

Non-geometric fluxes in higher dimensions: I

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arXiv:1106.4015 by D. A., M. Larfors, D. Lüst, P. Patalong
arXiv:1202.3060 by D. A., O. Hohm, M. Larfors, D. Lüst, P. Patalong
arXiv:1204.1979 by D. A., O. Hohm, M. Larfors, D. Lüst, P. Patalong
and work in progress...

Relations to many recent papers...

26/06/2012, String Phenomenology 2012,
Isaac Newton Institute, Cambridge, UK

Introduction

Non-geometry in 4d and 10d SUGRA

Restrict to NSNS sector: g_{mn} , b_{mn} , ϕ .

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Generated by integers Q_k^{mn} , R^{kmn} : “non-geometric fluxes”.

(\Leftrightarrow specific gaugings in gauged SUGRA).

[hep-th/0508133](#) by J. Shelton, W. Taylor, B. Wecht

[hep-th/0210209](#), [hep-th/0512005](#) by A. Dabholkar, C. Hull

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Terms: \checkmark for phenomenology (stab. of moduli, de Sitter sol. ...)

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A 10d non-geometric config. of fields:

typically looks ill-defined: not single-valued, global issues

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- Progress in relating 4d/10d non-geometry:
 - Field redefinition on NSNS fields $\Rightarrow Q, R$ in 10d Lag.
 - New fields globally defined (in some examples)
 \Rightarrow compactify, get scalar potential \checkmark

Field redefinition

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Key object: β : antisymmetric bivector β^{mn} .

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Arguments in GCG: β related to non-geometry / to Q_k^{mn} , R^{kmn}

[hep-th/0609084](#), [arXiv:0708.2392](#) by P. Grange, S. Schäfer-Nameki

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β appears via a reparametrization of the gen. metric \mathcal{H} :

$$\mathcal{H} = \begin{pmatrix} g - bg^{-1}b & bg^{-1} \\ -g^{-1}b & g^{-1} \end{pmatrix} = \begin{pmatrix} \tilde{g} & -\tilde{g}\beta \\ \beta\tilde{g} & \tilde{g}^{-1} - \beta\tilde{g}\beta \end{pmatrix}, \quad \tilde{g} : \text{new metric}$$

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Apply it on NSNS Lagrangian?

 β could be related to non-geo. fluxes \Rightarrow would they appear?

Rewriting of the NSNS Lagrangian

$$\mathcal{L} = e^{-2\phi} \sqrt{|g|} \left(\mathcal{R} + 4(\partial\phi)^2 - \frac{1}{2.3!} H_{kmn} H_{pqr} g^{kp} g^{mq} g^{nr} \right)$$

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(assumption: $\beta^{km} \partial_m \cdot = 0$)

$$\begin{aligned} \mathcal{R} = & \tilde{\mathcal{R}} - \partial_k \tilde{g}_{su} \partial_m \tilde{g}_{pq} \left(2\tilde{g}^{km} \tilde{g}^{uq} \tilde{g}^{ps} + 2\tilde{g}^{pq} \tilde{g}^{ks} \tilde{g}^{mu} + \frac{1}{2} \tilde{g}^{uq} \tilde{g}^{sm} \tilde{g}^{kp} \right) \\ & - \tilde{g}_{pq} \partial_k \beta^{pk} \partial_m \beta^{qm} - \frac{1}{2} \tilde{g}_{pq} \partial_k \beta^{qm} \partial_m \beta^{pk} \\ & + 2\tilde{g}^{km} \tilde{g}^{pq} \partial_k \partial_m \tilde{g}_{pq} + 2\tilde{g}^{km} (G^{-1})_{pq} \partial_k \partial_m G^{qp} \\ & + \partial_m G^{vl} \left(-2\tilde{g}^{mr} \tilde{g}^{ks} (G^{-1})_{lv} \partial_k \tilde{g}_{rs} - \tilde{g}^{rs} \tilde{g}^{km} (G^{-1})_{lv} \partial_k \tilde{g}_{vr} \right. \\ & \quad \left. + \tilde{g}^{ms} \tilde{g}^{ru} (G^{-1})_{lv} \partial_v \tilde{g}_{rs} - \tilde{g}^{km} \tilde{g}^{rs} (G^{-1})_{ls} \partial_k \tilde{g}_{vr} \right) \\ & + \partial_m G^{vl} \left((G^{-1})_{lq} \partial_v G^{qm} + \frac{1}{2} g_{lq} \partial_v G^{mq} \right) \\ & - \partial_m G^{vl} \partial_k G^{ps} \frac{1}{2} \tilde{g}^{km} \left(2(G^{-1})_{lv} (G^{-1})_{sp} + 5(G^{-1})_{sv} (G^{-1})_{lp} + g_{sl} \tilde{g}_{pv} \right) \end{aligned}$$

where $G = \tilde{g}^{-1} + \beta$.

