

Ising Model on K_n : Mixing time for Glauber dynamics at critical β

David A. Levin, *U. Oregon*
Malwina Luczak, *London School of Economics*
Yuval Peres, *Berkeley and Microsoft Research*

Newton Institute, March 25, 2008

Mean-field Ising model

Take $G_n = K_n$, the complete graph on the n vertices: $V_n = \{1, \dots, n\}$, and E_n contains all $\binom{n}{2}$ possible edges.

Ising measure on $\{-1, 1\}^{K_n}$ is

$$\pi(\sigma) = \frac{\exp\left(\frac{\beta}{n} \sum_{1 \leq i \neq j \leq n} \sigma(i)\sigma(j)\right)}{Z(\beta)}.$$

Glauber dynamics: Pick $v \in \{1, 2, \dots, n\}$ uniformly at random, update $\sigma(v)$.

The probability of updating to a +1 is then

$$\frac{e^{\beta(S-\sigma(v)/n)}}{e^{\beta(S-\sigma(v)/n)} + e^{-\beta(S-\sigma(v)/n)}}$$

where S is the *average magnetization*

$$S = \frac{1}{n} \sum_{i=1}^n \sigma(i).$$

Mean-field Ising model

Take $G_n = K_n$, the complete graph on the n vertices: $V_n = \{1, \dots, n\}$, and E_n contains all $\binom{n}{2}$ possible edges.

Ising measure on $\{-1, 1\}^{K_n}$ is

$$\pi(\sigma) = \frac{\exp\left(\frac{\beta}{n} \sum_{1 \leq i \neq j \leq n} \sigma(i)\sigma(j)\right)}{Z(\beta)}.$$

Glauber dynamics: Pick $v \in \{1, 2, \dots, n\}$ uniformly at random, update $\sigma(v)$.

The probability of updating to a +1 is then

$$\frac{e^{\beta(S-\sigma(v)/n)}}{e^{\beta(S-\sigma(v)/n)} + e^{-\beta(S-\sigma(v)/n)}}$$

where S is the *average magnetization*

$$S = \frac{1}{n} \sum_{i=1}^n \sigma(i).$$

Mean-field Ising model

Take $G_n = K_n$, the complete graph on the n vertices: $V_n = \{1, \dots, n\}$, and E_n contains all $\binom{n}{2}$ possible edges.

Ising measure on $\{-1, 1\}^{K_n}$ is

$$\pi(\sigma) = \frac{\exp\left(\frac{\beta}{n} \sum_{1 \leq i \neq j \leq n} \sigma(i)\sigma(j)\right)}{Z(\beta)}.$$

Glauber dynamics: Pick $v \in \{1, 2, \dots, n\}$ uniformly at random, update $\sigma(v)$.

The probability of updating to a +1 is then

$$\frac{e^{\beta(S-\sigma(v)/n)}}{e^{\beta(S-\sigma(v)/n)} + e^{-\beta(S-\sigma(v)/n)}}$$

where S is the *average magnetization*

$$S = \frac{1}{n} \sum_{i=1}^n \sigma(i).$$

Theorem (D.A.L., Luczak, Peres '07)

Let (X_t^n) be the Glauber dynamics for the Ising model on K_n . If $\beta = 1$, then there are constants c_1 and c_2 so that

$$c_1 n^{3/2} \leq t_{\text{mix}} \leq c_2 n^{3/2}.$$

Contrast with (Cf. M. Luczak's talk)

- 1 $n \log n$ mixing at $\beta < 1$
- 2 $n \log n$ mixing at $\beta > 1$ with truncated dynamics,
- 3 Exponentially slow mixing at $\beta > 1$.

Theorem (D.A.L., Luczak, Peres '07)

Let (X_t^n) be the Glauber dynamics for the Ising model on K_n . If $\beta = 1$, then there are constants c_1 and c_2 so that

$$c_1 n^{3/2} \leq t_{\text{mix}} \leq c_2 n^{3/2}.$$

Contrast with (Cf. M. Luczak's talk)

- 1 $n \log n$ mixing at $\beta < 1$
- 2 $n \log n$ mixing at $\beta > 1$ with truncated dynamics,
- 3 Exponentially slow mixing at $\beta > 1$.

