

Renormalized relaxed Lagrangian solutions for the Semi-Geostrophic system in physical space

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- 1 The Semi-Geostrophic system
 - SG in physical space
 - SG in a more readable form
 - SG in dual space
- 2 Lagrangian solutions
 - Yet another system
 - Weak Lagrangian solutions in physical space
 - Case of singular measures in dual space
- 3 Renormalized relaxed solutions
 - Introduction
 - Definition
 - Solutions in dual space revisited
 - Stability, existence and continuity in physical space
 - Return to dual space; energy conservation

Standard formulation (3D with rigid boundary)

Let $\Omega \subset \mathbb{R}^3$ be an open bounded set and $T > 0$. The semigeostrophic (SG) system ([Eliassen 1948](#), [Hoskins 1974](#)) in $(0, T) \times \Omega$:

$$\begin{aligned} D_t \mathbf{v}_g + \mathbf{e}_3 \times \mathbf{v} &= -\bar{\nabla} p, \\ \nabla p + \mathbf{e}_3 \times \mathbf{v}_g &= \theta \mathbf{e}_3, \\ \nabla \cdot \mathbf{u} &= 0, \\ D_t \theta &= 0. \end{aligned}$$

The initial and boundary data:

$$\begin{aligned} \mathbf{u} \cdot \nu &= 0, \quad \text{on } [0, T) \times \partial\Omega, \\ p(0, x) &= p_0(x) \quad \text{in } \Omega. \end{aligned}$$

Notation: $\mathbf{v}_g = \mathbf{e}_3 \times \bar{\nabla} p$ is the *geostrophic wind*, $\mathbf{u} = \mathbf{v}_g + \mathbf{v}_a = \mathbf{v} + u_3 \mathbf{e}_3 \in \mathbb{R}^3$ is the total wind velocity (\mathbf{v}_a is the *ageostrophic wind*), p is the pressure, θ is the potential temperature; $\bar{\nabla} = (\partial_{x_1}, \partial_{x_2}, 0)$, $D_t = \partial_t + \mathbf{u} \cdot \nabla$.

SG in terms of the modified pressure

- Introducing the function

$$P(t, x) = p(t, x) + (x_1^2 + x_2^2)/2,$$

we rewrite SG as the following system of equations for (P, \mathbf{u}) :

$$D_t X = J(X - x),$$

$$\nabla \cdot \mathbf{u} = 0,$$

$$X = \nabla P,$$

$$\mathbf{u} \cdot \nu = 0, \quad \text{on } [0, T) \times \partial\Omega,$$

$$P(0, x) = P_0(x) \quad \text{in } \Omega,$$

- Here $J = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$.

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Cullen-Purser stability

- Let (P, \mathbf{u}) be a solution and set $\mathcal{X}(t) := \{Y : \Omega \rightarrow \mathbb{R}^3 \text{ Borel} : Y_{\#}\chi = \nabla P(t, \cdot)_{\#}\chi\}$.
- (Cullen & Shutts) Then $X(t, \cdot) = \nabla P(t, \cdot)$ is a critical point for

$$I(Y) = \int_{\Omega} |Y(x) - x|^2 dx$$

among all maps $Y \in \mathcal{X}(t)$, for all $t \in (0, T)$.

- (Cullen & Purser) We are looking for minimizers.
- In the language of Optimal Transport, this means $X(t, \cdot)$ must be the gradient of a convex function, i.e. $P(t, \cdot)$ must be convex for all $t \in (0, T)$.

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Do we have solutions?

- In the general case, \mathbf{u} may be a singular measure. We need to understand what we mean by “solution”.

Theorem (Feldman & T.)

Let (P, \mathbf{u}) be a distributional solution for SG in the physical space such that $\nabla P \in H^1(0, T; L^2(\Omega; \mathbb{R}^3))$. Then $\alpha_t := \nabla P_{t\#} \chi$ is atom-free for \mathcal{L}^1 -a.e. $t \in (0, T)$.

- We would like the model to accommodate solutions for which ∇P_t could be locally constant.

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A change of variable and a new system

- Let $\nabla P(t, \cdot)_{\#} \chi =: \alpha(t, \cdot)$. If $\mathcal{L}^3(\Omega) = 1$, then α_t is a Borel probability on \mathbb{R}^3 .
- The path of measures $t \mapsto \alpha_t$ solves:

$$\partial_t \alpha + \nabla \cdot (U \alpha) = 0 \quad \text{in } [0, T) \times \mathbb{R}^3, \quad (1)$$

$$\nabla P(t, \cdot)_{\#} \chi_{\Omega} = \alpha(t, \cdot) \quad \text{for any } t \in [0, T);$$

$$U(t, X) = J[X - \bar{\gamma}(t, X)], \quad (2)$$

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where $\bar{\gamma}(t, \cdot)$ is the barycentric projection of $\nabla P^*(t, \cdot)$.