Rewriting of the NSNS Lagrangian

$$\begin{aligned}\mathcal{L} &= e^{-2\phi} \sqrt{|g|} \left(\mathcal{R} + 4(\partial\phi)^2 - \frac{1}{2 \cdot 3!} H_{kmn} H_{pqr} g^{kp} g^{mq} g^{nr} \right) \\ &= e^{-2\tilde{\phi}} \sqrt{|\tilde{g}|} \left(\tilde{\mathcal{R}} + 4(\partial\tilde{\phi})^2 - \frac{1}{2} |Q|^2 \right) + \partial(\dots) = \tilde{\mathcal{L}} + \partial(\dots)\end{aligned}$$

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\Rightarrow also get $R^{mnp} = 3 \beta^{k[m} \partial_k \beta^{np]}$, $|R|^2 = \frac{1}{3!} R^{kmn} R^{pqr} \tilde{g}_{kp} \tilde{g}_{mq} \tilde{g}_{nr}$

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Q -, R -fluxes appear in 10d NSNS via field redefinition
Relation to 4d Q -, R -fluxes/non-geo. terms?
 \Rightarrow dimensional reduction

The dimensional reduction

The 4d scalar potential

Split 10d \Rightarrow 4d max. sym. space-time \times 6d compact \mathcal{M}

Compactification ansatz: $ds_{10}^2 = ds_4^2 + ds_6^2$ (no warp factor),

b, β purely internal

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Only two 4d scalar fields: volume ρ and dilaton σ :

$$g_{6ij} = \rho g_{6ij}^{(0)}, \quad e^{-\phi} = e^{-\phi^{(0)}} \sigma \rho^{-\frac{3}{2}}, \quad e^{\phi^{(0)}} = g_s$$

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\hookrightarrow in $\mathcal{S}_{\text{NSNS}} = \frac{1}{2\kappa^2} \int d^{10}x \mathcal{L}$, integrate 6d, go to 4d Einstein fr.:

$$S_E = M_4^2 \int d^4x \sqrt{|g^E|} \left(\mathcal{R}_4^E + \text{kin} - \frac{1}{M_4^2} V(\rho, \sigma) \right)$$

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arXiv:0712.1196 by E. Silverstein

where $V(\rho, \sigma) = \sigma^{-2} \left(\rho^{-3} V_H^0 + \rho^{-1} V_{\mathcal{R}}^0 \right)$

$$V_H^0 = \frac{M_4^2}{v_0} \int d^6x \sqrt{|g_6^{(0)}|} \frac{1}{12} H_{ijk}^{(0)} H_{lmn}^{(0)} g_6^{il(0)} g_6^{jm(0)} g_6^{kn(0)}$$

$$V_{\mathcal{R}}^0 = -\frac{M_4^2}{v_0} \int d^6x \sqrt{|g_6^{(0)}|} \mathcal{R}_6^{(0)}$$

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$$g_{6ij} = \rho g_{6ij}^{(0)}, \quad e^{-\phi} = e^{-\phi^{(0)}} \sigma \rho^{-\frac{3}{2}}, \quad e^{\phi^{(0)}} = g_s$$

\hookrightarrow in $\mathcal{S}_{\text{NSNS}} = \frac{1}{2\kappa^2} \int d^{10}x \mathcal{L}$, integrate 6d, go to 4d Einstein fr.:

$$S_E = M_4^2 \int d^4x \sqrt{|g^E|} \left(\mathcal{R}_4^E + \text{kin} - \frac{1}{M_4^2} V(\rho, \sigma) \right)$$

arXiv:0712.1196 by E. Silverstein

where $V(\rho, \sigma) = \sigma^{-2} \left(\rho^{-3} V_H^0 + \rho^{-1} V_{\mathcal{R}}^0 \right)$

$$V_H^0 = \frac{M_4^2}{v_0} \int d^6x \sqrt{|g_6^{(0)}|} \frac{1}{12} H_{ijk}^{(0)} H_{lmn}^{(0)} g_6^{il(0)} g_6^{jm(0)} g_6^{kn(0)}$$

$$V_{\mathcal{R}}^0 = -\frac{M_4^2}{v_0} \int d^6x \sqrt{|g_6^{(0)}|} \mathcal{R}_6^{(0)}$$

Non-geometric terms? Most general (NSNS) potential:

$$V(\rho, \sigma) = \sigma^{-2} \left(\rho^{-3} V_H^0 + \rho^{-1} V_{\mathcal{R}}^0 + \rho V_Q^0 + \rho^3 V_R^0 \right)$$

arXiv:0711.2512 by M. P. Hertzberg, S. Kachru, W. Taylor, M. Tegmark

The dimensional reduction

The 4d scalar potential

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Split 10d \Rightarrow 4d max. sym. space-time \times 6d compact \mathcal{M}

