

On the regularity of Lagrangian trajectories for suitable weak solutions of the Navier–Stokes equations

Topological Fluid Dynamics (IUTAM Symposium)

July 25, 2012

Setting of the problem

- The Navier-Stokes equations

$$\begin{cases} u_t - \Delta u + u \cdot \nabla u + \nabla p = 0 \\ \nabla \cdot u = 0 \end{cases}$$

- Standard regularity of weak solutions

$$u \in L^2(0, T; V) \cap L^\infty(0, T; H)$$

- Regularity in $L^w L^s$ spaces: $u \in L^w(0, T; L^s(\Omega))$ for

$$\frac{2}{w} + \frac{3}{s} \leq \frac{3}{2}, \quad 2 \leq s \leq 6$$

- Serrin's condition

$$\frac{2}{w} + \frac{3}{s} \leq 1, \quad s \geq 3$$

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How weak is a weak solution?

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- uniqueness?
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Foias *et al*

- Lagrange trajectories

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- Foias *et al.* 1981: $u \in L^{2/3}(0, T; D(A)) \cap L^2(0, T; V)$

$$\Rightarrow u \in L^1(0, T; L^\infty(\Omega))$$

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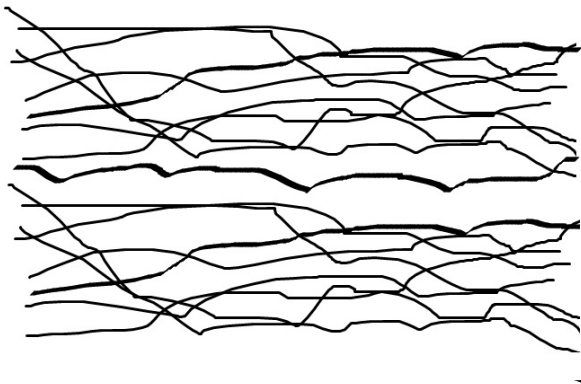
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Foias *et al* 1985

There exists at least one function $\Phi : \Omega \times [0; T] \rightarrow \Omega$ such that

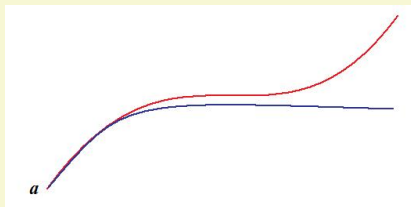
- 1 the function $\phi_a(\cdot) = \Phi(a, \cdot)$ satisfies

$$\phi(t) = a + \int_0^t u(\phi(s), s) ds$$

- 2 $\phi_a(\cdot) \in W^{1,1}(0, T)$
- 3 the mapping $a \rightarrow \Phi(a, \cdot)$ belongs to $L^\infty(\Omega; C([0, T], \bar{\Omega}))$
- 4 Φ is volume-preserving: for any Borel set $B \subset \Omega$,

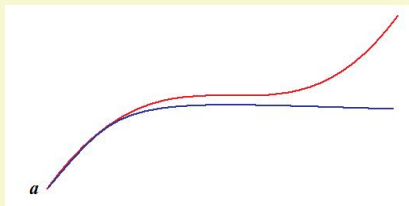
$$\mu(\Phi(\cdot, t)^{-1}(B)) = \mu(B)$$

- ① Question: Are the particle trajectories unique for almost all $a \in \Omega$?



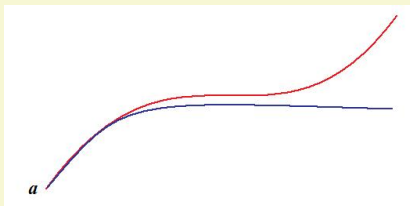
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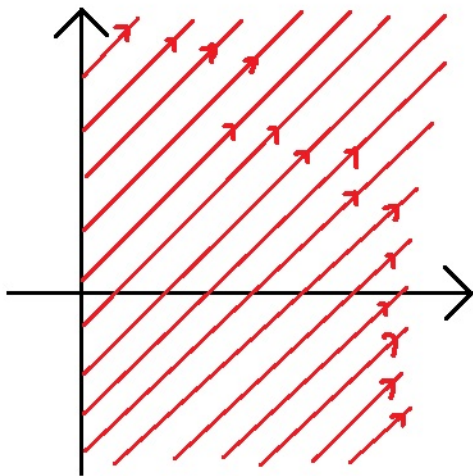