- 1 First show that can couple together two copies of the magnetization chain (S_t) so that they agree in $O(n^{3/2})$ steps.
 - Show that one copy of the chain (S_t) hits zero in order $O(n^{3/2})$ steps.
 - Use the *reflection coupling* to show that (S_t) can be coupled to a second (\tilde{S}_t) so that $S_t = \tilde{S}_t$ by the time (S_t) hits zero.
- 2 Then appeal to the general lemma that once the magnetizations agree, can couple together the full chains in order $n \log n$ more steps.

- 1 First show that can couple together two copies of the magnetization chain (S_t) so that they agree in $O(n^{3/2})$ steps.
 - Show that one copy of the chain (S_t) hits zero in order $O(n^{3/2})$ steps.
 - Use the *reflection coupling* to show that (S_t) can be coupled to a second (\tilde{S}_t) so that $S_t = \tilde{S}_t$ by the time (S_t) hits zero.
- 2 Then appeal to the general lemma that once the magnetizations agree, can couple together the full chains in order $n \log n$ more steps.

- 1 First show that can couple together two copies of the magnetization chain (S_t) so that they agree in $O(n^{3/2})$ steps.
 - Show that one copy of the chain (S_t) hits zero in order $O(n^{3/2})$ steps.
 - Use the *reflection coupling* to show that (S_t) can be coupled to a second (\tilde{S}_t) so that $S_t = \tilde{S}_t$ by the time (S_t) hits zero.
- 2 Then appeal to the general lemma that once the magnetizations agree, can couple together the full chains in order $n \log n$ more steps.

- 1 First show that can couple together two copies of the magnetization chain (S_t) so that they agree in $O(n^{3/2})$ steps.
 - Show that one copy of the chain (S_t) hits zero in order $O(n^{3/2})$ steps.
 - Use the *reflection coupling* to show that (S_t) can be coupled to a second (\tilde{S}_t) so that $S_t = \tilde{S}_t$ by the time (S_t) hits zero.
- 2 Then appeal to the general lemma that once the magnetizations agree, can couple together the full chains in order $n \log n$ more steps.

$\beta = 1$. Why $n^{3/2}$?

The time τ_0 for one copy of (S_t) to hit zero is $O(n^{3/2})$.

To prove, two phases: First drive expectation $s_t = E(S_t)$ down, then allow for the fluctuations of (S_t) to finish the job.

Phase 1: (S_t) has drift

$$\mathbb{E}[S_{t+1} - S_t \mid \mathcal{F}_t] \approx - \left[\frac{S_t}{n} - \tanh(\beta S_t/n) \right].$$

Here $\beta = 1$. Expanding $\tanh(x) = x - x^3/3 + \dots$ in the key equation yields

$$\mathbb{E}[S_{t+1} - S_t \mid \mathcal{F}_t] \approx -\frac{1}{3} \left(\frac{S_t}{n} \right)^3.$$

$\beta = 1$. Why $n^{3/2}$?

The time τ_0 for one copy of (S_t) to hit zero is $O(n^{3/2})$.

To prove, two phases: First drive expectation $s_t = E(S_t)$ down, then allow for the fluctuations of (S_t) to finish the job.

Phase 1: (S_t) has drift

$$\mathbb{E}[S_{t+1} - S_t \mid \mathcal{F}_t] \approx - \left[\frac{S_t}{n} - \tanh(\beta S_t/n) \right].$$

Here $\beta = 1$. Expanding $\tanh(x) = x - x^3/3 + \dots$ in the key equation yields

$$\mathbb{E}[S_{t+1} - S_t \mid \mathcal{F}_t] \approx -\frac{1}{3} \left(\frac{S_t}{n} \right)^3.$$

Thus,

$$s_{t+1} - s_t \approx -\frac{s_t^3}{3n}.$$

Let $b_k = 2^{-k}$. If $s_t \leq b_k$, then the number of additional steps u needed for $s_{t+u} \leq b_{k+1}$ is not more than

$$\frac{3n(b_{k+1} - b_k)}{b_k^3} = \frac{c_1 n}{b_k^2}.$$

The total number of steps needed to get below b_{k_0} is then

$$\sum_{k=1}^{k_0} \frac{c_1 n}{b_k^2} = \frac{c_2 n}{b_{k_0}^2}.$$

Taking $b_{k_0} = n^{\alpha-1}$, we need

$$\frac{c_2 n}{n^{2\alpha-2}} = c_2 n^{3-2\alpha}$$

steps.