- Note that equation (1) represents the fact that the dual flow $t \mapsto \alpha_t$ is weakly (in the sense of distributions) transported by the dual velocity U defined by (2).

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- Solutions in “dual variables” [Benamou & Brenier](#), [Cullen & Gangbo](#), [Cullen & Maroofi](#).
- Hamiltonian system treatment [Gangbo & Pacini](#), [Ambrosio & Gangbo](#).
- Dual-space solutions are obtained either through variational techniques, or by iterative methods.
- Lagrangian solutions in physical space in the case of nonsingular dual-space solutions [Cullen & Feldman](#).
- Eulerian solutions in physical space in special cases (“nice” dual-space solutions) [Ambrosio](#), [Colombo](#), [De Philippis & Figalli](#).
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Formal Lagrangian flow

- Assume one has a classical solution to SG in physical space.
- If \mathbf{u} has a flow $F : [0, T] \times \Omega \rightarrow \Omega$ defined by

$$\partial_t F(t, x) = \mathbf{u}(t, F(t, x)), \quad F(0, x) = x,$$

then F can replace \mathbf{u} as an unknown. Let $Z(t, x) := \nabla P(t, F(t, x))$.

- The system for (P, F) becomes

$$\begin{aligned} \partial_t Z(t, x) &= J[Z(t, x) - F(t, x)] && \text{in } [0, T) \times \Omega, \\ Z(0, x) &= \nabla P_0(x) && \text{in } \Omega. \end{aligned} \tag{LE}$$

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Definition in the general case

Let $P_0 \in W^{1,\infty}(\Omega)$, $F : [0, T) \times \Omega \rightarrow \Omega$ and $P : [0, T) \times \Omega \rightarrow \mathbb{R}$ be as before. Set $\alpha_t := \nabla P_{t\#\chi}$, $\mu_t := \bar{\gamma}_{t\#\alpha_t}$. Then the pair (P, F) is called a weak Lagrangian solution for SG in $[0, T) \times \Omega$ if

- $F(0, x) = x$ for μ_0 -a.e. $x \in \Omega$, $P(0, x) = P_0(x)$ for a.e. $x \in \Omega$,
- for any $t > 0$ the mapping $F_t = F(t, \cdot) : \Omega \rightarrow \Omega$ satisfies $F_{t\#\mu_0} = \mu_t$;
- there exists a Borel mapping $F^* : [0, T) \times \Omega \rightarrow \Omega$ satisfying $F_{t\#\mu_t}^* = \mu_0$ and $F_t^* \circ F_t = \text{Id}$ μ_0 -a.e., $F_t \circ F_t^* = \text{Id}$ μ_t -a.e. ;
- The function $Z : (0, T) \times \Omega \rightarrow \mathbb{R}^3$ defined by $Z_t = \nabla P_t \circ F_t$ lies, along with F , in $L^1((0, T) \times \Omega; \mathbb{R}^3)$ and is a distributional solution of (LE).

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Definition in the general case

Let $P_0 \in W^{1,\infty}(\Omega)$, $F : [0, T) \times \Omega \rightarrow \Omega$ and $P : [0, T) \times \Omega \rightarrow \mathbb{R}$ be as before. Set $\alpha_t := \nabla P_{t\#}\chi$, $\mu_t := \bar{\gamma}_{t\#}\alpha_t$. Then the pair (P, F) is called a weak Lagrangian solution for SG in $[0, T) \times \Omega$ if