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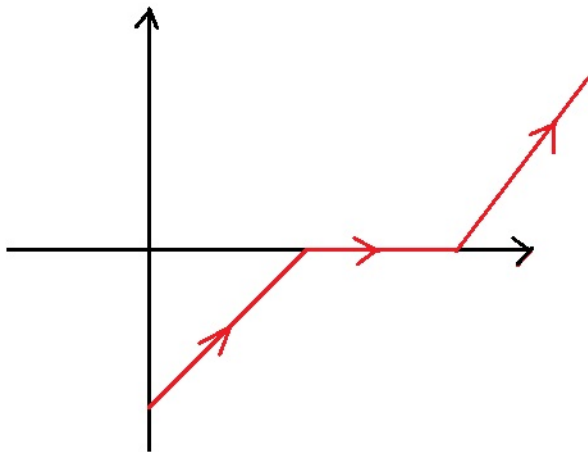
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- 2 How regular are they?



$$f(x) = 1$$



Suitable weak solutions

- $p \in L^{5/3}$
- local energy inequality

\Rightarrow

- small singular set: the Hasdorff dimension no greater than 1
- small singular set: the box-counting dimension no greater than $5/3$

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Definition

- 1 A point (x, t) is called regular if u is Hölder continuous in some neighbourhood of (x, t) . A point is singular if it is not regular. The set of singular points in $(0, T) \times \Omega$ we denote by S .

CKN - 1982

- 1 The set S of singular points of a suitable weak solution u has 1-dimensional parabolic Hausdorff measure zero: for any $n \in \mathbb{N}$ it can be covered by cylinders $Q_k = (t_k, t_k + r_k^2) \times B(x_k, r_k)$ such that

$$\sum_{k=1}^{\infty} r_k < \frac{1}{n}$$

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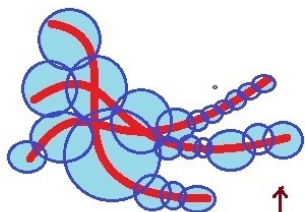
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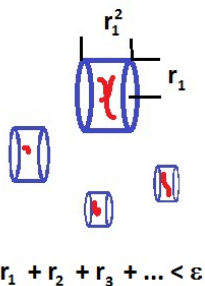
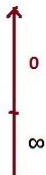
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Hausdorff dimension

$$r_1^s + r_2^s + r_3^s + \dots$$



$$r_1 + r_2 + r_3 + \dots < \varepsilon$$

Conditional result

Theorem 1. James Robinson and WS 2008

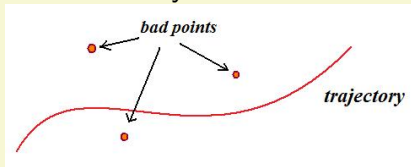
If u is a suitable weak solution with

$$u \in L^{6/5}(0, T; L^\infty)$$

corresponding to $u_0 \in H \cap H^{1/2}$, then almost every initial condition $a \in \Omega$ gives rise to a unique particle trajectory, which is C^1 function of time.

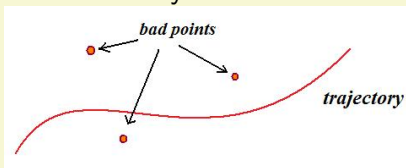
Idea of proof

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Avoiding "bad points"

Theorem 2

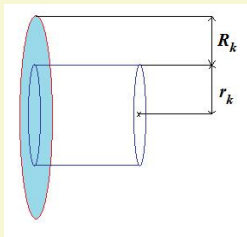
If u is a suitable weak solution with $u \in L^{6/5}(0, T; L^\infty)$ then the set of initial conditions $a \in \Omega$ that give rise to trajectories intersecting the singular set S is of Lebesgue measure zero.