Thus,

$$s_{t+1} - s_t \approx -\frac{s_t^3}{3n}.$$

Let $b_k = 2^{-k}$. If $s_t \leq b_k$, then the number of additional steps u needed for $s_{t+u} \leq b_{k+1}$ is not more than

$$\frac{3n(b_{k+1} - b_k)}{b_k^3} = \frac{c_1 n}{b_k^2}.$$

The total number of steps needed to get below b_{k_0} is then

$$\sum_{k=1}^{k_0} \frac{c_1 n}{b_k^2} = \frac{c_2 n}{b_{k_0}^2}.$$

Taking $b_{k_0} = n^{\alpha-1}$, we need

$$\frac{c_2 n}{n^{2\alpha-2}} = c_2 n^{3-2\alpha}$$

steps.

Thus,

$$s_{t+1} - s_t \approx -\frac{s_t^3}{3n}.$$

Let $b_k = 2^{-k}$. If $s_t \leq b_k$, then the number of additional steps u needed for $s_{t+u} \leq b_{k+1}$ is not more than

$$\frac{3n(b_{k+1} - b_k)}{b_k^3} = \frac{c_1 n}{b_k^2}.$$

The total number of steps needed to get below b_{k_0} is then

$$\sum_{k=1}^{k_0} \frac{c_1 n}{b_k^2} = \frac{c_2 n}{b_{k_0}^2}.$$

Taking $b_{k_0} = n^{\alpha-1}$, we need

$$\frac{c_2 n}{n^{2\alpha-2}} = c_2 n^{3-2\alpha}$$

steps.

Thus,

$$s_{t+1} - s_t \approx -\frac{s_t^3}{3n}.$$

Let $b_k = 2^{-k}$. If $s_t \leq b_k$, then the number of additional steps u needed for $s_{t+u} \leq b_{k+1}$ is not more than

$$\frac{3n(b_{k+1} - b_k)}{b_k^3} = \frac{c_1 n}{b_k^2}.$$

The total number of steps needed to get below b_{k_0} is then

$$\sum_{k=1}^{k_0} \frac{c_1 n}{b_k^2} = \frac{c_2 n}{b_{k_0}^2}.$$

Taking $b_{k_0} = n^{\alpha-1}$, we need

$$\frac{c_2 n}{n^{2\alpha-2}} = c_2 n^{3-2\alpha}$$

steps.

Phase 2: fluctuations

The fluctuations of nS_t are similar to those of simple unbiased random walk.

Need additional $n^{2\alpha}$ steps to hit zero.

Total time to hit zero (with high probability) is

$$O(n^{3-2\alpha}) + O(n^{2\alpha})$$

The choice $\alpha = 3/4$ makes the powers equal, and gives hitting time $O(n^{3/2})$.

Phase 2: fluctuations

The fluctuations of nS_t are similar to those of simple unbiased random walk.

Need additional $n^{2\alpha}$ steps to hit zero.

Total time to hit zero (with high probability) is

$$O(n^{3-2\alpha}) + O(n^{2\alpha})$$

The choice $\alpha = 3/4$ makes the powers equal, and gives hitting time $O(n^{3/2})$.

Phase 2: fluctuations

The fluctuations of nS_t are similar to those of simple unbiased random walk.

Need additional $n^{2\alpha}$ steps to hit zero.

Total time to hit zero (with high probability) is

$$O(n^{3-2\alpha}) + O(n^{2\alpha})$$

The choice $\alpha = 3/4$ makes the powers equal, and gives hitting time $O(n^{3/2})$.

Phase 2: fluctuations

The fluctuations of nS_t are similar to those of simple unbiased random walk.