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Existence and consistency

- **Cullen & Feldman** used **Ambrosio's** theory of regular Lagrangian flows to construct weak Lagrangian solutions in physical space in the case $\alpha_0 \in L^p(\mathbb{R}^3)$ for some $p > 1$.
- Recall $U_t = J[\text{Id} - \nabla P_t^*]$. There exists (**Ambrosio**) a unique Borel bounded, \mathcal{L}^3 -measure-preserving (for a.e. t) vector-map $\Phi(t, X)$ satisfying $\dot{\Phi} = U(\cdot, \Phi)$, \mathcal{L}^4 -a.e. in $(0, T) \times \nabla P_0(\Omega)$ and $\Phi(0, \cdot) = \text{Id}$, α_0 -a.e. in \mathbb{R}^3 .
- Set $F_t := \nabla P_t^* \circ \Phi_t \circ \nabla P_0$. Then (P, F) is a weak Lagrangian solution of SG in physical space (**Cullen & Feldman**).
- **Cullen & Feldman** proved that by setting $u(t, x) := \partial_t F(t, F^*(t, x))$ under the assumption $\partial_t F \in L^\infty((0, T) \times \Omega; \mathbb{R}^3)$, we obtain that the pair (P, u) is a weak (Eulerian) solution of SG in physical space.
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When α_0 is singular; an example

- If $\alpha_0 = \delta_{z_0}$ for some $z_0 \in \mathbb{R}^3$, one can readily check that if $\dot{z}(t) = Jz(t)$, $z(0) = z_0$, then $\alpha_t := \delta_{z(t)}$ solves the SG in dual space.
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Dual flow gives weak Lagrangian solutions in physical space

Theorem (Feldman & T.)

Assume there exists a Borel flow-map $\Phi : [0, T] \times \mathbb{R}^3 \rightarrow \mathbb{R}^3$ for U such that $\Phi_{t\#}\alpha_0 = \alpha_t$ for all $t \in [0, T)$. Then the pair (P, F) is a weak Lagrangian solution in physical space satisfying all properties but the existence of F^* from the previous definition. Furthermore, if there also exists a Borel map $\Phi^* : [0, T] \times \mathbb{R}^3 \rightarrow \mathbb{R}^3$ such that $\Phi_{\#}^*\alpha_t = \alpha_0$ for all $t \in [0, T)$, then (P, F) is a weak Lagrangian solution in the physical space.

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- To obtain Φ in the general case.
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- If $\alpha_0 \in L^p(\mathbb{R}^3)$ for some $1 \leq p < \infty$, we have (P, F) a weak Lagrangian solution with F_t volume preserving for all $t \in [0, T)$ and $F_0 = \text{Id}_\Omega$.
- Set $\sigma_t := (\text{Id} \times F_t)_{\#} \chi$. This is a probability measure on $\Omega \times \Omega$ with both marginals χ . Also, $\sigma_0 = (\text{Id} \times \text{Id})_{\#} \chi$.
- Set $\sigma(dt, dx, dy) := \sigma_t(dx, dy)dt$.
- Let $\xi \in C_b^1(\mathbb{R}^3)$. Then $t \mapsto \xi(Z(t, x))$ is absolutely continuous for χ -a.e. $x \in \Omega$ and

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Renormalized Relaxed Lagrangian solutions

Definition. Let P_0 be convex on Ω . Let $P : [0, T) \times \Omega \rightarrow \mathbb{R}$, and let $\sigma = \int_0^T \sigma_t dt$ be a Borel measure on $[0, T) \times \Omega \times \Omega$. Then (P, σ) is a renormalized relaxed Lagrangian solution of SG with initial data P_0 if

- (i) P_t is convex for each $t \in [0, T)$;
- (ii) $\pi_{k\#}\sigma_t = \chi$ for each $t \in [0, T)$, where $\pi_k(\mathbf{x}) = x_k$ for $\mathbf{x} = (x_1, x_2) \in \Omega \times \Omega$, $k = 1, 2$;
- (iii) for any $\xi \in C_b^1(\mathbb{R}^3)$, $\zeta \in C_c^1([0, T) \times \Omega)$

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- Equation (iii) well-defined for $\nabla P \in L^\infty([0, T); L^1(\Omega))$ by (ii).

Renormalized Relaxed Lagrangian solutions

Definition. Let P_0 be convex on Ω . Let $P : [0, T) \times \Omega \rightarrow \mathbb{R}$, and let $\sigma = \int_0^T \sigma_t dt$ be a Borel measure on $[0, T) \times \Omega \times \Omega$. Then (P, σ) is a **renormalized relaxed Lagrangian solution** of SG with initial data P_0 if

- (i) P_t is convex for each $t \in [0, T)$;
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Global Lipschitz solutions in dual space. Weak stability

Theorem (Ambrosio & Gangbo, Feldman & T.)