Proof of Theorem 2 - part one

- 1 We cover the singular set by cylinders Q_k
- 2 We define the numbers R_k :

$$R_k = \int_{t_k}^{t_k + r_k^2} \|u\|_\infty$$

- 3 Then we consider balls $\hat{B}_k = (x_k, r_k + R_k)$

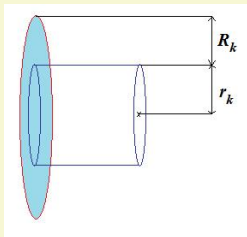


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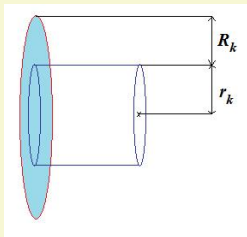


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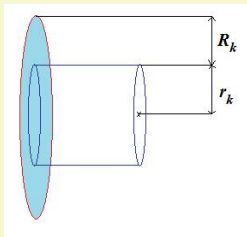


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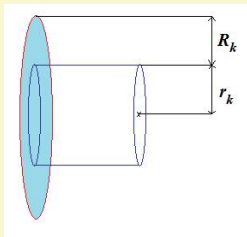


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Proof of Theorem 2 - part two

- ① using inequality:

$$|\varphi_a(t) - \varphi_a(t_k)| \leq R_k$$

we prove that trajectories that do not meet \hat{B}_k at time t_k cannot enter cylinder Q_k

- ② we estimate the volume of the family of balls \hat{B}_k :

$$\sum_{k=1}^{\infty} \mu(\hat{B}_k) \leq \sum_{k=1}^{\infty} c(r_k + R_k)^3 \leq C \left[\sum_{k=1}^{\infty} r_k^3 + \sum_{k=1}^{\infty} R_k^3 \right]$$

- ③ Proof

$$R_k \leq \left(\int_{t_k}^{t_k+r_k^2} ds \right)^{1/6} \left(\int_{t_k}^{t_k+r_k^2} \|u\|_{\infty}^{6/5} \right)^{5/6} \leq Cr_k^{1/3}$$

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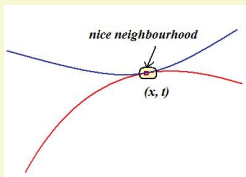
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Theorem 2 implies Theorem 1

Proof

- 1 since $u_0 \in H^{1/2}$ trajectories are unique on some interval $[0, \varepsilon)$
- 2 assume that for some $t > 0$ and $x \in \Omega$ (x, t) is not singular and there are two trajectories passing through (x, t)

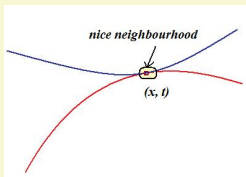


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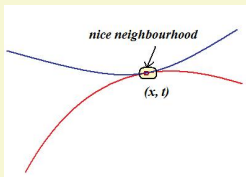


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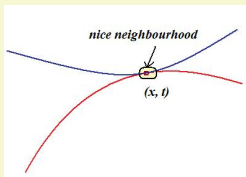


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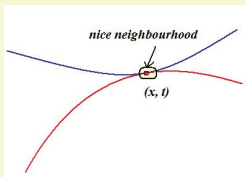


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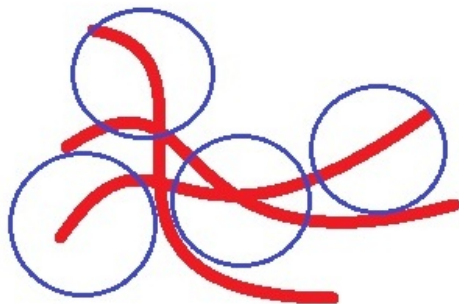
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Theorem 3. (James Robinson and WS 2009)

If u is a suitable weak solution with $p \in L^{5/3}((0, T) \times \Omega)$ corresponding to $u_0 \in H \cap H^{1/2}(\Omega)$ then almost every initial condition $a \in \Omega$ gives rise to a unique particle trajectory, which is a C^1 function of time.



box-counting dimension

$$N \sim cr^{-d}$$

Fractal dimension

Definition.