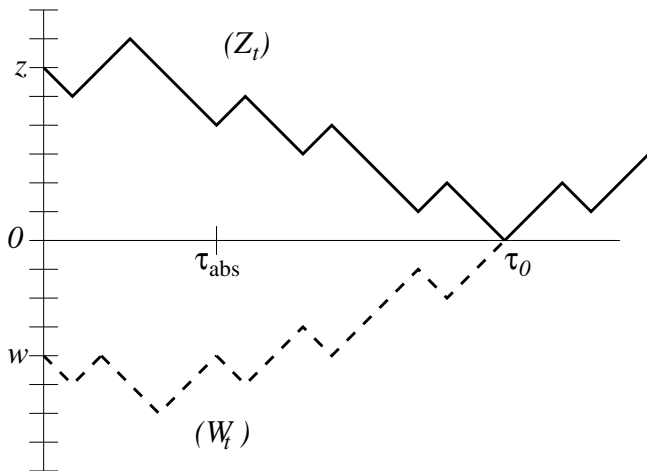
Need additional $n^{2\alpha}$ steps to hit zero.

Total time to hit zero (with high probability) is

$$O(n^{3-2\alpha}) + O(n^{2\alpha})$$

The choice $\alpha = 3/4$ makes the powers equal, and gives hitting time $O(n^{3/2})$.

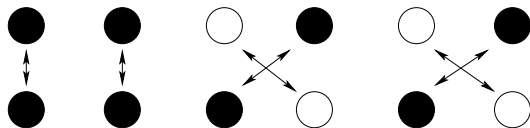
Reflection coupling



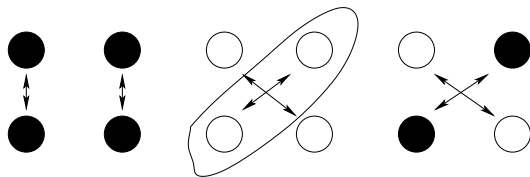
Wait until absolute values meet (they could jump over each other, but that happens with fixed probability), then reflect.

Coupling from same magnetization

Match vertices so that agreements are matched together. Update matched vertices together.

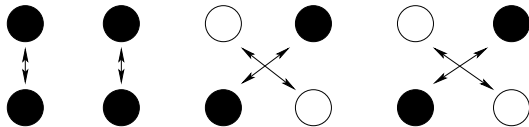


When different vertices are updated together to a different spin, the number of disagreements decreases by two.

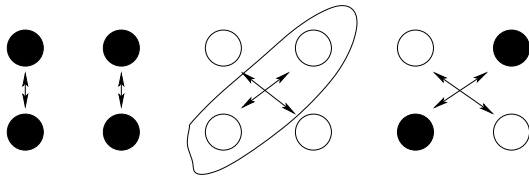


Coupling from same magnetization

Match vertices so that agreements are matched together. Update matched vertices together.



When different vertices are updated together to a different spin, the number of disagreements decreases by two.



$$E[D_{t+1} - D_t \mid \mathcal{F}_t] \leq \frac{-c}{n} D_t,$$

and so

$$E[D_t] \leq \left(1 - \frac{c}{n}\right)^t \leq e^{-ct/n}.$$

Taking $t = \gamma n \log n$ for γ large makes this small.

Under stationary measure,

$n^{1/4}S_n$ converges weakly to non-trivial distribution.

So,

$$P(S_n > An^{-1/4}) \rightarrow 0 \quad \text{as } A \rightarrow \infty.$$

(Simon and Griffiths 1973)

If $t = \gamma n^{3/2}$, then $E(S_t) \approx c(\gamma)n^{1/4}$, where $c(\gamma) \rightarrow \infty$ as $\gamma \rightarrow 0$.

This if $t = \gamma n^{3/2}$, then the distribution $P(S_t \in \cdot)$ and $\mu(S \in \cdot)$ are separated.

Under stationary measure,

$n^{1/4}S_n$ converges weakly to non-trivial distribution.

So,

$$P(S_n > An^{-1/4}) \rightarrow 0 \quad \text{as } A \rightarrow \infty.$$

(Simon and Griffiths 1973)

If $t = \gamma n^{3/2}$, then $E(S_t) \approx c(\gamma)n^{1/4}$, where $c(\gamma) \rightarrow \infty$ as $\gamma \rightarrow 0$.

This if $t = \gamma n^{3/2}$, then the distribution $P(S_t \in \cdot)$ and $\mu(S \in \cdot)$ are separated.

Consider the d -dimensional torus $(\mathbb{Z}/n\mathbb{Z})^d$. Let β_c be the critical temperature.

For $\beta = \beta_c$, the mixing time is polynomial in n .

Stronger: $t_{\text{mix}} = O(|V_n|^{3/2})$ for $d > d_c$.