Let $\alpha_0 \in \mathcal{P}_2(\mathbb{R}^3)$ be given. Then there exists a distributional solution $\alpha \in AC^\infty(0, \infty; \mathcal{P}_2(\mathbb{R}^3))$ for SG in dual space with $\alpha(0) = \alpha_0$, where $\bar{\gamma}_t$ is the barycentric projection of the optimal plan between α_t and χ . In other words, the curve α is globally Lipschitz and satisfies SG in dual space in the sense of distributions in $(0, \infty) \times \mathbb{R}^3$. Furthermore, the Hamiltonian energy $t \rightarrow H(\alpha_t) := -W_2^2(\alpha_t, \chi)/2$ is conserved.

- Since $[0, \infty) \ni t \rightarrow \alpha_t \in \mathcal{P}_2(\mathbb{R}^3)$ is continuous, we conclude that there is a family $P \in C([0, \infty); H^1(\Omega))$ of convex functions $P(t, \cdot)$ such that $\nabla_x P(t, \cdot) =: \nabla P_t$ pushes χ forward to α_t optimally for all $t \geq 0$.

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Weak stability in dual space

Theorem (Feldman & T.)

Let $\{\alpha_0\} \cup \{\alpha_0^n\}_{n \geq 1} \subset \mathcal{P}_2(\mathbb{R}^3)$ be such that

$$\alpha_0^n \rightharpoonup \alpha_0 \text{ weakly } \star \text{ as measures, and } \sup_{n \geq 1} W_2(\alpha_0^n, \alpha_0) < \infty.$$

Let α^n be solutions as in previous theorem corresponding to the initial data α_0^n . Then, possibly up to a subsequence, α^n converges to a solution $\alpha \in AC^\infty(0, \infty; \mathcal{P}_2(\mathbb{R}^3))$ with initial data α_0 . More precisely, we have

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Stability of optimal maps

Theorem (Feldman & T.)

Let $\Omega \in \mathbb{R}^d$ be open and bounded and consider the Borel measures $\{\mu_n\}_{n \geq 1}$, μ in \mathbb{R}^d with total mass equal to $\mathcal{L}^d(\Omega)$ and with finite second moments. Denote by M_n , M the Brenier optimal maps for pushing $\mathcal{L} := \mathcal{L}^d|_{\Omega}$ forward to μ_n , μ for the quadratic cost (i.e., gradients of convex functions such that $M_n\#\mathcal{L} = \mu_n$, $M\#\mathcal{L} = \mu$). Then, for any $1 < p \leq 2$, we have

$$\lim_{n \rightarrow \infty} W_p(\mu_n, \mu) = 0 \text{ is equivalent to } \lim_{n \rightarrow \infty} \|M_n - M\|_{L^p(\Omega; \mathbb{R}^d)} = 0.$$

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Weak stability for renormalized relaxed solutions

Corollary

Let $P_0 \in H^1(\Omega)$, $\{P_0^n\}_n \subset H^1(\Omega)$ be convex and such that $\nabla P_0^n \rightarrow \nabla P_0$ in $L^2(\Omega; \mathbb{R}^3)$. Assume (P^n, σ^n) are renormalized relaxed solutions for SG in physical space corresponding to the initial data P_0^n . Then, possibly up to a subsequence, (P^n, σ^n) converges to a renormalized relaxed solution (P, σ) corresponding to the initial datum P_0 . This convergence is in the following sense:

- (i) $\nabla P_t^n \rightarrow \nabla P_t$ in $L^p(\Omega; \mathbb{R}^3)$ for all $t \in [0, T)$ and all $1 \leq p < 2$;
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Furthermore, the dual space solutions satisfy $W_p(\alpha_t^n, \alpha_t) \rightarrow 0$ for all $t \in [0, T)$ and all $1 \leq p < 2$.

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Existence of renormalized relaxed solutions

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Let $P_0 \in H^1(\Omega)$. Then there exists a renormalized relaxed solution (P, σ) corresponding to the initial datum P_0 .

- Idea of proof: approximate P_0 in $H^1(\Omega)$ by “nice” convex functions P_0^n such that $\nabla P_0^n \# \chi =: \alpha_0^n \in L^1(\mathbb{R}^3)$, then use **Cullen & Feldman**'s solutions (P^n, F^n) to construct $\sigma_t^n := (\text{Id} \times F_t^n) \# \chi$. Then use the *a priori* stability result (previous corollary) to conclude.

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Time-continuity

- Define $G_t : \Omega \times \Omega \rightarrow \mathbb{R}^3 \times \mathbb{R}^3$ by $G_t(x, y) = (x, \nabla P_t(y))$, then $t \rightarrow G_{t\#}\sigma_t$ is continuous with respect to narrow convergence on $\Omega \times \mathbb{R}^3$: i.e., from the equation, for any $\xi, \zeta \in C_c^1(\mathbb{R}^3)$

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is continuous.

