{2}(1 - e_b), \quad \frac{1}{g_2} = \frac{1}{g_1} + \frac{1}{g_3}, \quad g_3 = \frac{\pi}{2}(1 - e_g)$$

DIMENSIONS OF ELECTROMAGNETIC OPERATORS

- In basis where coupling constant tensor g^α is diagonal:

$$\Delta_{\mathbf{e},\mathbf{m}} = \frac{1}{4\pi} \left[\frac{1}{g_\alpha} \mathbf{e}_\alpha (\mathbf{e}_\alpha - 2\mathbf{e}_{0\alpha}) + g_\alpha (m^\alpha)^2 \right]$$

MARGINALITY REQUIREMENT

- Most relevant \mathbf{e} -charges in \mathcal{R}_w^* are marginal

$$g_1 = \frac{\pi}{2}(1 - \mathbf{e}_b), \quad \frac{1}{g_2} = \frac{1}{g_1} + \frac{1}{g_3}, \quad g_3 = \frac{\pi}{2}(1 - \mathbf{e}_g)$$

Summary

- 1 Hamiltonian circuits and walks
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 Coulomb Gas construction
 - Compactification with respect to a lattice
 - Liouville field theory
- 3 **Critical exponents**
 - **Central charge and watermelon exponents**
 - Standard exponents of polymer physics

$$\text{CENTRAL CHARGE } c = D + 12\Delta_{\mathbf{e}_0, \mathbf{0}}$$

$$c = \begin{cases} 2 - \frac{6e_0^2}{1-e_0} & \text{for the FPL model} \\ 3 - 6 \left(\frac{e_b^2}{1-e_b} + \frac{e_g^2}{1-e_g} \right) & \text{for the FPL}^2 \text{ model} \end{cases}$$

WATERMELON EXPONENTS (FROM VECTOR DEFECTS $\mathbf{m} \in \mathbb{R}^D$)

- For FPL model on honeycomb lattice

$$\Delta_\ell = \begin{cases} \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} & \text{for } \ell \text{ even} \\ \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} + \frac{3g}{8} & \text{for } \ell \text{ odd} \end{cases}$$

- For FPL² model on square lattice (with $n_g = 1$, and $g \equiv g_1$)

$$\Delta_\ell = \begin{cases} \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} & \text{for } \ell \text{ even} \\ \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} + \frac{g}{3g+2} & \text{for } \ell \text{ odd} \end{cases}$$

CENTRAL CHARGE $c = D + 12\Delta_{\mathbf{e}_0, \mathbf{0}}$

$$c = \begin{cases} 2 - \frac{6e_0^2}{1-e_0} & \text{for the FPL model} \\ 3 - 6 \left(\frac{e_b^2}{1-e_b} + \frac{e_g^2}{1-e_g} \right) & \text{for the FPL}^2 \text{ model} \end{cases}$$

WATERMELON EXPONENTS (FROM VECTOR DEFECTS $\mathbf{m} \in \mathbb{R}^D$)

- For FPL model on honeycomb lattice

$$\Delta_\ell = \begin{cases} \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} & \text{for } \ell \text{ even} \\ \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} + \frac{3g}{8} & \text{for } \ell \text{ odd} \end{cases}$$

- For FPL² model on square lattice (with $n_g = 1$, and $g \equiv g_1$)

$$\Delta_\ell = \begin{cases} \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} & \text{for } \ell \text{ even} \\ \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} + \frac{g}{3g+2} & \text{for } \ell \text{ odd} \end{cases}$$

CENTRAL CHARGE $c = D + 12\Delta_{\mathbf{e}_0, \mathbf{0}}$

$$c = \begin{cases} 2 - \frac{6e_0^2}{1-e_0} & \text{for the FPL model} \\ 3 - 6 \left(\frac{e_b^2}{1-e_b} + \frac{e_g^2}{1-e_g} \right) & \text{for the FPL}^2 \text{ model} \end{cases}$$

WATERMELON EXPONENTS (FROM VECTOR DEFECTS $\mathbf{m} \in \mathbb{R}^D$)

- For FPL model on honeycomb lattice

$$\Delta_\ell = \begin{cases} \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} & \text{for } \ell \text{ even} \\ \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} + \frac{3g}{8} & \text{for } \ell \text{ odd} \end{cases}$$

- For FPL² model on square lattice (with $n_g = 1$, and $g \equiv g_1$)

$$\Delta_\ell = \begin{cases} \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} & \text{for } \ell \text{ even} \\ \frac{1}{8}g\ell^2 - \frac{(1-g)^2}{2g} + \frac{g}{3g+2} & \text{for } \ell \text{ odd} \end{cases}$$

Summary

- 1 Hamiltonian circuits and walks
 - $O(n)$ model revisited
 - FPL² model on the square lattice
- 2 Coulomb Gas construction
 - Compactification with respect to a lattice
 - Liouville field theory
- 3 **Critical exponents**
 - Central charge and watermelon exponents
 - **Standard exponents of polymer physics**

NUMBER OF SELF-AVOIDING WALKS AND POLYGONS

- Length l , with one monomer attached to a fixed point

$$\mathcal{N}_{\text{SAW}} \sim \mu^l l^{\gamma-1} \quad \mathcal{N}_{\text{SAP}} \sim \mu^l l^{-\nu d}$$

- Standard scaling relations

$$\Delta_1 = 1 - \frac{\gamma}{2\nu} \quad \Delta_2 = 2 - \frac{1}{\nu}$$

RESULTS FOR CONFORMATIONAL EXPONENTS

Model	ν	γ
Dilute $O(n \rightarrow 0)$	$\frac{3}{4}$	$\frac{43}{32}$
Dense $O(n \rightarrow 0)$	$\frac{1}{2}$	$\frac{19}{16}$
FPL honeycomb	$\frac{1}{2}$	1
FPL square	$\frac{1}{2}$	$\frac{117}{112}$

NUMBER OF SELF-AVOIDING WALKS AND POLYGONS

- Length l , with one monomer attached to a fixed point

$$\mathcal{N}_{\text{SAW}} \sim \mu^l l^{\gamma-1} \quad \mathcal{N}_{\text{SAP}} \sim \mu^l l^{-\nu d}$$

- Standard scaling relations

$$\Delta_1 = 1 - \frac{\gamma}{2\nu} \quad \Delta_2 = 2 - \frac{1}{\nu}$$

RESULTS FOR CONFORMATIONAL EXPONENTS

Model	ν	γ
Dilute $O(n \rightarrow 0)$	$\frac{3}{4}$	$\frac{43}{32}$
Dense $O(n \rightarrow 0)$	$\frac{1}{2}$	$\frac{19}{16}$
FPL honeycomb	$\frac{1}{2}$	1
FPL square	$\frac{1}{2}$	$\frac{117}{112}$