The upper box-counting dimension of a set X is given by

$$\limsup_{\epsilon \rightarrow 0} \frac{\log N(X, \epsilon)}{-\log \epsilon},$$

where $N(X, \epsilon)$ is

- 1 the minimum number of balls of radius ϵ required to cover X , or
- 2 the maximum number of ϵ -separated points in X .

Remark. We always have

$$d_H(X) \leq d_F(X)$$

and for some X $d_H(X) < d_F(X)$.

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Fractal dimension of a singular set

Theorem

The upper box-counting dimension of a putative singular set S is less or equal $5/3$.

Proof is based on Ladyzhenskaya-Seregin condition:

$$\frac{1}{r^2} \int_{Q_r(x,t)} |u|^3 + |p|^{3/2} \leq \varepsilon$$

from which we deduce that at a singular point we must have

$$\int_{Q_r(x,t)} |u|^{10/3} + |p|^{5/3} > cr^{5/3}.$$

Fractal dimension of a singular set

Theorem

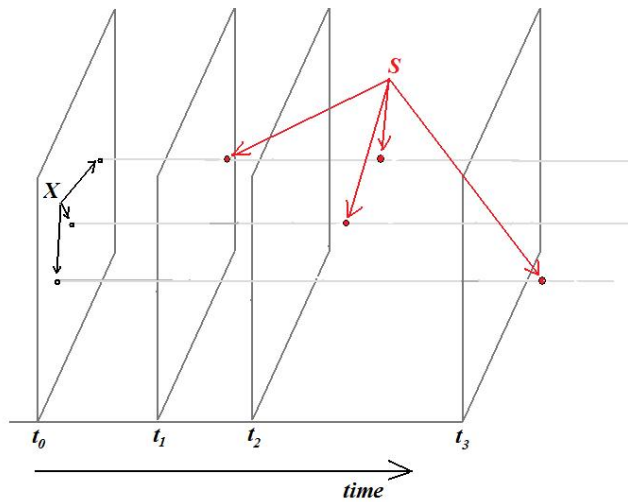
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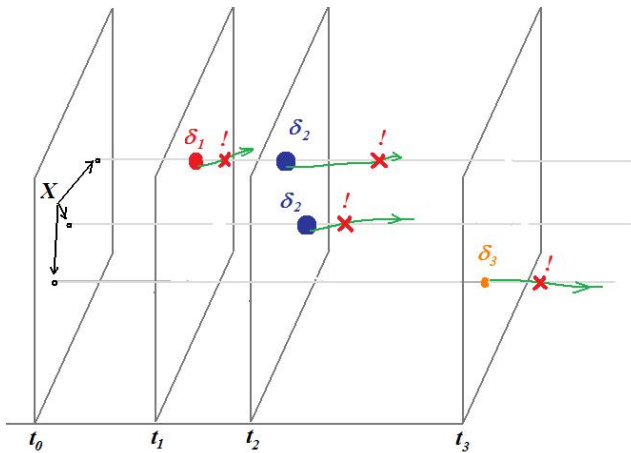
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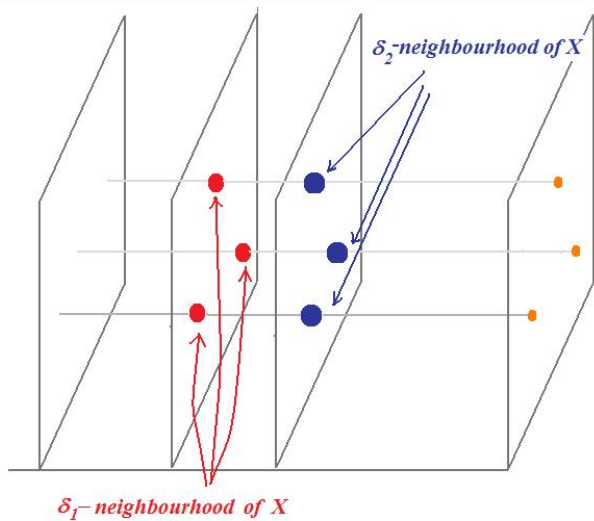
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Definition of deltas