- Also, for any $s \in [0, T]$: $t \rightarrow G_{s\#}\sigma_t$ is continuous at $t = s$ with respect to narrow convergence on $\Omega \times \mathbb{R}^3$.
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Under-determinacy in the singular case; an example

- Let $\Omega = B(0, 1)$, $P_0 \equiv 0$ and $T = 2\pi$.
- The solution obtained by the described procedure is $d\tilde{\sigma}dt$, where

$$\begin{aligned} \int_{\Omega^2} \zeta d\tilde{\sigma} &= \int_0^1 \int_0^{2\pi} \int_0^{2\pi} r \zeta(r \cos \theta, r \sin \theta, r \cos \vartheta, r \sin \vartheta) d\vartheta d\theta dr \\ &= \int_{\mathcal{S}} \frac{\zeta(\mathbf{x})}{|\mathbf{x}|} dS(\mathbf{x}), \end{aligned}$$

where \mathcal{S} is the hypersurface in $\Omega \times \Omega$ of equation $x^2 = y^2$ for $\mathbf{x} := (x, y) \in \Omega \times \Omega$.

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Renormalization property enables return to dual space

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- Let (P, σ) be a renormalized relaxed Lagrangian solution and $\alpha_t = \nabla P_{t\#}\chi$. Then (P, α) is a distribution solution in dual space. In particular, the geostrophic energy is conserved.
- Renormalization property is used:** for test functions $\varphi(t)\xi(X)$ in dual space with $\varphi(0) = 0$ (for simplicity), we have

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 \int_0^T \int_{\mathbb{R}^3} \dot{\varphi}(t)\xi(X)\alpha_t(dX)dt &= \int_0^T \int_{\Omega} \dot{\varphi}(t)\xi(\nabla P_t(y))dydt \\
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- $I_1 = \int_0^T \int_{\Omega} \varphi(t) \nabla \xi(\nabla P_t(y)) \cdot J \nabla P_t(y) dy dt = \int_0^T \int_{\mathbb{R}^3} \varphi(t) \nabla \xi(X) \cdot JX \alpha_t(dX) dt.$
- Denote $\gamma_t = (\text{Id} \times \nabla P_t) \# \chi.$

$$\begin{aligned} I_2 &= - \int_0^T \int_{\Omega} \varphi(t) \nabla \xi(\nabla P_t(y)) \cdot Jy dy dt = - \int_0^T \int_{\Omega \times \mathbb{R}^3} \varphi(t) \nabla \xi(X) \cdot Jy \gamma_t(dy, dX) dt \\ &= - \int_0^T \int_{\mathbb{R}^3} \varphi(t) \nabla \xi(X) \cdot J\bar{\gamma}_t(X) \alpha_t(dX) dt \end{aligned}$$

- We have obtained $\int_0^T \int_{\mathbb{R}^3} [\partial_t \zeta + \nabla \zeta \cdot U] \alpha_t(dX) dt = 0,$ for $U(t, X) = J(X - \bar{\gamma}_t(X)),$
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 I_2 &= - \int_0^T \int_{\Omega} \varphi(t) \nabla \xi(\nabla P_t(y)) \cdot Jy dy dt = - \int_0^T \int_{\Omega \times \mathbb{R}^3} \varphi(t) \nabla \xi(X) \cdot Jy \gamma_t(dy, dX) dt \\
 &= - \int_0^T \int_{\mathbb{R}^3} \varphi(t) \nabla \xi(X) \cdot J\bar{\gamma}_t(X) \alpha_t(dX) dt
 \end{aligned}$$

- We have obtained $\int_0^T \int_{\mathbb{R}^3} [\partial_t \zeta + \nabla \zeta \cdot U] \alpha_t(dX) dt = 0$, for $U(t, X) = J(X - \bar{\gamma}_t(X))$,
 $\zeta(t, x) = \varphi(t) \xi(X).$

Renormalization property enables return to dual space

(2)

- $I_1 = \int_0^T \int_{\Omega} \varphi(t) \nabla \xi(\nabla P_t(y)) \cdot J \nabla P_t(y) dy dt = \int_0^T \int_{\mathbb{R}^3} \varphi(t) \nabla \xi(X) \cdot JX \alpha_t(dX) dt.$
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Thank you!