

# Moduli Stabilisation of Heterotic String Theory with Flux

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Work done in collaboration with:  
Michael Klaut, Andre Lukas, Cyril Matti,  
arXiv:1005.5302, arXiv:1107.3573, upcoming

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$$ds^2 = ds_{2+1}^2 + dy^2 + g_{uv}(x^m)dx^u dx^v.$$

This allows the internal space  $X$  to have torsion (a non-Calabi-Yau compactification), which may help with moduli stabilisation. Drawback: it singles out a preferred direction (the  $y$ -direction), and we must work to get rid of the  $y$ -dependence of our fields.

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- Supersymmetry then implies that  $\mathbb{R}_y \times X$  has a  $G_2$ -structure, which gives  $X$  a half-flat  $SU(3)$ -structure. We will take  $X$  to be a coset space of the form  $X = G/H$  (a homogeneous space).

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- Can write down explicitly  $g_{uv}$ ,  $J$  and  $\Omega$ .
- No complex structure moduli:  $H^3(X) = 0$ .





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- On  $X = SU(3)/U(1)^2$  there are three  $G$ -invariant two-forms  $\{\omega_1, \omega_2, \omega_3\}$ , dual four-forms  $\{\tilde{\omega}_1, \tilde{\omega}_2, \tilde{\omega}_3\}$ , and two "symplectic"  $G$ -invariant three-forms  $\{\alpha_0, \beta^0\}$ .

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- The Kähler form, and the complex three-form may be expanded in terms of these

$$J = v^i \omega_i, \quad \Omega = \mathcal{Z} \alpha_0 - \mathcal{G}_{\mathcal{Z}} \beta^0,$$

Similar to the Calabi-Yau setting. Here  $\mathcal{G}$  is a second order homogeneous polynomial in  $\mathcal{Z}$  (the prepotential in the Calabi-Yau setting), and  $\mathcal{G}_{\mathcal{Z}} = \partial_{\mathcal{Z}} \mathcal{G}$ .

- In Michael Klaut's talk, we heard that this domain wall solution solves the supersymmetry equations and the Bianchi identity exactly, and the equations of motion are solved to  $\mathcal{O}(\frac{\alpha'^2}{R^4})$ . We now want to match this to a 4d supergravity.

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- We also expand the NS-NS two-form and three forms in terms of the internal geometry:

$$B = b^i \omega_i,$$

$$H = dB + \frac{\alpha'}{4}(\omega_{YM} - \omega_L) = b^i e_i \beta - \mu \alpha,$$

where  $\mu$  is a constant to  $\mathcal{O}(\frac{\alpha'^2}{R^4})$ , and is determined by calculating the Chern-Simons forms of the connections. The constant  $\mu$  will thus depend on bundle parameters.

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- We have a Kahler potential, and the Gukov-Vafa-Witten superpotential:

$$W \sim \int_X \Omega \wedge (H + idJ) = \sqrt{8}(e_i T^i - \mu \frac{\mathcal{G}_Z}{Z}).$$



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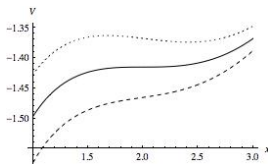
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- Trying to set  $D_3 W = D_5 W = 0$ , we find that this is not possible. We get either a constant  $v^3$  and the dilaton blowing up as  $|y| \rightarrow \infty$  ( $\mu < 0$ ), or  $v^3$  blowing up and a constant dilaton ( $\mu \geq 0$ ). This is as in the 10d case. We conclude that supersymmetric solutions with all fields constant do not exist perturbatively. We do however find that including  $\alpha'$  effects makes it possible to stabilise all internal moduli consistently when  $\mu < 0$ .

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- Note that  $W$  is independent of  $\phi$ , and so  $\phi$  is a flat direction. Solved by adding non-perturbative effects, e.g. a gauging condensate  $W_g = ke^{-cS}$ . Here  $k$  is a constant of  $\mathcal{O}(1)$  (in units of  $\alpha'$ ), while  $c$  is a number depending on the condensing gauge group (its Coxeter number).

# Moduli stabilisation: Non-perturbative effects

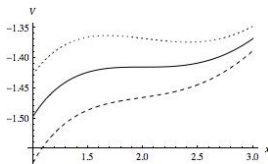
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- For  $\mu < 0$ , we find that there is a minimum value of  $k$ ,  $k = k_{min}$ , where a stable AdS solution to the supergravity potential may be obtained:



**Figure:** The supergravity potential as a function of  $x = cs$ . We have set  $v^1 = v^2 = 0$ , and  $v^3$  to the value corresponding to  $\partial_{v^3} V = 0$ . The dashed line corresponds to  $k < k_{min}$ , the full line corresponds to  $k = k_{min}$ , and the dotted line corresponds to  $k > k_{min}$ .

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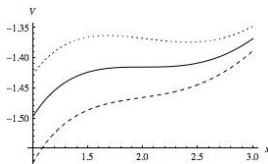


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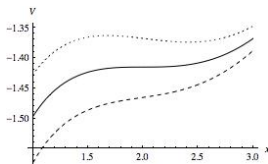
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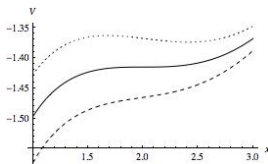
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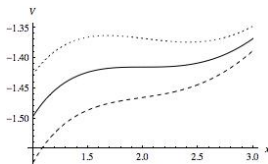
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- Example 2:  $v \approx 15\alpha'$ ,  $\alpha_{YM} = 0.63$ ,  $k_{min} \approx 18\alpha'$ .

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- Worst case, we can still learn something about the physics of non-Kähler compactifications (e.g. Strominger system).



- Thank you for your attention!