$$\delta_k = \int_{t_k}^{t_{k+1}} \|u(s)\|_{\infty} ds$$



Proposition

If $X \in R^n$ has a box-counting dimension d , then for any $d' > d$ there exists an $\epsilon_0 > 0$ such that

$$\mu(O(X, \epsilon)) \leq c_n \epsilon^{n-d'} \text{ for all } 0 < \epsilon < \epsilon_0$$

Corollary

- 1 $V_1 =$ total volume of δ_1 -neighbourhood of $X \leq c\delta_1^{6/5}$
- 2 $V_2 =$ total volume of δ_2 -neighbourhood of $X \leq c\delta_2^{6/5}$
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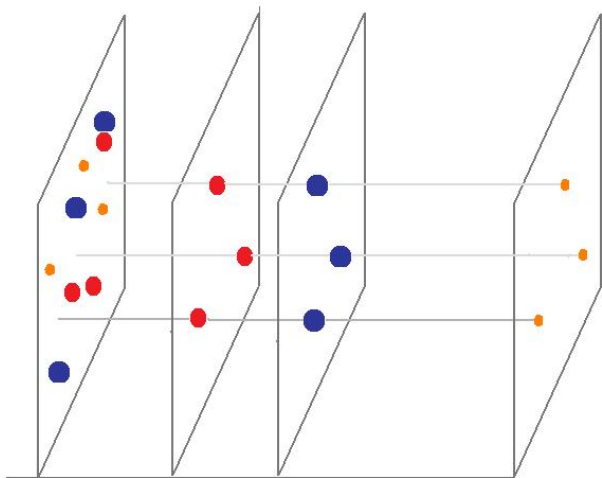
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$$\begin{aligned}\mu(K) &\leq V_1 + V_2 + V_3 + \dots + V_N \leq c(\delta_0^{6/5} + \delta_1^{6/5} + \delta_2^{6/5} + \dots + \delta_N^{6/5}) \leq \\ &\leq c[\epsilon^{1/5} \int_{t_0}^{t_1} \|u\|_\infty + \epsilon^{1/5} \int_{t_1}^{t_2} \|u\|_\infty + \dots + \epsilon^{1/5} \int_{t_{N-1}}^{t_N} \|u\|_\infty] \leq \\ &\leq c\epsilon^{1/5} \|u\|_{L^1(0,T;L^\infty(\Omega))}\end{aligned}$$

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- ① Aizenman M. *A sufficient condition for the avoidance of sets by measure preservng flows in R^n* , Duke Math. J. 45 809-812, 1978
- ② Caffarelli, Kohn and Nirenberg *Partial regularity of suitable weak solutions of the Navier-Stokes equations*, Commu. Pure Appl. Math. 35 771-831, 1982
- ③ Foias, Guillope, Temam, *Lagrangian representation of a flow*, Journal of Differential Equations 57, 440-449 (1985)
- ④ Robinson, Sadowski *A criterion for uniqueness of Lagrangian trajectories for weak solutions of the 3D Navier-Stokes equations*, Commun. Math. Phys., 290, 15-22 (2009)
- ⑤ Robinson, Sadowski *Almost everywhere uniqueness of Lagrangian trajectories for suitable weak solutions of the three dimensional Navier-Stokes equations*, Nonlinearity 22, 2093-2099, 2009.

Open problems



- 1 In what class $C^{1,\gamma}$ are the trajectories?
- 2 What about 'not suitable' weak solutions?
- 3 How is Lagrangian description related to GRP?

THANK YOU!

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