Compactification ansatz: $ds_{10}^2 = ds_4^2 + ds_6^2$ (no warp factor),

b, β purely internal

Only two 4d scalar fields: volume ρ and dilaton σ :

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[arXiv:0711.2512](#) by M. P. Hertzberg, S. Kachru, W. Taylor, M. Tegmark

With $\tilde{\mathcal{L}}$ instead of \mathcal{L} , we get the \checkmark 4d potential

We get $V(\rho, \sigma) = \sigma^{-2} (\rho^{-1} V_{\mathcal{R}}^0 + \rho V_Q^0 + \rho^3 V_R^0)$ where

$$V_{\mathcal{R}}^0 = -\frac{M_4^2}{v_0} \int d^6x \sqrt{|\tilde{g}_6^{(0)}|} \tilde{\mathcal{R}}_6^{(0)}$$

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$$V_Q^0 = -\frac{M_4^2}{v_0} \int d^6x \sqrt{|\tilde{g}_6^{(0)}|} \left(-\frac{1}{4} \tilde{g}_{ik} \tilde{g}_{jl} \tilde{g}^{rs} Q_r^{kl} Q_s^{ij} + \frac{1}{2} \tilde{g}_{pq} Q_k^{lp} Q_l^{kq} \right. \\ \left. + \tilde{g}_{jl} \tilde{g}_{pq} \beta^{jm} (Q_k^{lp} \partial_m \tilde{g}^{kq} + \partial_k \tilde{g}^{lp} Q_m^{kq}) \right. \\ \left. - \frac{1}{4} \tilde{g}_{ik} \tilde{g}_{jl} \tilde{g}_{pq} (\beta^{pr} \beta^{qs} \partial_r \tilde{g}^{kl} \partial_s \tilde{g}^{ij} - 2\beta^{ir} \beta^{js} \partial_r \tilde{g}^{lp} \partial_s \tilde{g}^{kq}) \right. \\ \left. + \frac{1}{2\sqrt{|\tilde{g}|}} \tilde{g}^{pq} \partial_k \tilde{g}_{pq} \partial_m \left(\sqrt{|\tilde{g}|} \tilde{g}_{ij} \beta^{ik} \beta^{jm} \right) \right)^{(0)}_6$$

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$\tilde{\mathcal{L}}$ and 10d Q, R give the \checkmark 4d potential (give a 10d origin)

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$\tilde{\mathcal{L}}$ and 10d Q, R give the \checkmark 4d potential (give a 10d origin)

4d Q is not clearly identified...

\hookrightarrow DFT brings an interpretation for the Q -terms, the role of Q
(see talk of Peter Patalong)

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Global aspects

We have shown

$$\mathcal{L}(g, b, \phi) = \tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi}) + \partial(\dots)$$

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For a 10d NSNS non-geometric configuration: **global issues**
 \hookrightarrow no $\int d^6 x$, no $V_H^0 \dots$

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In some examples, $\tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi}) \checkmark$

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In some examples, $\tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi})$ ✓

We propose to discard $\partial(\dots)$, use $\tilde{\mathcal{L}}$ as the good low-energy effective description of string theory

Global aspects

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\hookrightarrow integrate $\tilde{\mathcal{L}} \checkmark$, get V_Q^0, V_R^0

Global aspects

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$$\mathcal{L}(g, b, \phi) = \tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi}) + \partial(\dots)$$

For a 10d NSNS non-geometric configuration: **global issues**
 \hookrightarrow no $\int d^6x$, no V_H^0 ...

In some examples, $\tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi}) \checkmark$

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Relation between 4d/10d non-geometry:
for a 10d non-geometric configuration,
field redefinition + dimensional reduction
 \Rightarrow generates 4d non-geometric terms

Conclusion and outlook

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- GCG \Rightarrow field redefinition $(g, b, \phi) \leftrightarrow (\tilde{g}, \beta, \tilde{\phi})$
Rewriting NSNS Lag.: $\mathcal{L}(g, b, \phi) = \tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi}) + \partial(\dots)$
10d $Q_k{}^{mn} = \partial_k \beta^{mn}$ (for $\beta^{km} \partial_m \cdot = 0$), $R^{mnp} = 3 \beta^{k[m} \partial_k \beta^{np]}$
- 10d NSNS non-geometry $\Rightarrow \tilde{\mathcal{L}}(\tilde{g}, \beta, \tilde{\phi})$ has no global issue...
 \hookrightarrow dim. reduction of $\tilde{\mathcal{L}} \Rightarrow$ 4d non-geometric potential \checkmark
Relation between 4d/10d non-geometry
- Extend to RR sector (S-duality)
and D-brane/O-plane sources (new objects?)
 \hookrightarrow new interesting backgrounds, 10d de Sitter solutions...
- Relations between non-commutativity and non-geometry